



Collisions

Ramsauer effect

Electron collisions

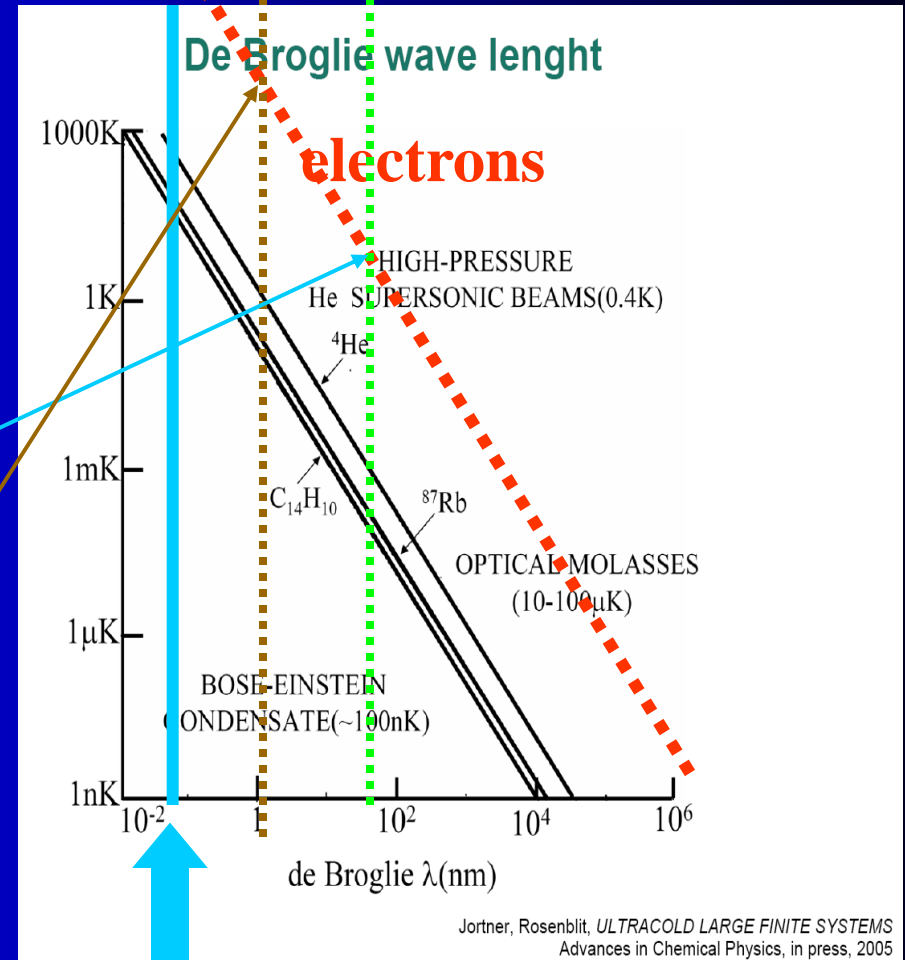
Low energy collisions with molecules

De Broglie wave length

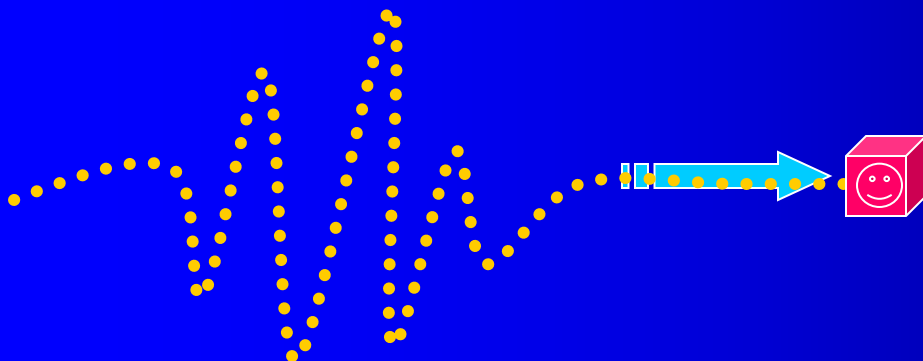
$$\lambda = \frac{h}{p} = \frac{h}{mv} \sqrt{1 - \frac{v^2}{c^2}}$$

$$\lambda_e(4K) \sim 540 \text{ \AA} \sim 54 \times 10^{-9} \text{ m}$$

$$\lambda_e(1\text{eV}) \sim 11.6 \text{ \AA} \sim 1.16 \times 10^{-9} \text{ m}$$



a_0



Collisions of electrons with atoms

Classical or quantum approach?

Electron:

$$\begin{aligned} 1\text{eV} &\rightarrow v = 5.9 \times 10^7 \text{ cm s}^{-1} \\ \tau &\sim a_0/v \sim 10^{-8} / 5.9 \times 10^7 = 2 \times 10^{-16} \text{ s} \\ \lambda &\sim 2\text{\AA} = 2 \times 10^{-8} \text{ cm} \text{ de Broglie} \end{aligned}$$

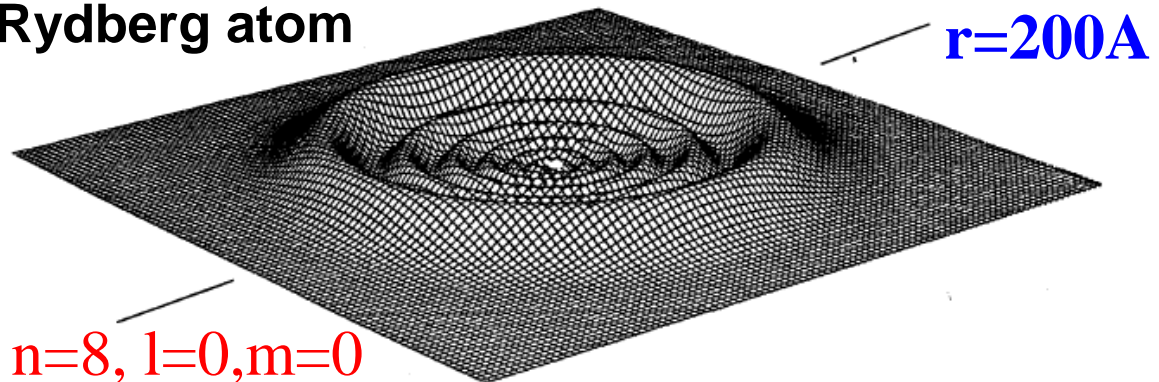
Ar+:

$$\begin{aligned} 1\text{eV} &\rightarrow v = 2 \times 10^5 \text{ cm s}^{-1} \\ \tau &\sim a_0/v \sim 10^{-8} / 2 \times 10^5 \sim 6 \times 10^{-14} \text{ s} \\ \lambda &\sim 9 \times 10^{-11} \text{ cm} \text{ de Broglie} \end{aligned}$$

$\text{H}_3^* + e$ at 10 K ????

$$\lambda_e(4\text{K}) \sim 540 \text{ \AA} \sim 54 \times 10^{-9} \text{ m}$$

Rydberg atom



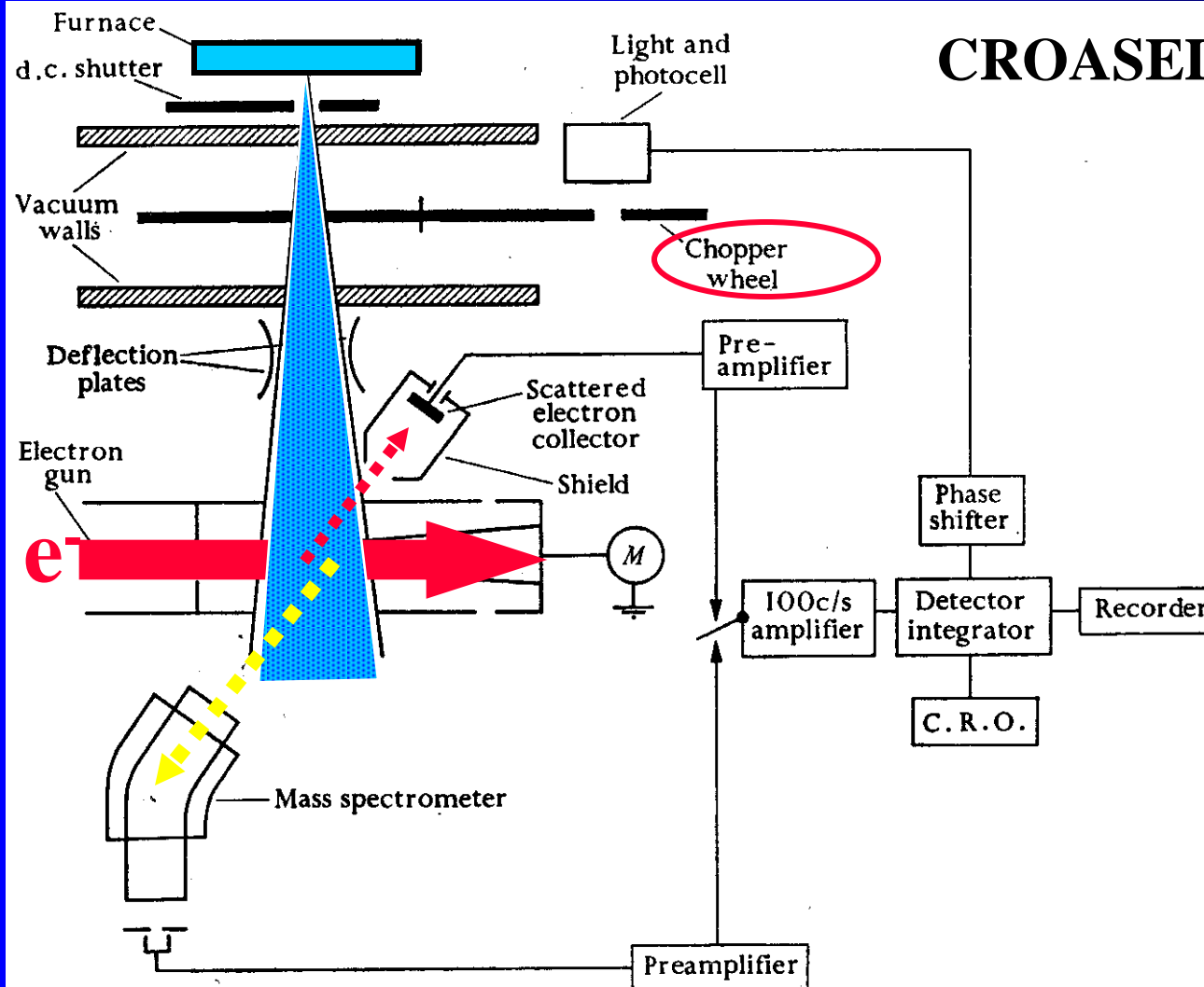
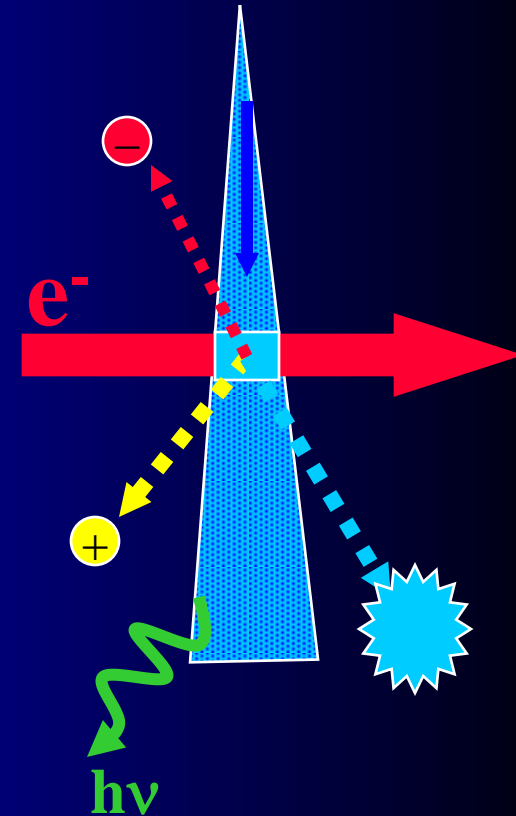


FIG. 1.2. Schematic diagram of the arrangement of apparatus used by Fite, Brackmann, and Neynaber for observation of elastic scattering of electrons by atomic hydrogen.

CROASED BEAM METHOD



Position (angle), mass and energy sensitive detectors

Partial cross section for excitation

Collisions with e

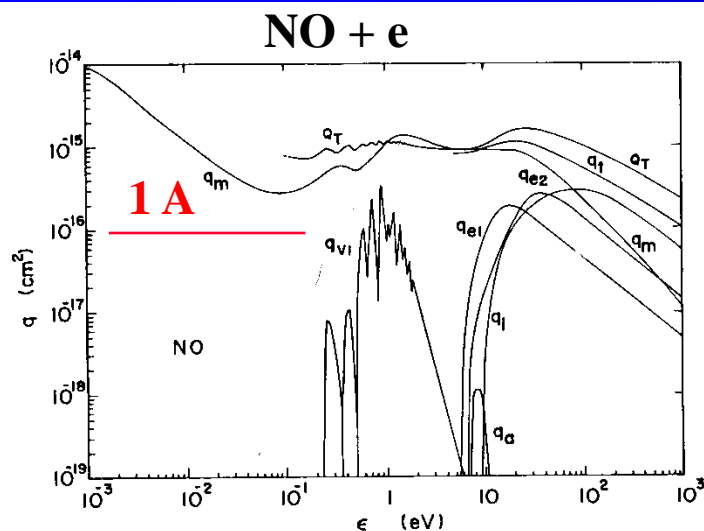


Fig. 4. Cross-section set for NO (1986).

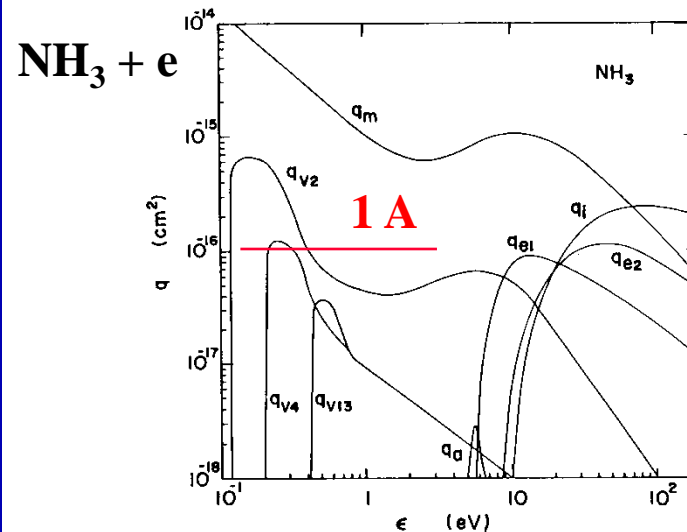


Fig. 6. Electron collision cross-section set for NH_3 (1986).

