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## Roadmap on photonic, electronic and atomic collision physics I. Light-matter interaction

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# Roadmap on photonic, electronic and atomic collision physics I. Light-matter interaction

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## 38 39 40 41 **Abstract**

42  
43  
44 We publish three Roadmaps on photonic, electronic and atomic collision physics in order to celebrate  
45 the 60th anniversary of the ICPEAC conference. In Roadmap I, we focus on light-matter interaction. In  
46 this area, studies of ultrafast electronic and molecular dynamics have been rapidly growing, with the  
47 advent of new light sources such as attosecond lasers and X-ray free electron lasers. In parallel,  
48 experiments with established synchrotron radiation sources and femtosecond lasers using cutting- edge  
49 detection schemes are revealing new scientific insights that have never been exploited. Relevant  
50 theories are also being rapidly developed. Target samples for photon-impact experiments are expanding  
51 from atoms and small molecules to complex systems such as biomolecules, fullerene, clusters and  
52 solids. This Roadmap aims at looking back along the road, explaining the development of these fields,  
53 and looking forward, collecting contributions from twenty leading groups from the field.  
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### INTRODUCTION

#### 1: ATTOSECOND ELECTRON DYNAMICS

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- 1.2: Attosecond Physics and Laser Control of Atomic Dynamics
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- 1.4: A solid state perspective on attosecond and High-Harmonic Generation
- 1.5: Attosecond coherent control of electronic dynamics using multi-color, phase controlled, extreme ultraviolet pulses

#### 2: FEMTOSECOND MOLECULAR DYNAMICS

- 2.1: Time-resolved molecular photoelectron spectroscopy at a seeded FEL
- 2.2: Chemical reaction dynamics from time-resolved X-ray scattering of free molecules
- 2.3: Mid-IR-driven attosecond electron wave packets for molecular imaging
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Accepted manuscript

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3 **INTRODUCTION** – Kiyoshi Ueda, Tohoku  
4 University, Emma Sokell, University College Dublin,  
5 Stefan Schippers, Justus-Liebig-University Giessen,  
6 Friedrich Aumayr, TU Wien, and Hossein Sadeghpour,  
7 Harvard University  
8  
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10 To celebrate the 60th anniversary of the ICPEAC  
11 conference, we publish a series of Roadmaps on  
12 photonic, electronic and atomic collision physics. One  
13 for each of the three classes of projectile that comprise  
14 the breadth of topics encompassed by ICPEAC; I. Light-  
15 matter interaction; II. Electron and antimatter  
16 interactions; and III. Heavy particles: with zero to  
17 relativistic speeds. Each of the Roadmaps is intended to  
18 provide an overview of the present status of the field,  
19 how it was arrived at and address current and future  
20 challenges faced by those working in the broad area of  
21 research. As with all IOP Roadmaps, the three articles  
22 have been authored collaboratively by leading  
23 researchers in the areas and each aims to provide an  
24 impression of current trends in the respective field of  
25 research.  
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28 Roadmap I. Light-matter interaction is dedicated to  
29 recent advances in photonic collision physics, one of the  
30 major scientific fields covered by ICPEAC. In this area,  
31 studies of attosecond electronic dynamics and  
32 femtosecond molecular dynamics have been rapidly  
33 growing. The natural time scales of electronic dynamics  
34 in any form of matter are attoseconds, while those of  
35 molecular dynamics are femtoseconds. The advent of  
36 new light sources, such as attosecond lasers and X-ray  
37 free electron lasers (FELs) with femtosecond pulse  
38 durations, are paving the way for such ultrafast science  
39 areas. Continuously evolving developments of new light  
40 sources such as high-repetition-rate X-ray FELs and  
41 super-intense lasers are opening new pathways for new  
42 challenges. Relevant theories have also been developing  
43 rapidly, helping to design the experiments with new-  
44 generation light sources and/or to interpret the  
45 experimental results. In parallel to these new directions  
46 in photon-impact experiments with the new generation  
47 light sources, experiments with established synchrotron  
48 radiation sources and femtosecond lasers with cutting-  
49 edge detection schemes such as multi-electron  
50 coincidence detections and momentum-resolved  
51 electron-multi-ion coincidence detections (often called  
52 COLTRIMS or reaction microscope) have been  
53 continuously revealing new scientific insights that have  
54 never been fully exploited. The target samples for these  
55 photon impact experiments are also expanding from  
56 atoms and small molecules to complex systems such as  
57 biomolecules, fullerenes, clusters and solids. The  
58 present Roadmap aims at looking back along the road  
59 briefly described above, explaining such developments  
60 of the fields, and looking forward. For this purpose, we

organize the contributions from twenty of the leading  
groups into five sections: section 1 on attosecond  
electronic dynamics, section 2 on femtosecond  
molecular dynamics, section 3 on photoionization,  