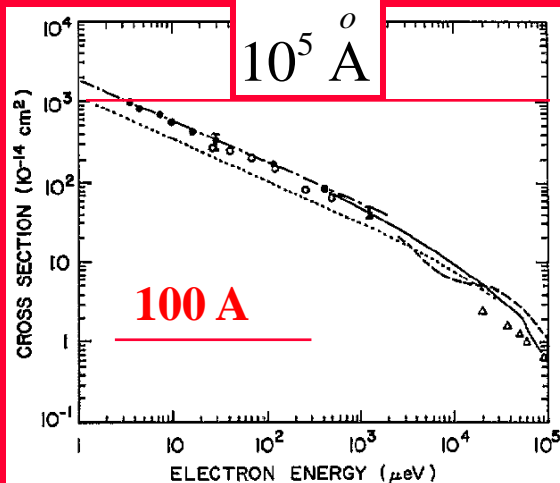
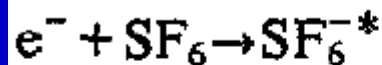
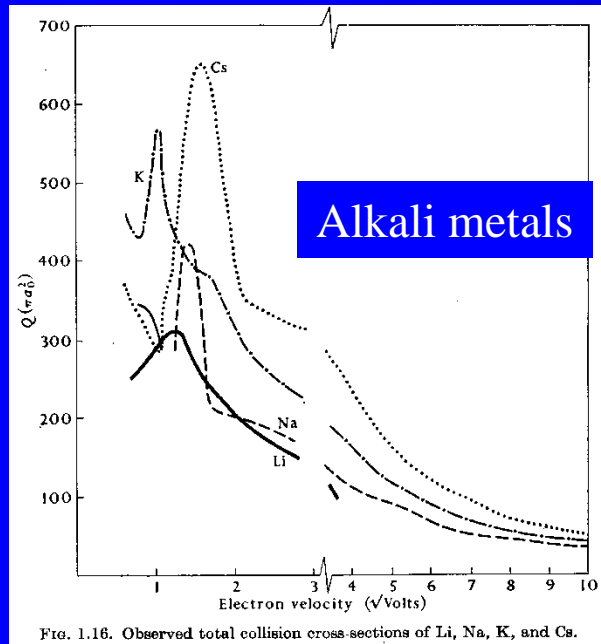


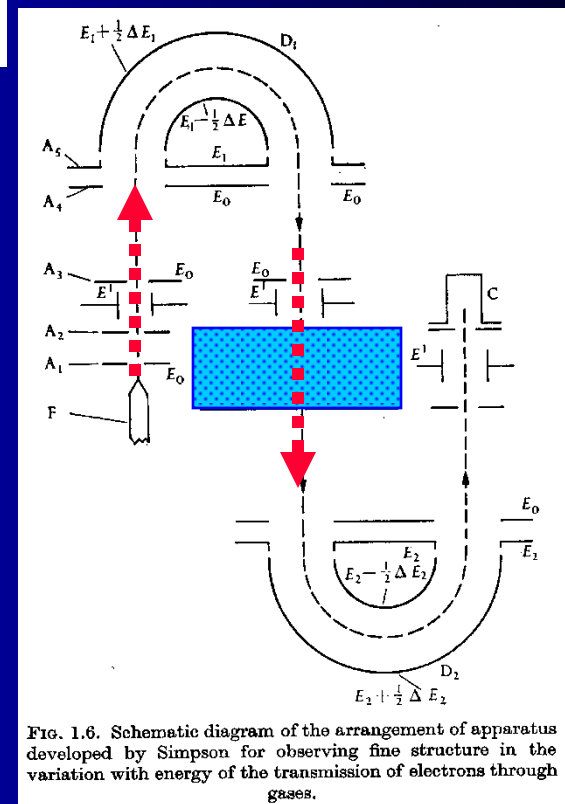
Figure 3. Cross sections for electron attachment to CCl_4 . \bullet , $\bar{\sigma}_e\text{-K}(np)$; $-\cdot-$, $\sigma_e(v)\text{-K}(np)$ (Frey *et al* 1994b); \circ , $\bar{\sigma}_e\text{-K}(np)$ (Ling *et al* 1992); $---$, free electrons (Hotop 1994); $----$, free electrons (Orient *et al* 1989); Δ , free electrons (Christodoulides and Christophorou (1971); $-----$, theory (Klots 1976).



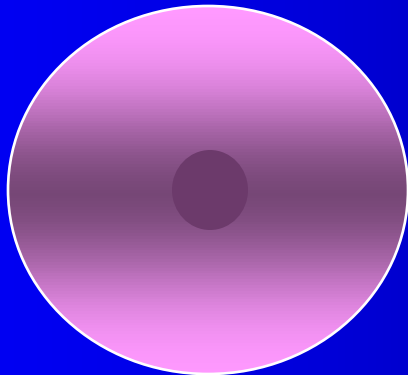
Total collision cross sections Na, K, Cs...



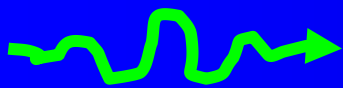
Collisions with e



Cs



$e^-_{(v)}$



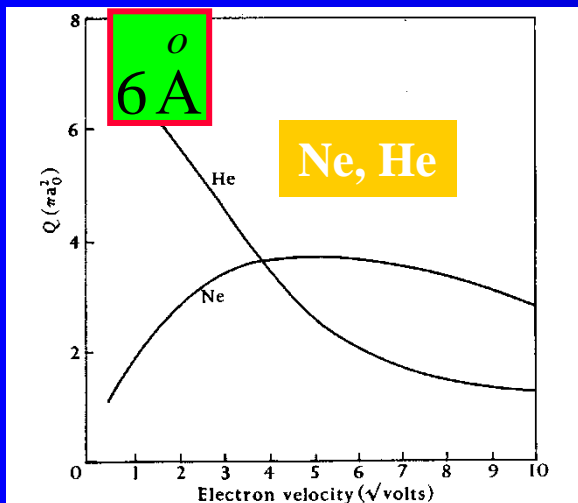


FIG. 1.10. Observed total collision cross-sections of He and Ne.

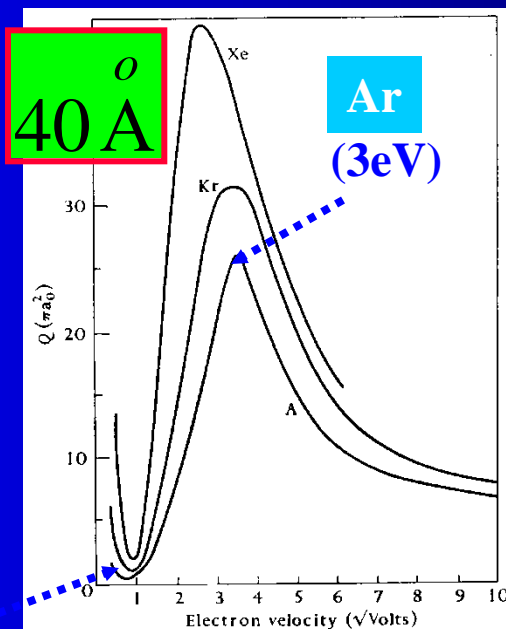


FIG. 1.9. Observed total collision cross-sections of Ar, Kr, and Xe.

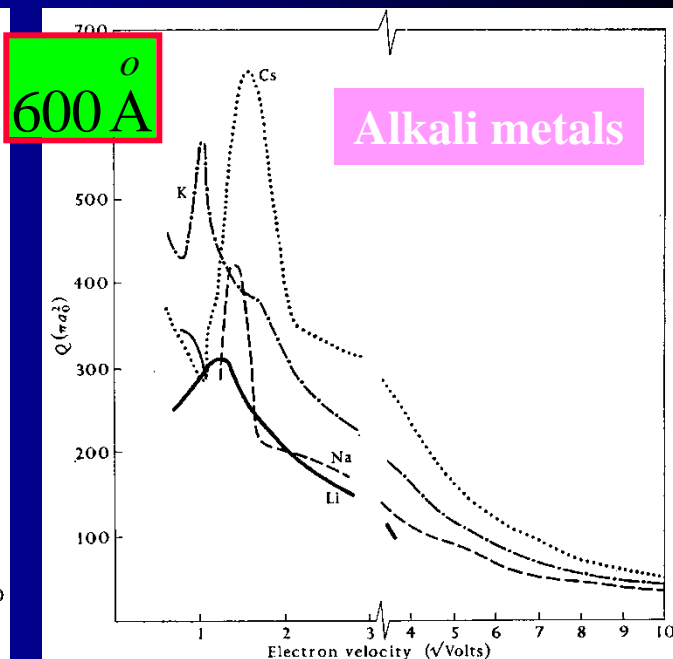


FIG. 1.16. Observed total collision cross-sections of Li, Na, K, and Cs.

(0,3eV)

$\sigma(v)$

Ar
(3eV)

Ne

Ar

(0.3eV)

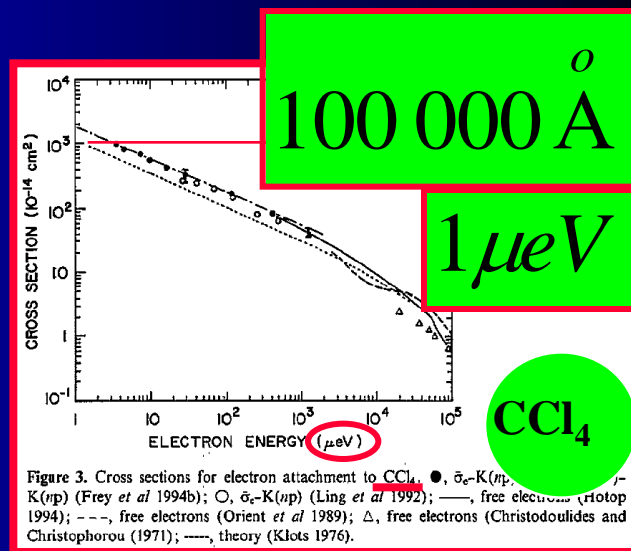
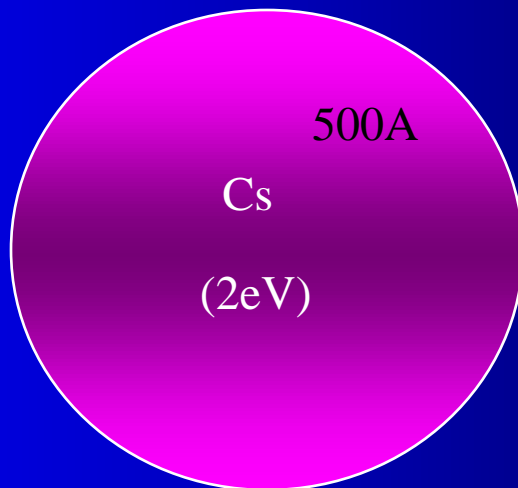


Figure 3. Cross sections for electron attachment to CCl_4 . \bullet , $\sigma_{\text{e}}\text{-K}(np)$; \circ , $\sigma_{\text{e}}\text{-K}(np)$ (Frey *et al* 1994b); Δ , $\sigma_{\text{e}}\text{-K}(np)$ (Ling *et al* 1992); —, free electrons (Hofop 1994); ---, free electrons (Orient *et al* 1989); Δ , free electrons (Christodoulides and Christophorou (1971); —, theory (Klots 1976).

Collisions of electrons with atoms – Ramsauer's method

Lenard 1903

Akesson 1916

Ramsauer 1921

ATTENUATION METHOD

$$\delta I = -N\sigma I_p \delta x$$

$$I_p = I_0 \exp(-\sigma N x)$$

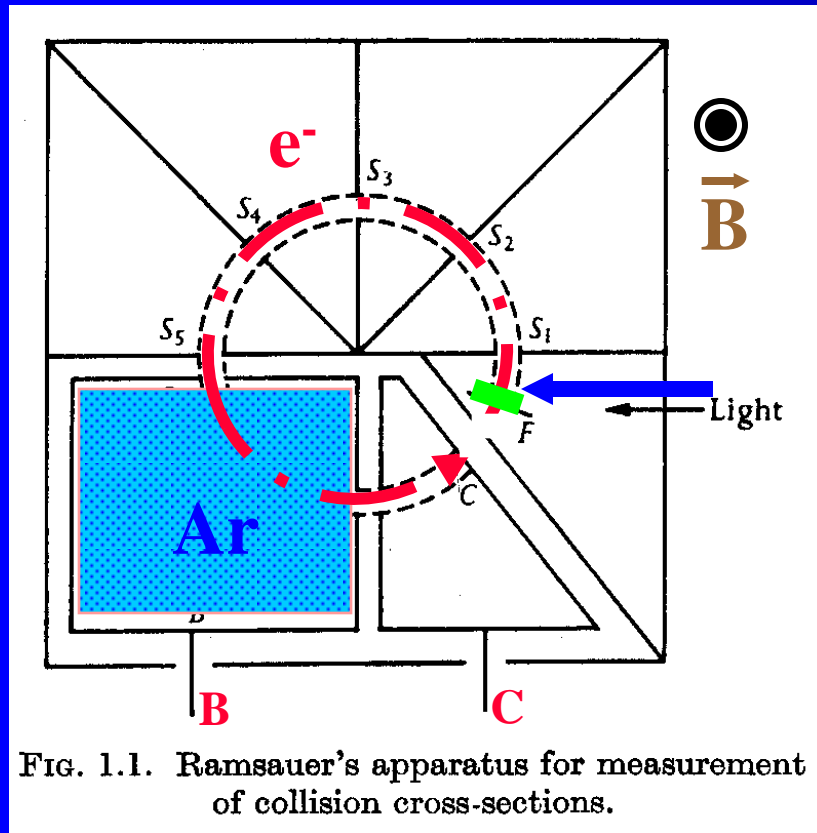
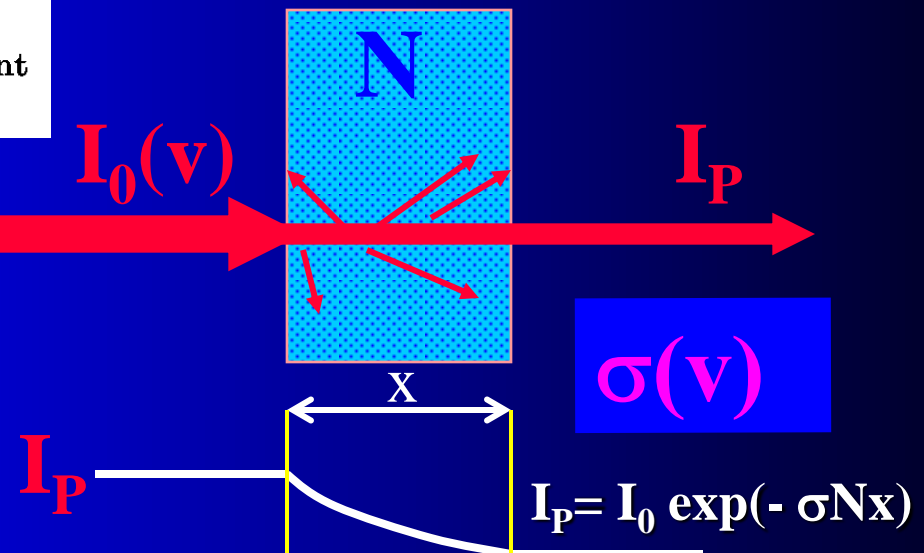
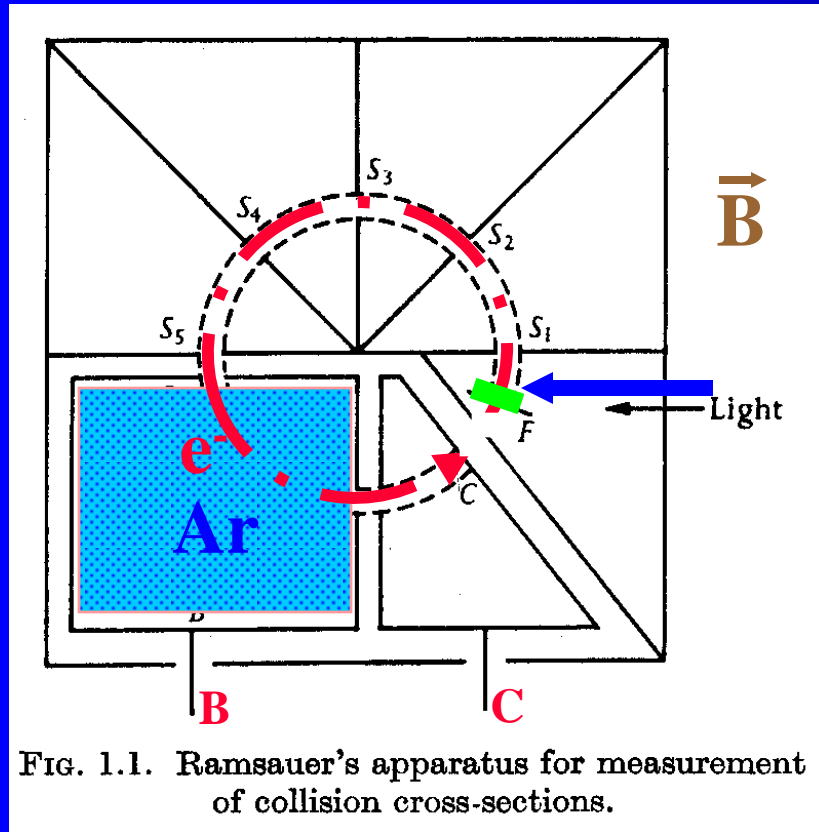


Photo cathode

Mono energetic electrons



Collisions of electrons with atoms – Ramsauer's method



Lenard 1903
Akesson 1916
Ramsauer 1921

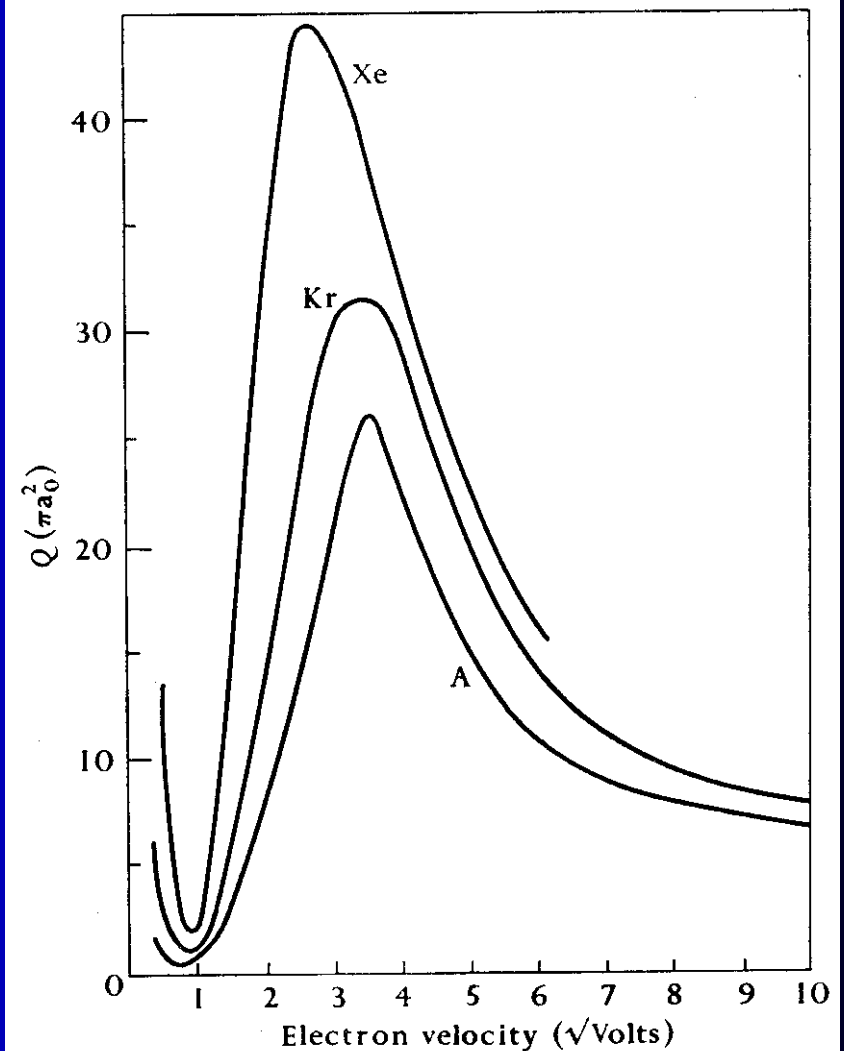


FIG. 1.9. Observed total collision cross-sections of A, Kr, and Xe.

Total collision cross section – e/atoms

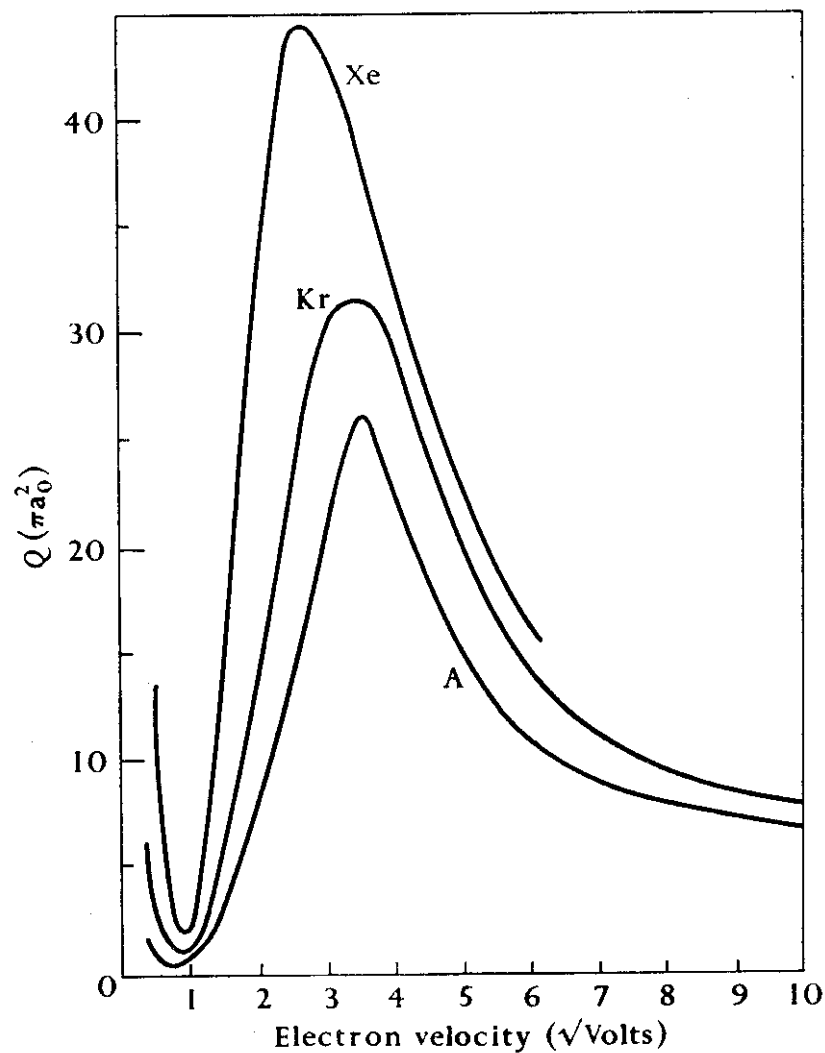


FIG. 1.9. Observed total collision cross-sections of A, Kr, and Xe.

$a_0 = 0.53 \times 10^{-8} \text{ cm} \sim 0.5 \text{ \AA}$

Radius of the first Bohr orbit of H atom

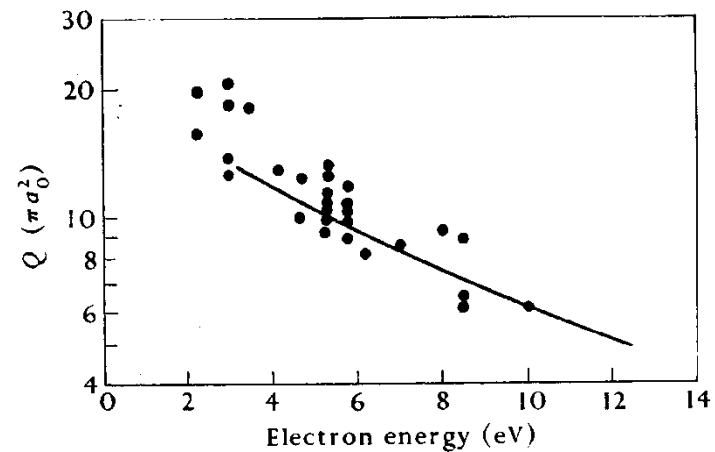


FIG. 1.11. Total collision cross-sections of atomic hydrogen. ● observed by Brackmann, Fite, and Neynaber; — observed by Neynaber, Marino, Rothe, and Trujillo.

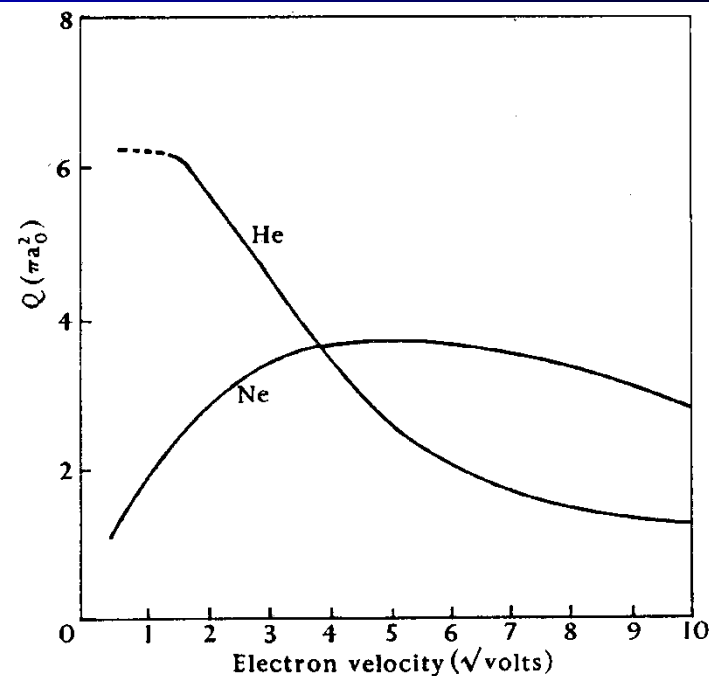


FIG. 1.10. Observed total collision cross-sections of He and Ne.

Details of Ramsauer effect

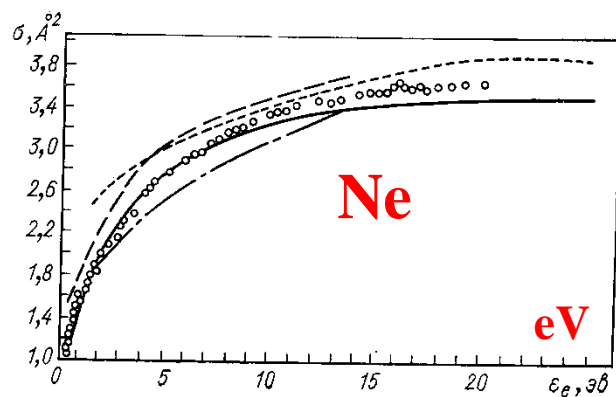


Рис. 5.8. Полное сечение рассеяния электрона на атоме неона.

Эксперимент (метод Рамзауэра): \circ — [101]; — [29]; — [92]; — [95]. Теория: — — — [109].

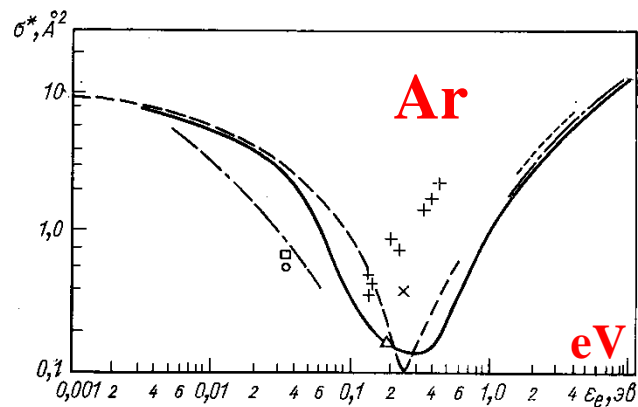


Рис. 5.9. Диффузионное сечение столкновения электрона с атомом аргона.

Эксперимент (подвижность электронов при малых полях и температурах): — [21]; — [47]; \times — [60]; \circ — [91]; \square — [112]; \triangle — [44]; — — — [16]; — — — — [108]; + — [43]. Теория: — — — [87].

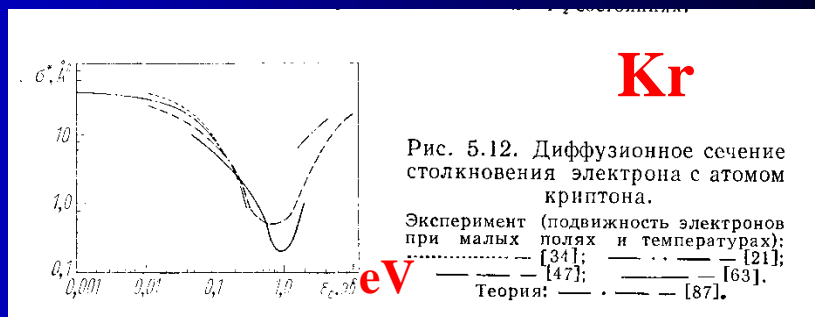


Рис. 5.12. Диффузионное сечение столкновения электрона с атомом криптона.

Эксперимент (подвижность электронов при малых полях и температурах): — [34]; — — — [21]; — — — [47]; — — — [63]. Теория: — — — [87].



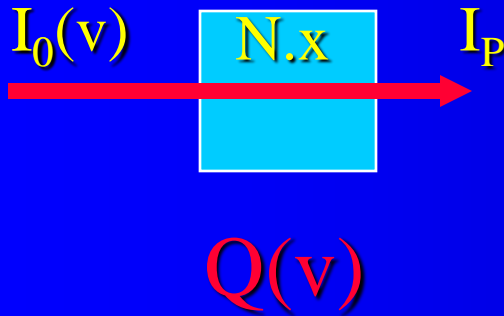
Рис. 5.3. Диффузионное сечение столкновения электрона с атомом гелия.

Эксперимент (подвижность электронов при малых полях и температурах): \square — [39]; \triangle — [73]; — — — [88]; — — — [91]; — [58]; — — — [13]; — — — [62]. Теория: — — — [75]; — — — [32]; — — — — расчет по формуле (5.37).

Frequencies of elastic collisions

$$\delta I = -NQI_p \delta x$$

$$I_p = I_0 \exp(-QNx)$$



$$a_0 = 0.53 \times 10^{-8} \text{ cm} \sim 0.5 \text{ \AA}$$

Radius of the first Bohr orbit of H atom

$$v \sim n v \sigma$$

Collision Frequencies

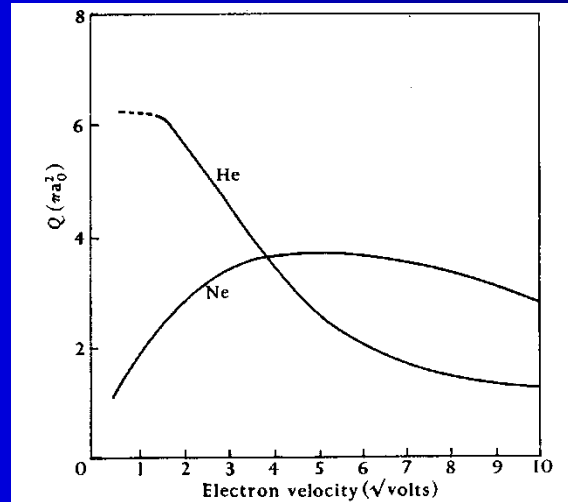


FIG. 1.10. Observed total collision cross-sections of He and Ne.

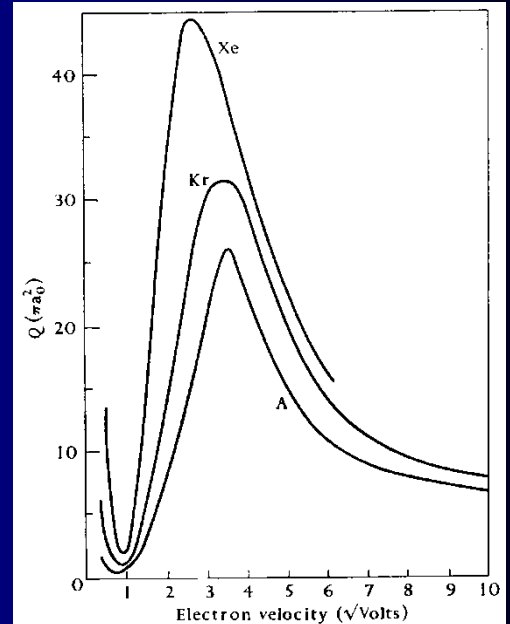
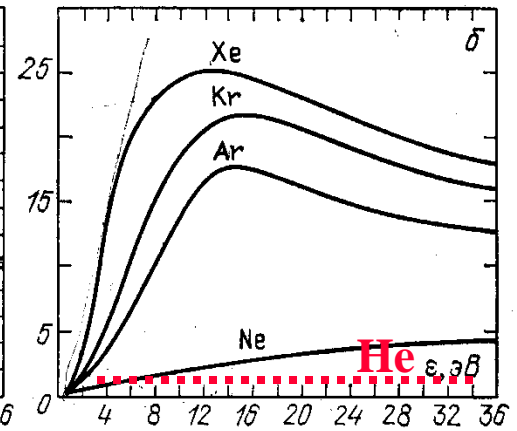
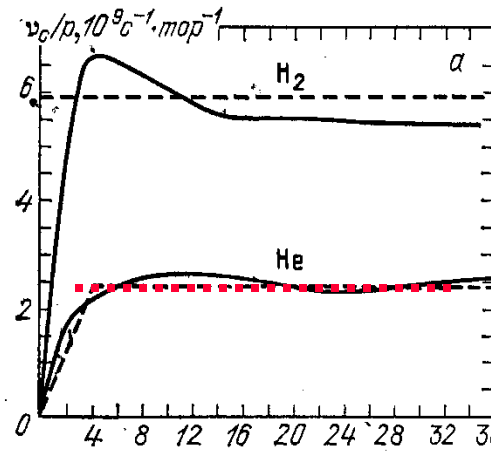


FIG. 1.9. Observed total collision cross-sections of Ar, Kr, and Xe.



Р и с. 2.5. Частоты упругих столкновений электронов, $p=1$ топ: а — в H_2 и He; б — в инертных газах; штриховые линии — удобная аппроксимация при расчетах [24]

Total collision and reactive cross sections comparison

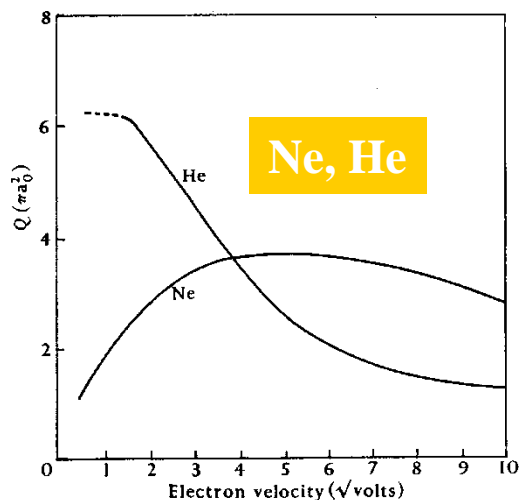


FIG. 1.10. Observed total collision cross-sections of He and Ne.

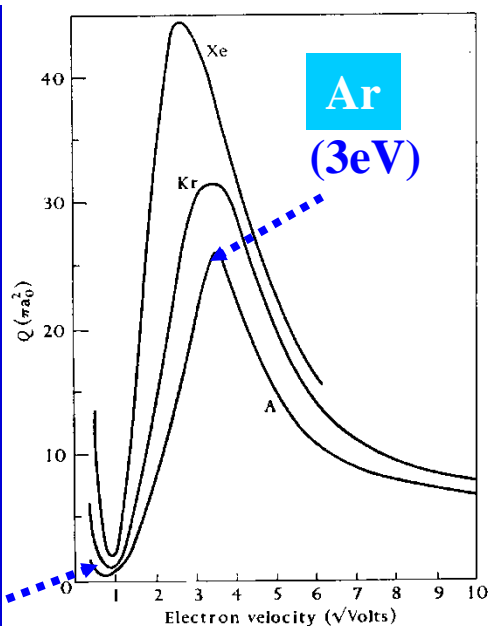


FIG. 1.9. Observed total collision cross-sections of Ar, Kr, and Xe.

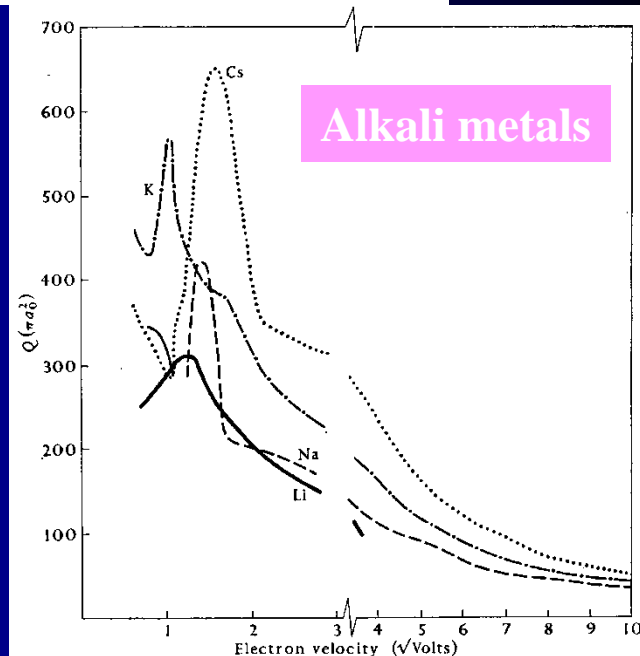


FIG. 1.16. Observed total collision cross-sections of Li, Na, K, and Cs.

$\sigma(v)$

(0,3eV)

Ar
(3eV)

Ne

Ar

(0.3eV)

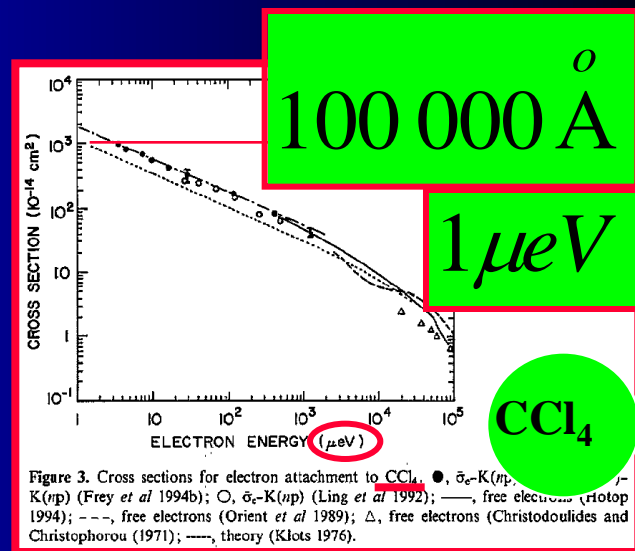
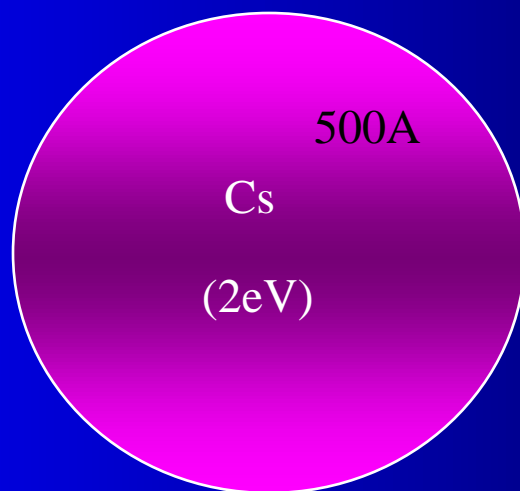
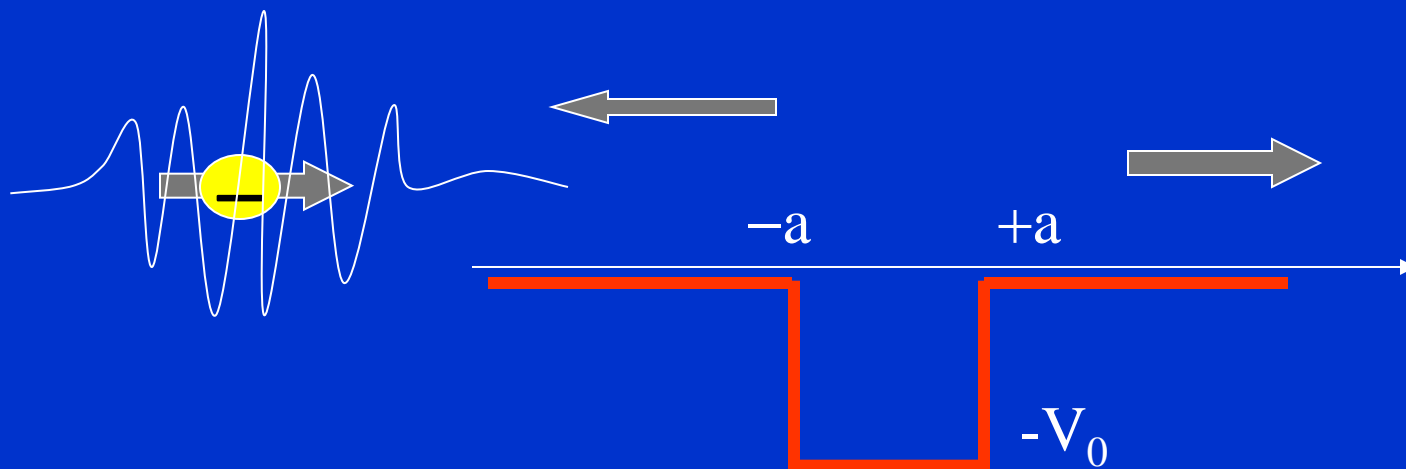


Figure 3. Cross sections for electron attachment to CCl_4 . \bullet , $\delta_e\text{-K}(np)$; \circ , $\delta_e\text{-K}(np)$ (Frey *et al* 1994b); Δ , free electrons (Hofop 1994); $---$, free electrons (Orient *et al* 1989); Δ , free electrons (Christodoulides and Christophorou (1971); $---$, theory (Klots 1976).

Kvantová mechanika

Jednorozměrný rozptyl



Kvantová mechanika I

J. Klíma B. Velický

MFF 1992

Jednorozměrný rozptyl

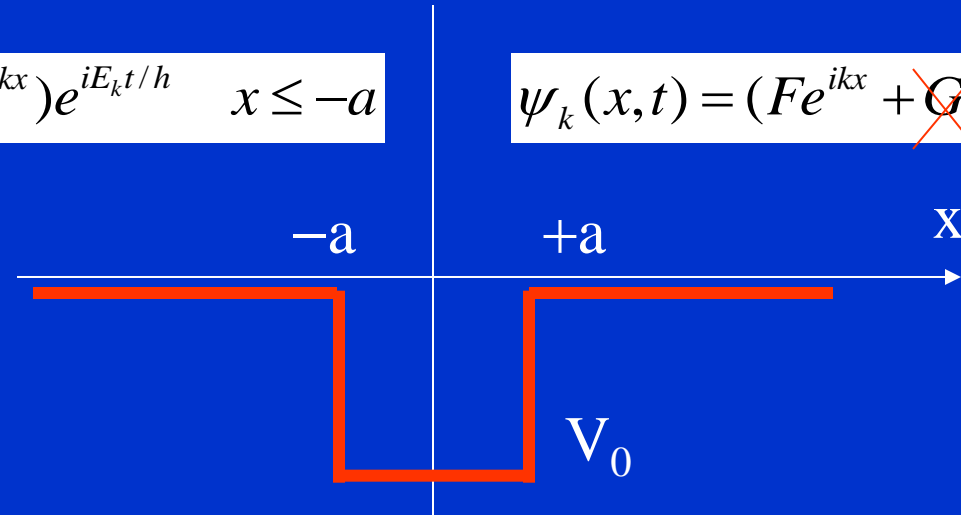
Vlnová funkce má tvar superposice Brogliových vln

$$k = \sqrt{2mE / \hbar^2}$$

$$\psi_k(x, t) = (Ae^{ikx} + Be^{-ikx})e^{iE_k t / \hbar} \quad x \leq -a$$

$$k = \sqrt{2mE / \hbar^2}$$

$$\psi_k(x, t) = (Fe^{ikx} + \cancel{G}e^{-ikx})e^{iE_k t / \hbar} \quad x > a$$



$$\psi_k(x, t) = (Ce^{ik'x} + De^{-ik'x})e^{iE_k t / \hbar} \quad |x| \leq a$$

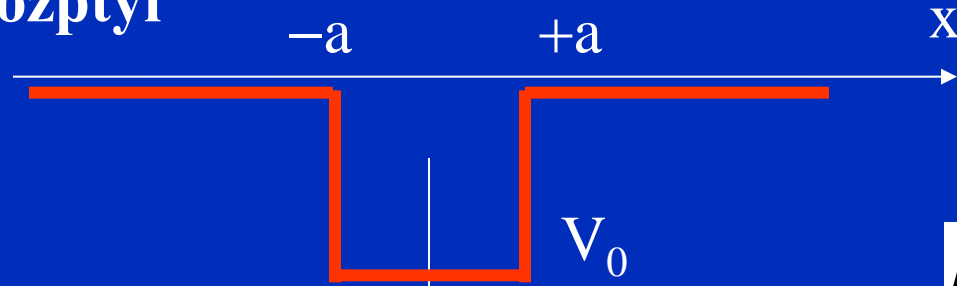
$$k' = \sqrt{2m(E + V_0) / \hbar^2}$$

a) dopadající částice $\rightarrow A$

b) odražená částice $\rightarrow B$

c) procházející částice $\rightarrow F \neq 0, G = 0$

Jednorozměrný rozptyl



$$k = \sqrt{2mE / \hbar^2}$$

$$k' = \sqrt{2m(E + V_0) / \hbar^2}$$

$$\psi_k(x, t) = (Ae^{ikx} + Be^{-ikx})e^{iE_k t / \hbar} \quad x \leq -a$$

$$\psi_k(x, t) = (Ce^{ik'x} + De^{-ik'x})e^{iE_k t / \hbar} \quad |x| \leq a$$

$$\psi_k(x, t) = (Fe^{ikx})e^{iE_k t / \hbar} \quad x > a$$

Parametry jsou **E**, **V₀**, **a**

Hladkost řešení v bodech $\pm a$
Urči konstanty B, C, D, G,
Hodnota A je vstupní parametr

Tok dopadajících částic

$$j_{in} = \frac{\hbar k}{m} |A|^2$$

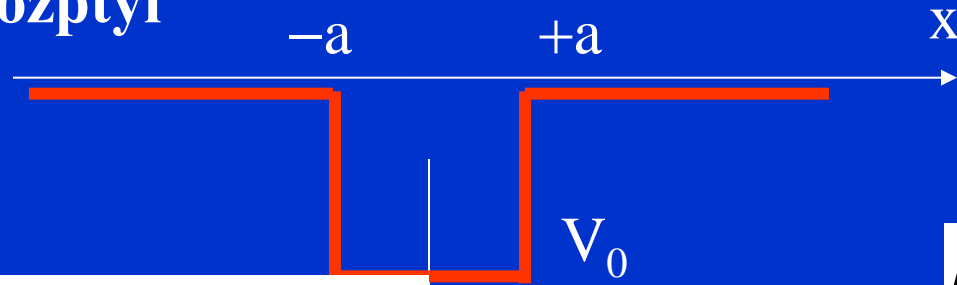
Tok odražených částic

$$j_{rf} = \frac{\hbar k}{m} |B|^2$$

Tok prošlých částic

$$j_{tr} = \frac{\hbar k}{m} |F|^2$$

Jednorozměrný rozptyl



$$k = \sqrt{2mE / \hbar^2}$$

$$k' = \sqrt{2m(E + V_0) / \hbar^2}$$

$$\psi_k(x, t) = (Ae^{ikx} + Be^{-ikx})e^{iE_k t / \hbar} \quad x \leq -a$$

$$\psi_k(x, t) = (Ce^{ik'x} + De^{-ik'x})e^{iE_k t / \hbar} \quad |x| \leq a$$

$$\psi_k(x, t) = (Fe^{ikx})e^{iE_k t / \hbar} \quad x > a$$

Parametry jsou **E**, **V₀**, **a**

Hladkost řešení v bodech ±a
Urči konstanty B, C, D, G,
Hodnota A je vstupní parametr

Tok dopadajících částic

$$j_{in} = \frac{\hbar k}{m} |A|^2$$

$$C = \frac{F}{2} \left(1 + \frac{k}{k'}\right) e^{i(k-k')a}$$

Tok odražených částic

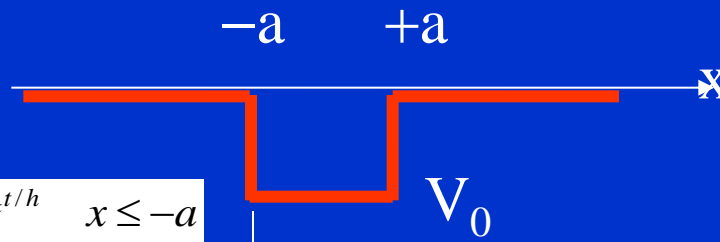
$$j_{rf} = \frac{\hbar k}{m} |B|^2$$

$$D = \frac{F}{2} \left(1 - \frac{k}{k'}\right) e^{i(k+k')a}$$

Tok prošlých částic

$$j_{tr} = \frac{\hbar k}{m} |F|^2$$

Jednorozměrný rozptyl



$$k = \sqrt{2mE} / \hbar$$

$$k' = \sqrt{2m(E + V_0)} / \hbar$$

$$\psi_k(x,t) = (Ae^{ikx} + Be^{-ikx})e^{iE_k t / \hbar} \quad x \leq -a$$

$$\psi_k(x,t) = (Ce^{ik'x} + De^{-ik'x})e^{iE_k t / \hbar} \quad |x| \leq a$$

$$\psi_k(x,t) = (Fe^{ikx})e^{iE_k t / \hbar} \quad x > a$$

Parametry jsou **E**, **V₀**, **a**

$$j_{in} = \frac{\hbar k}{m} |A|^2$$

$$j_{rf} = \frac{\hbar k}{m} |B|^2$$

$$j_{tr} = \frac{\hbar k}{m} |F|^2$$

Hladkost řešení v bodech $\pm a$
Urči konstanty B, C, D, F,
Hodnota A je vstupní parametr

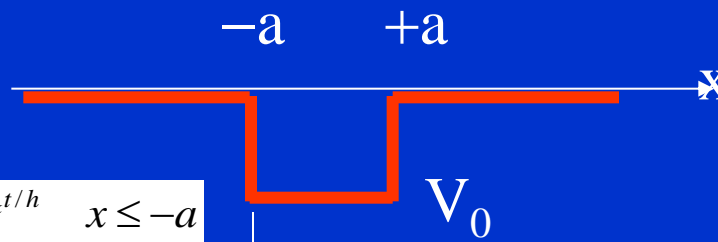
$$\varepsilon = \frac{k'}{k} + \frac{k}{k'}$$

$$A = e^{2ika} (\cos(2k'a) - i(\varepsilon/2) \sin(2k'a)) F$$

Koeficient průchodu T, koeficient odrazu R

$$\frac{1}{T} = \left| \frac{A}{F} \right|^2 = 1 + \frac{V_0^2}{4E(E + V)} \sin^2(2k'a)$$

Jednorozměrný rozptyl



$$k = \sqrt{2mE} / \hbar$$

$$k' = \sqrt{2m(E + V_0)} / \hbar$$

$$\psi_k(x,t) = (Ae^{ikx} + Be^{-ikx})e^{iE_k t / \hbar} \quad x \leq -a$$

$$\psi_k(x,t) = (Ce^{ik'x} + De^{-ik'x})e^{iE_k t / \hbar} \quad |x| \leq a$$

$$\psi_k(x,t) = (Fe^{ikx})e^{iE_k t / \hbar} \quad x > a$$

Parametry jsou **E**, **V₀**, **a**

$$j_{in} = \frac{\hbar k}{m} |A|^2$$

$$j_{rf} = \frac{\hbar k}{m} |B|^2$$

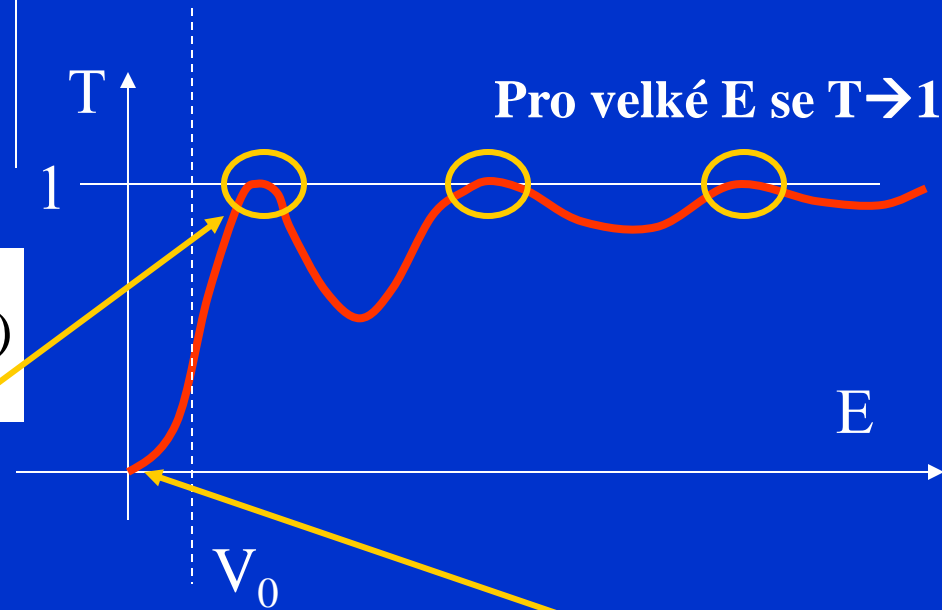
$$j_{tr} = \frac{\hbar k}{m} |F|^2$$

Koeficient průchodu T, koeficient odrazu R

$$\frac{1}{T} = \left| \frac{A}{F} \right|^2 = 1 + \frac{V_0^2}{4E(E + V_0)} \sin^2(2k'a)$$

T=1 pro

$$2k'_n a = n\pi$$



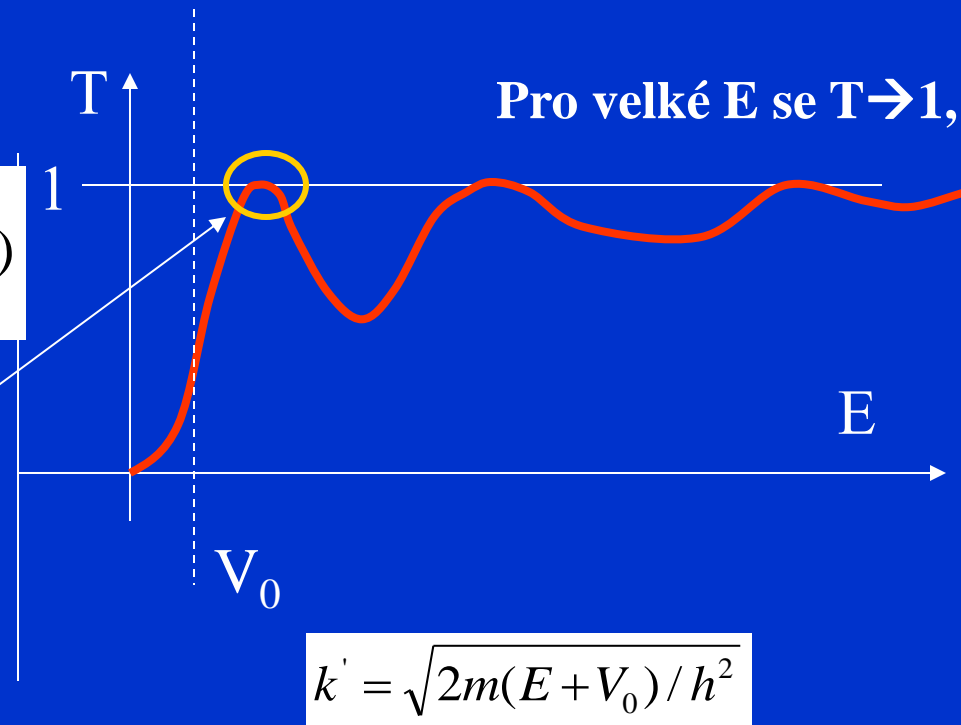
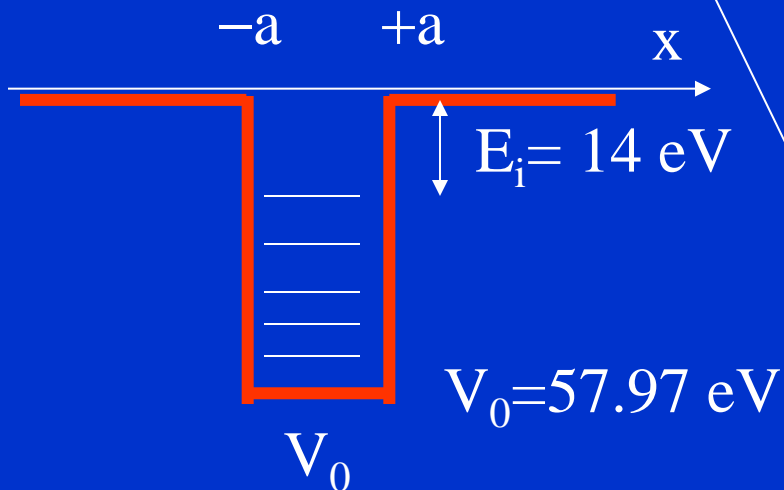
$$\lim_{E \rightarrow 0} \frac{1}{T} \sim 1 + \frac{V_0^2}{4EV_0} \sin^2(2k'a) \sim 1 + \frac{V_0}{4E} \sin^2(2\sqrt{2mV_0}/\hbar a) \sim 1 + \frac{V_0}{4E} \text{const} \sim \infty \quad \lim_{E \rightarrow 0} T \sim 0$$

Efekt Ramsauera - Kr

Parametry jsou **E**, **V₀**, **a**

$$\frac{1}{T} = \left| \frac{A}{F} \right|^2 = 1 + \frac{V_0^2}{4E(E + V_0)} \sin^2(2k'a)$$

T=1 pro $2k_n'a = n\pi$



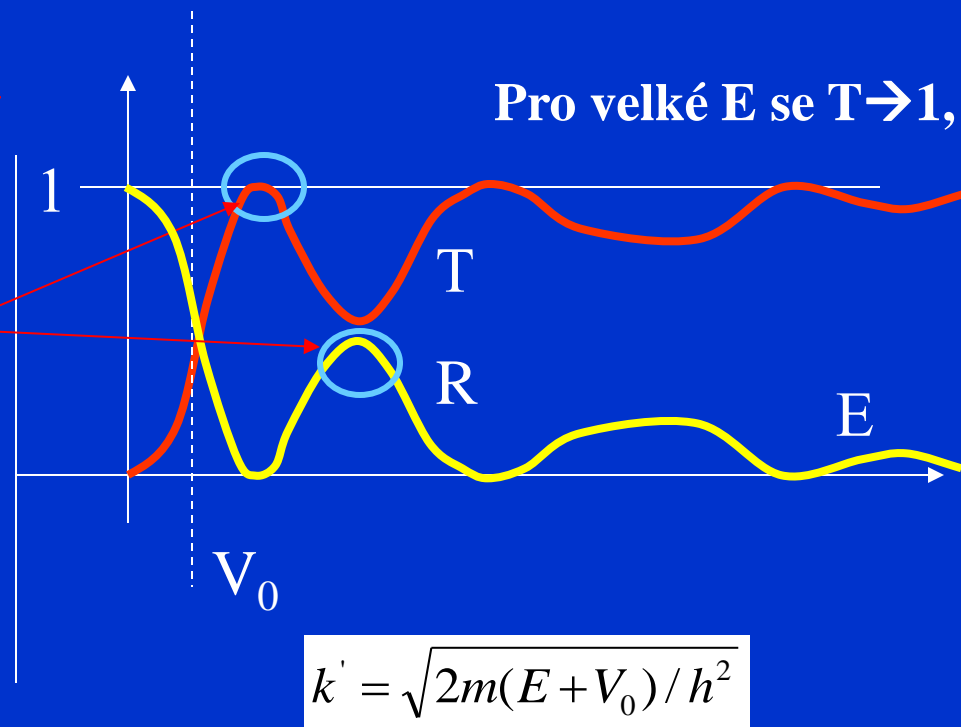
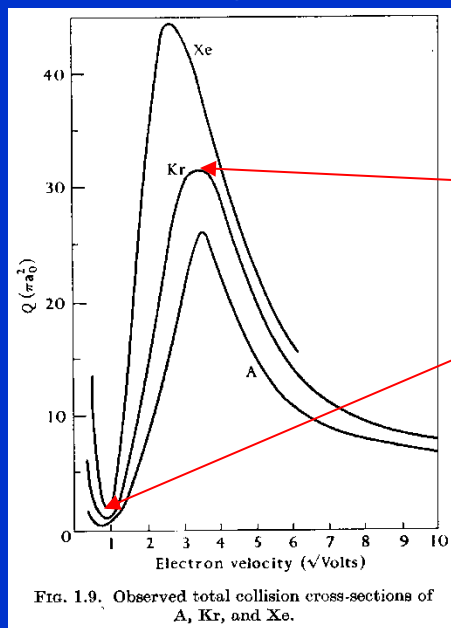
$$k' = \sqrt{2m(E + V_0) / \hbar^2}$$

Kr; $a = 2 \text{ \AA}$
 $E_i = 14 \text{ eV} \rightarrow V_0 = 57.97 \text{ eV}$

$E = 0.013 \text{ eV}$ $V_0 = 0.75 \text{ eV}$

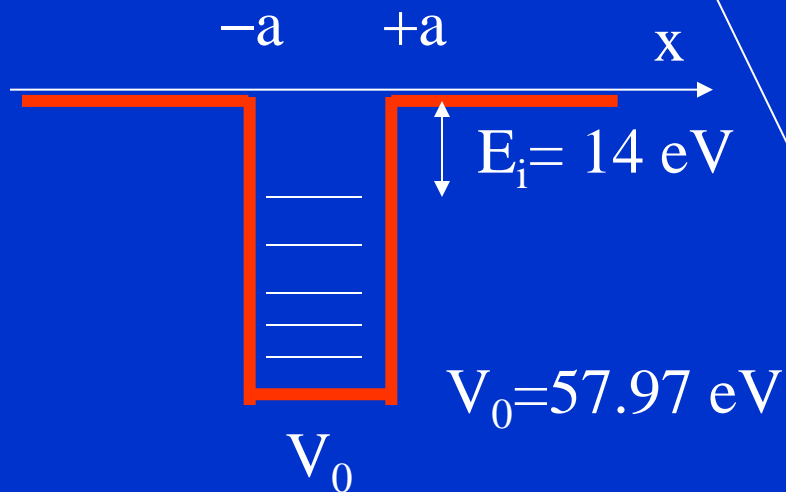
Jednorozměrný rozptyl

Parametry jsou E, V_0, a $T+R=1$



$$2k_n' a = n\pi$$

$$k' = \sqrt{2m(E + V_0) / \hbar^2}$$



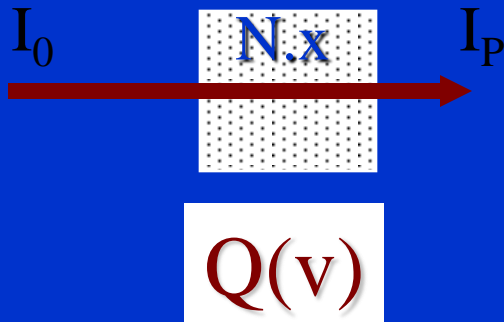
Kr; $a = 2 \text{ \AA}$
 $E_i = 14 \text{ eV} \rightarrow V_0 = 57.97 \text{ eV}$

$$E = 0.013 \text{ eV} \quad V_0 = 0.75 \text{ eV}$$

Frequencies of elastic collisions

$$\delta I = -NQI_p \delta x$$

$$I_p = I_0 \exp(-QNx)$$



**Collision
Frequencies**

$$\nu \sim n v \sigma$$

$$a_0 = 0.53 \times 10^{-8} \text{ cm} \sim 0.5 \text{ \AA}$$

Radius of the first Bohr orbit of H atom

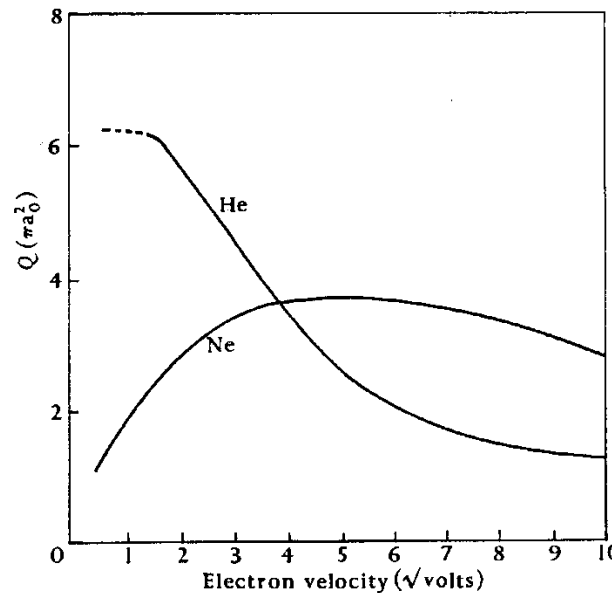


FIG. 1.10. Observed total collision cross-sections of He and Ne.

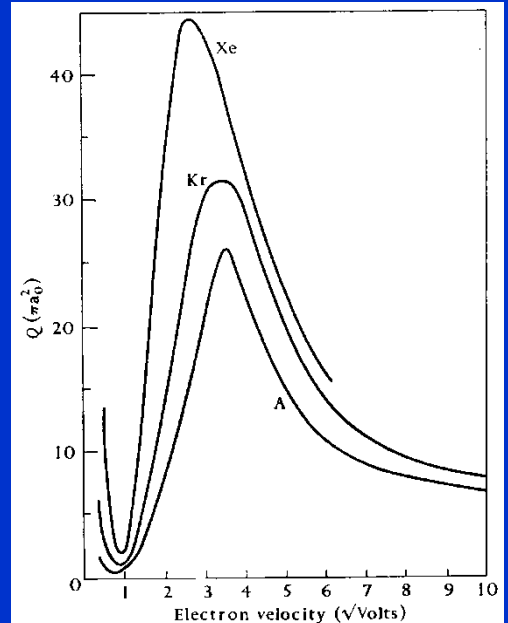
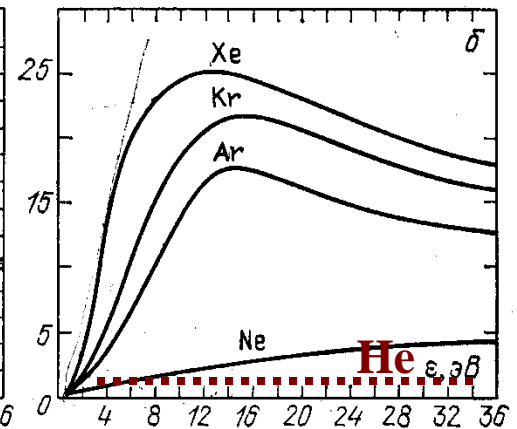
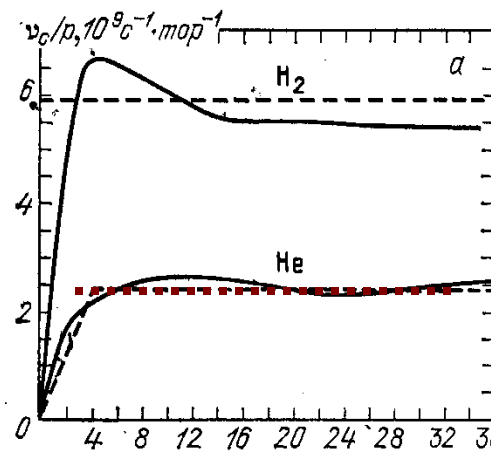


FIG. 1.9. Observed total collision cross-sections of A, Kr, and Xe.



Р и с. 2.5. Частоты упругих столкновений электронов, $p=1$ топ: а — в H_2 и He; б — в инертных газах; штриховые линии — удобная аппроксимация при расчетах [24]

Koniec rosprávký 27 10 2022

Very low collision energies

1995

TOPICAL REVIEW

Electron-molecule collisions at very low electron energies

F B Dunning

Department of Physics and the Rice Quantum Institute, Rice University, PO Box 1892, Houston, TX 77251, USA

J. Phys. B: At. Mol. Opt. Phys. 28 (1995) 1645-1672. Printed in the UK

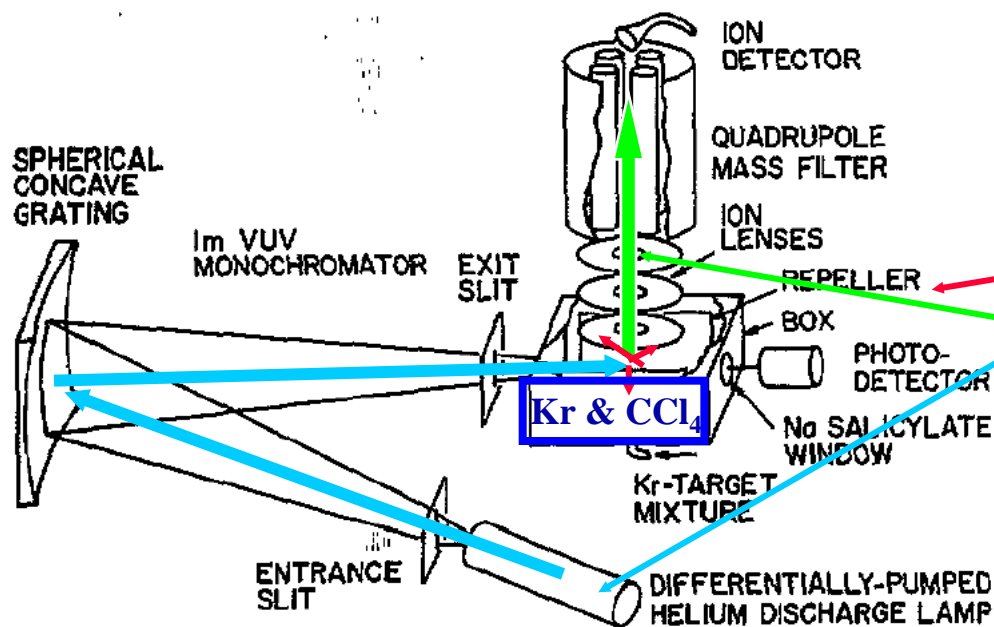
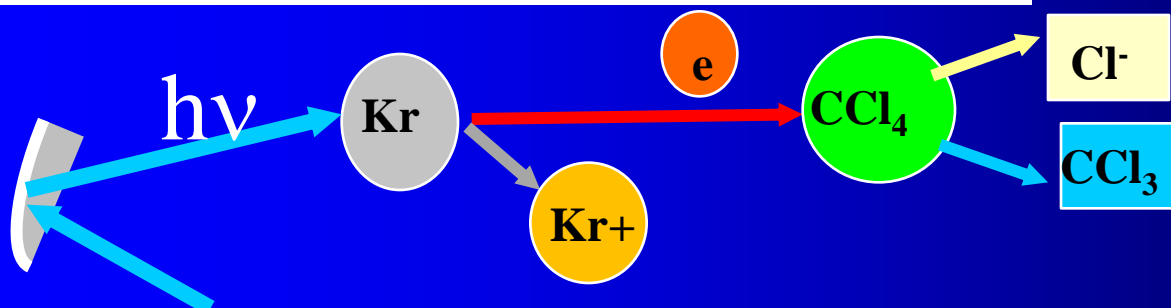
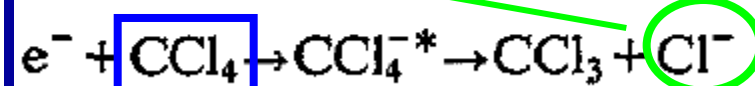
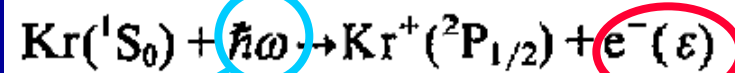


Figure 1. Schematic diagram of the VUV photoionization apparatus used for attachment studies (Chutjian and Alajajian 1985a, b).



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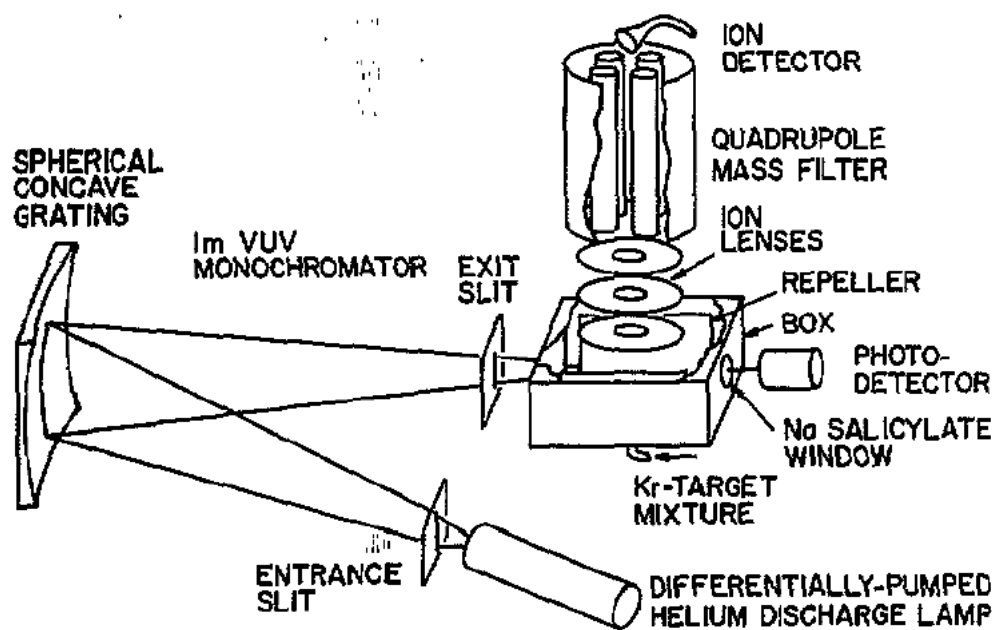
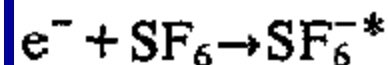


Figure 1. Schematic diagram of the VUV photoionization apparatus used for attachment studies (Chutjian and Alajajian 1985a, b).



Very low collision energies

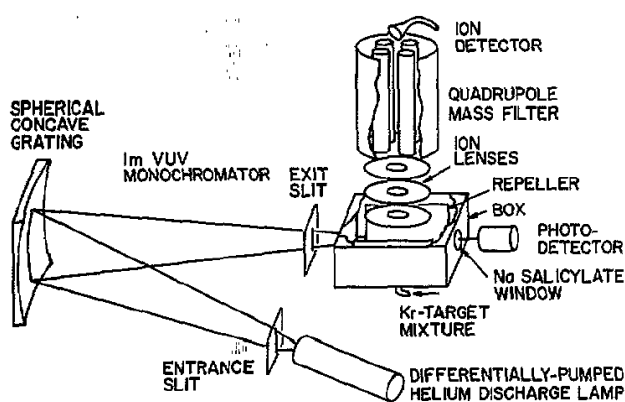
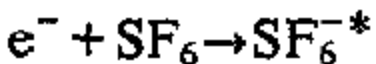
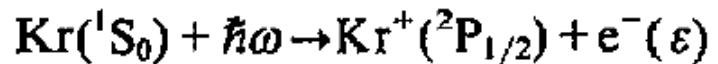


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TOPICAL REVIEW

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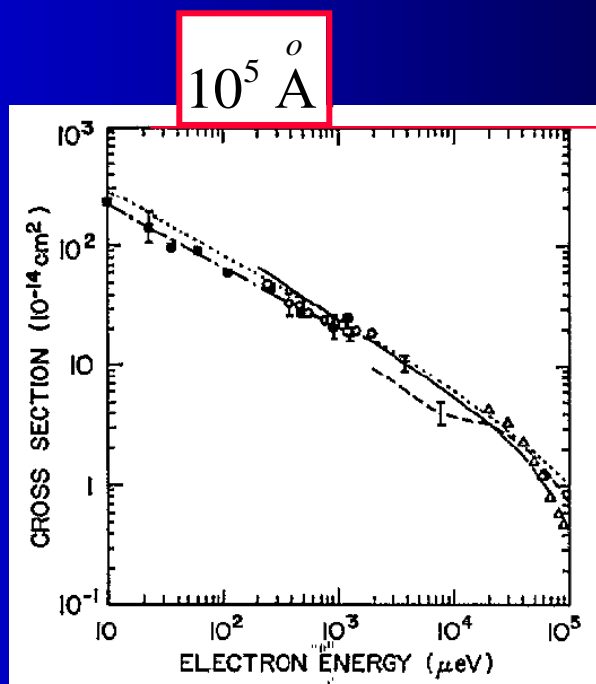


Figure 2. Cross section for electron attachment to SF_6 . ■, $\sigma_e\text{-K}(np)$; — · —, $\sigma_e(v)\text{-K}(np)$ (Ling *et al* 1992). ○, $\sigma_e\text{-Rb}(ns)$ (Zollars *et al* 1985); —, free electrons (Klar *et al* 1992a, b); ---, free electrons (Chutjian and Alajajian 1985); △, free electrons (Pai *et al* 1979, Chutjian and Alajajian 1985a); ----, theory (Klots 1976).

Electron attachment at very low electron energies

$10^5 \text{ } ^\circ \text{A}$

$10^5 \text{ } ^\circ \text{A}$

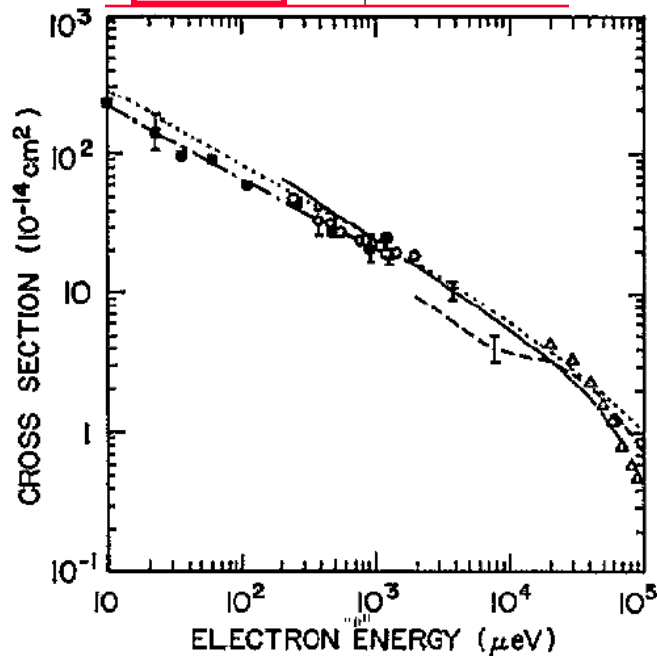


Figure 2. Cross section for electron attachment to SF_6 . ■, $\sigma_e\text{-K}(np)$; — · —, $\sigma_e(\nu)\text{-K}(np)$ (Ling *et al* 1992). ○, $\sigma_e\text{-Rb}(ns)$ (Zollars *et al* 1985); —, free electrons (Klar *et al* 1992a, b); ---, free electrons (Chutjian and Alajajian 1985); △, free electrons (Pai *et al* 1979, Chutjian and Alajajian 1985a); ----, theory (Klots 1976).

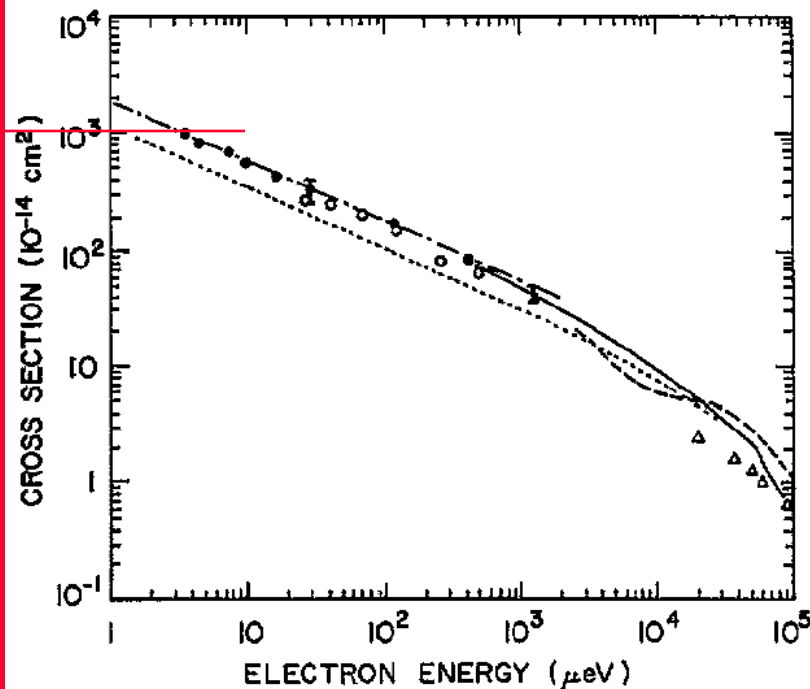


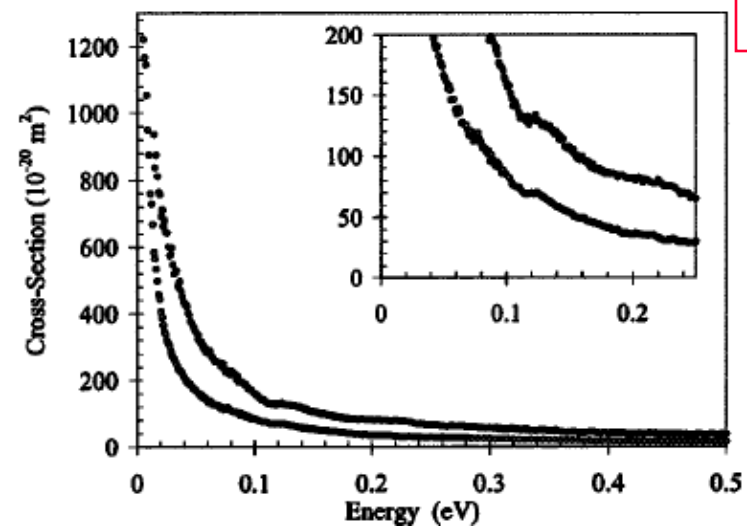
Figure 3. Cross sections for electron attachment to CCl_4 . ●, $\sigma_e\text{-K}(np)$; — · —, $\sigma_e(\nu)\text{-K}(np)$ (Frey *et al* 1994b); ○, $\sigma_e\text{-K}(np)$ (Ling *et al* 1992); —, free electrons (Hotop 1994); ---, free electrons (Orient *et al* 1989); △, free electrons (Christodoulides and Christophorou (1971); ----, theory (Klots 1976).

Cold electron scattering in SF₆ and C₆F₆: Bound and virtual state channels

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(Received 26 November 2003; published 20 May 2004)



The general process which we study involves electron attachment,

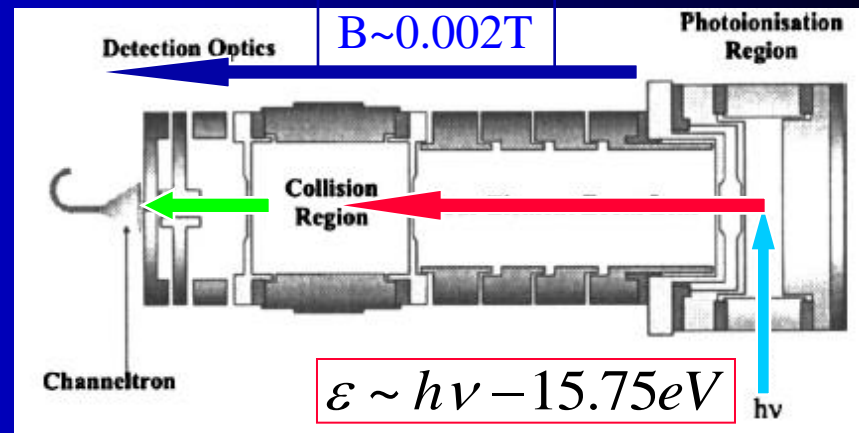
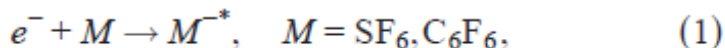
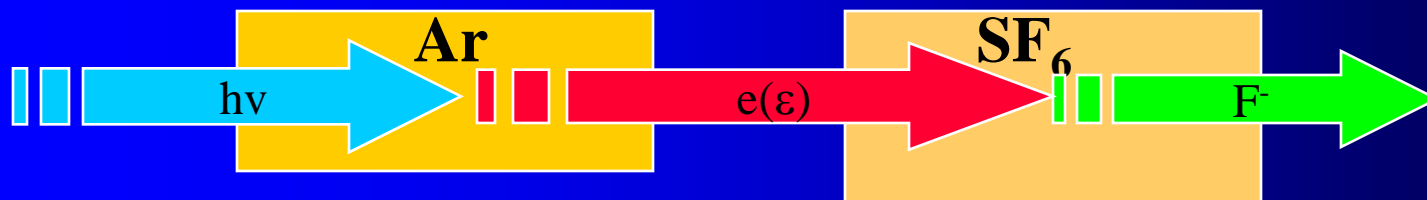
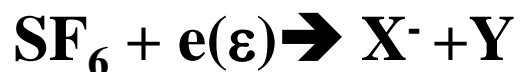


FIG. 1. A scale diagram of the apparatus. Monochromatic synchrotron radiation from ASTRID ($h\nu$) enters a photoionization region containing Ar. Photoelectrons, expelled by a weak electric field, are focused by a four-element lens [38] into a collision chamber containing the target gas. Transmitted electrons are detected at the channel electron multiplier (channeltron) situated beyond some further electron optics. The apparatus may be immersed in an axial magnetic field of 2×10^{-3} T.

B=0.002T



$$\epsilon \sim h\nu - 15.75 \text{ eV}$$



Scattering of cold electrons by ammonia, hydrogen sulfide, and carbonyl sulfide

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(Received 2 July 2008; published 29 October 2008)

Experimental data obtained with a high-resolution transmission experiment are presented for the scattering of electrons in the energy range 20 meV–10 eV for NH_3 , 25 meV–10 eV for H_2S , and 15 meV–2.5 eV for OCS. Data include cross sections for both integral scattering and scattering into the backward hemisphere, the latter up to 650 meV impact energy, with an electron energy resolution of between 1.6 and 3.5 meV. The new data allow the first detailed comparison with theory for the energy regime dominated by rotationally inelastic and elastic scattering for these species. It is evident that theory still lacks quantitative predictive power at low energy, although qualitative agreement is consistently good for all three species. A discussion is given of the possible presence of a virtual state in OCS scattering as recently proposed on theoretical grounds.

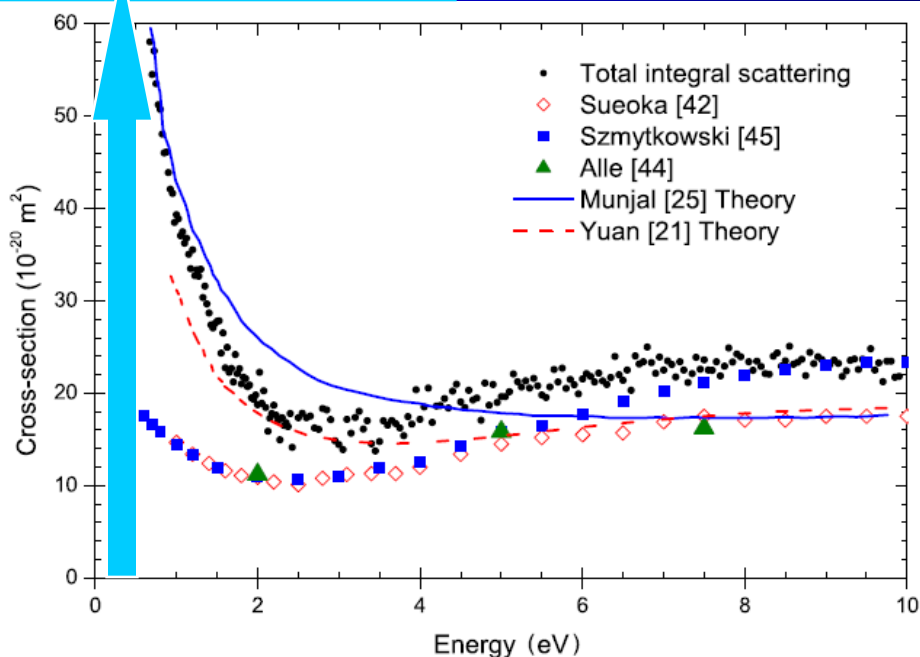
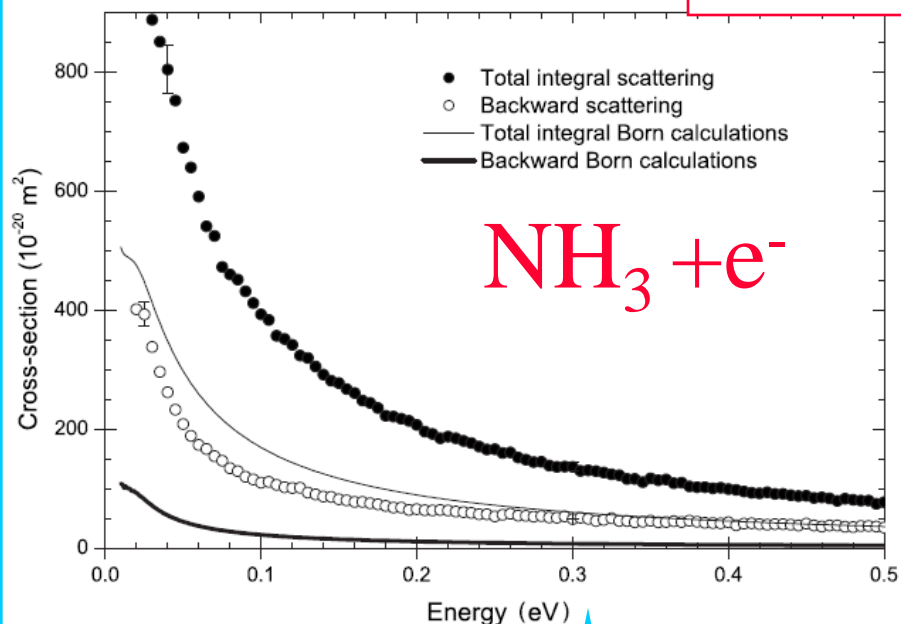


FIG. 1. (Color online) NH_3 : the variation of the sum of the integral elastic and inelastic cross sections, $\sigma_{T,I}$, between 0.675 and 10 eV. Also shown are experimental data from Sueoka *et al.* [42], Szymkowski *et al.* [45], and Alle *et al.* [44] and theoretical values from Munjal *et al.* [25] and Yuan *et al.* [21].

Molecules cross section for interaction with electrons

2008

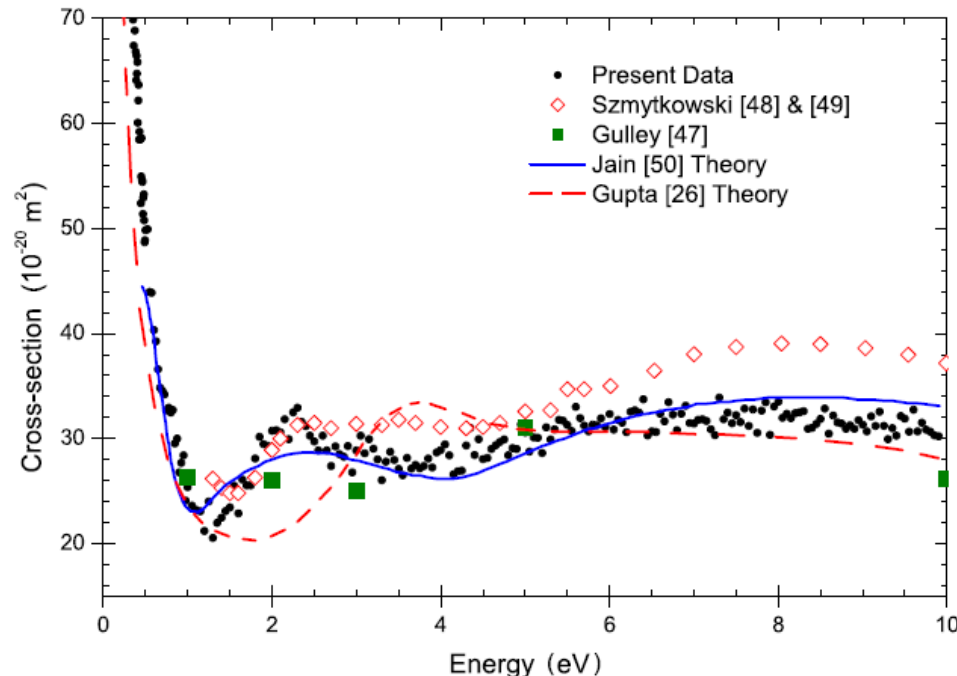


FIG. 3. (Color online) H_2S : the variation of the sum of the integral elastic and inelastic cross sections, $\sigma_{T,I}$, between 380 meV and 10 eV. Also shown are experimental data in Szmytkowski *et al.* [48,49] and Gulley *et al.* [47] and theoretical values from Jain *et al.* [50] and Gupta *et al.* [26].

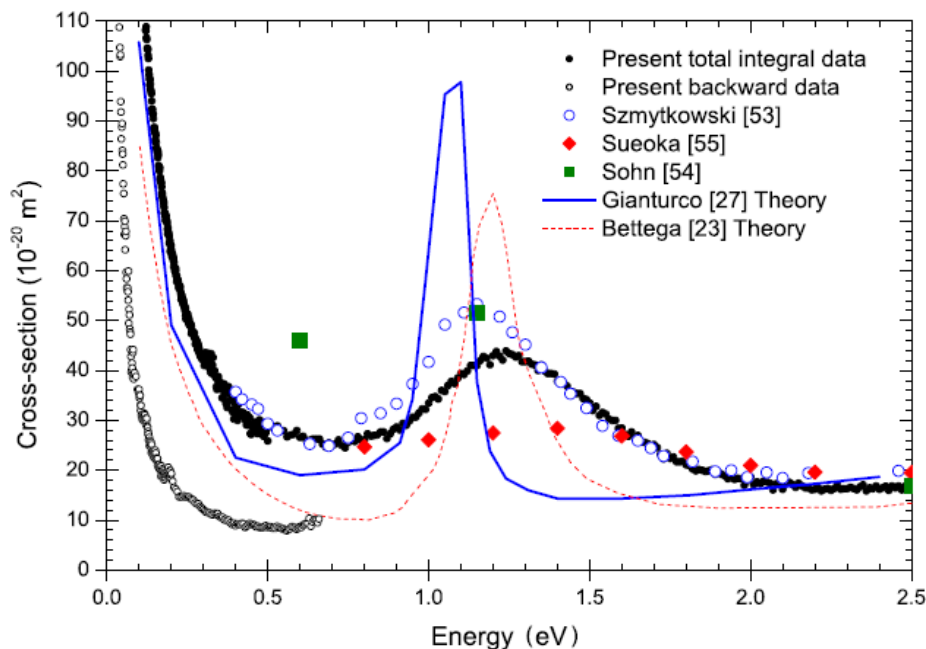
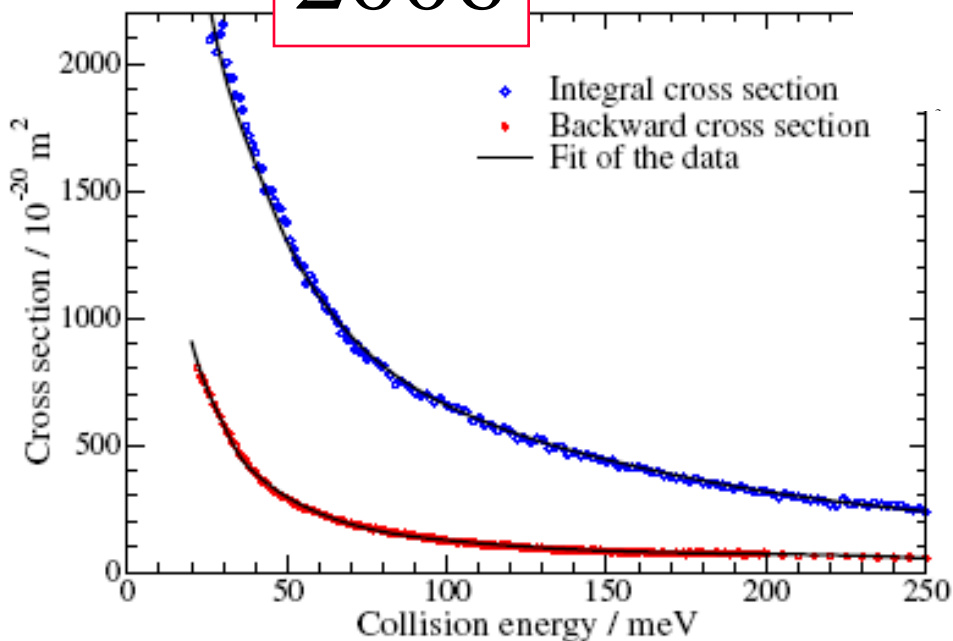
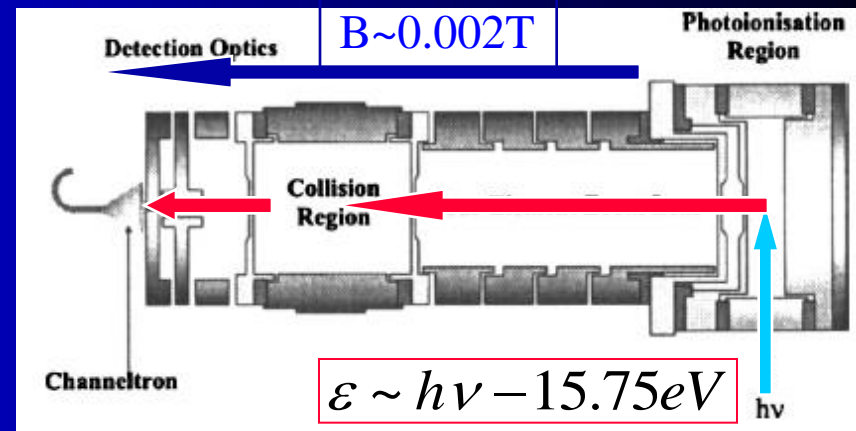
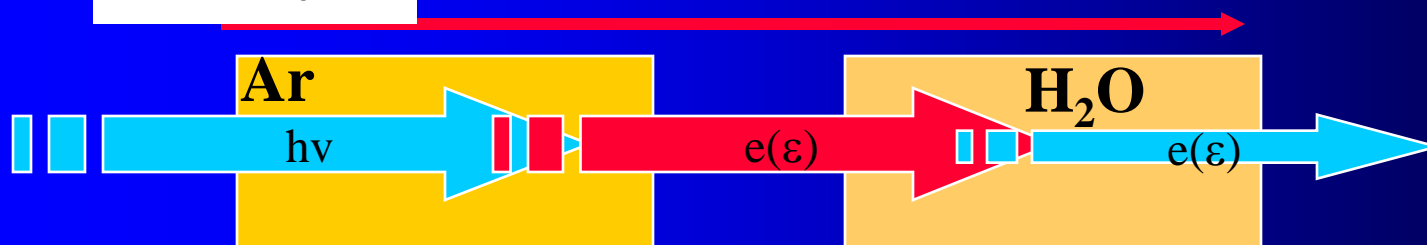


FIG. 5. (Color online) OCS : the variation of the sum of the integral elastic and inelastic cross sections, $\sigma_{T,I}$, between 120 meV and 2.5 eV, and the elastic and inelastic backward scattering cross section into the backward 2π sr, between 39 and 650 meV. Also shown are experimental values from Szmytkowski *et al.* [53], Sueoka *et al.* [55], and Sohn *et al.* [54] and theoretical values of integral cross sections from Gianturco *et al.* [27] and Bettega *et al.* [23].

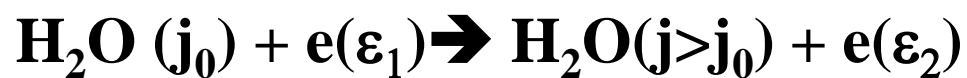
Rotational Excitation of H_2O by Cold ElectronsR. Čurík,¹ J. P. Ziesel,² N. C. Jones,³ T. A. Field,⁴ and D. Field^{3,*}¹J. Heyrovský Institute of Physical Chemistry, Dolejškova 3, 18223 Prague 8, Czech Republic

Experimental data are presented for the scattering of electrons by H_2O between 17 and 250 meV impact energy. These results are used in conjunction with a generally applicable method, based on a quantum defect theory approach to electron-polar molecule collisions, to derive the first set of data for state-to-state rotationally inelastic scattering cross sections based on experimental values.

$$B = 2 \times 10^{-3} \text{ T}$$



$$\varepsilon \sim h\nu - 15.75 \text{ eV}$$



Molecules

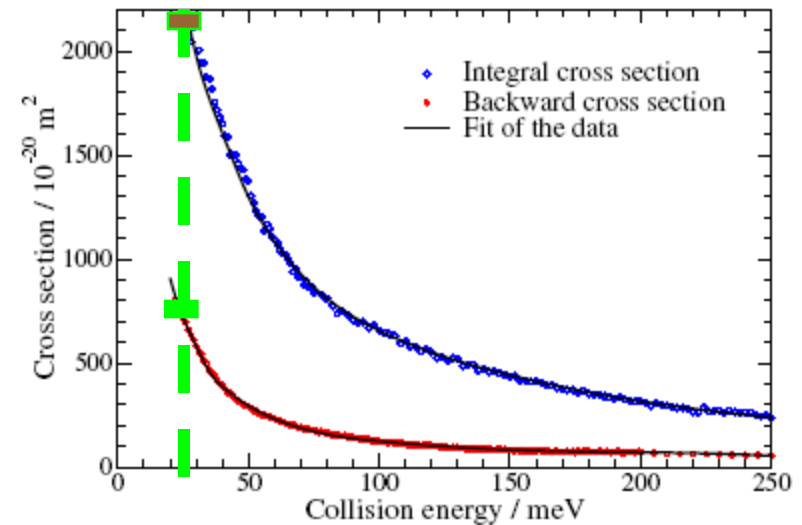
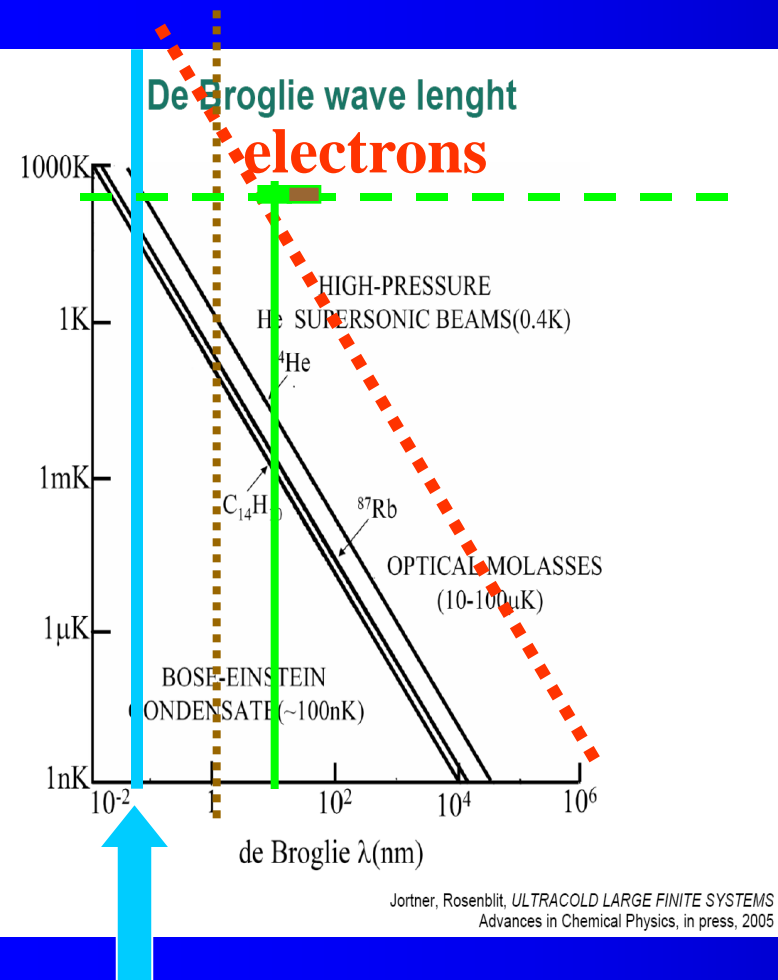


FIG. 1 (color online). Integral (upper set) and backward cross sections (lower set) for scattering of electrons by H_2O as a function of electron impact energy. Values are $\pm 5\%$. The solid lines are fits to theory (see text).

$$\lambda = \frac{h}{p} = \frac{h}{mv} \sqrt{1 - \frac{v^2}{c^2}}$$

$$\sigma \sim \pi \lambda^2 \sim 1/\varepsilon$$

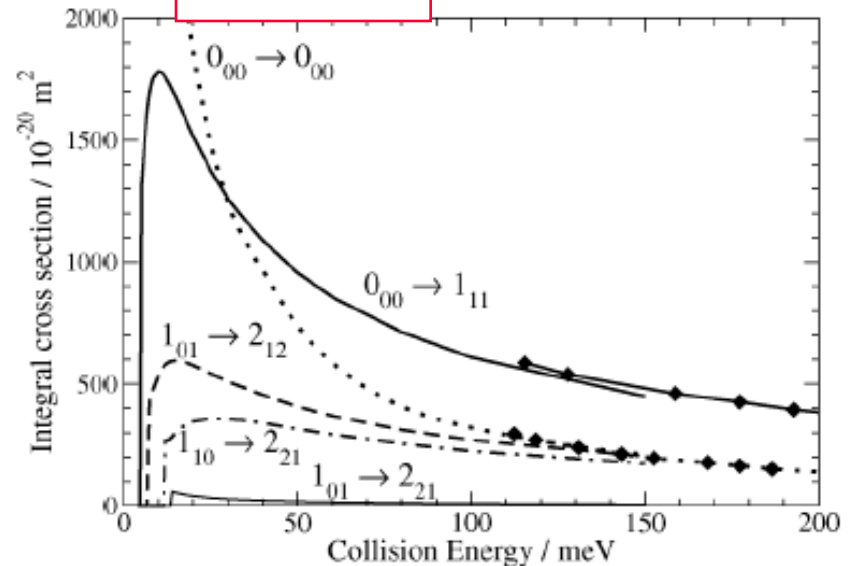
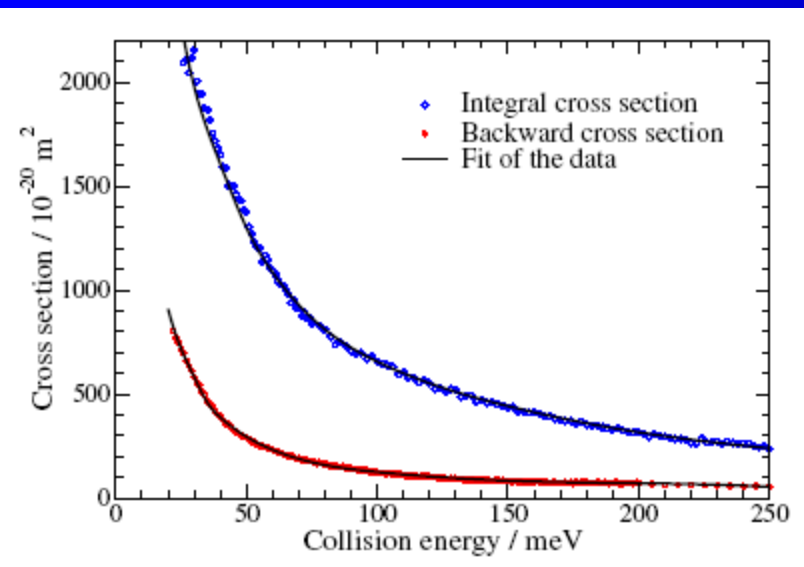


FIG. 3. Selected state-to-state integral cross sections for rotational excitation of the H_2O molecule determined from experimental data. Full curves represent results for para- H_2O and dashed for ortho- H_2O . The dotted curve represents elastic scattering for para- H_2O in its lowest rotational state. Curves with diamonds show the results of R -matrix calculations in Ref. [12].

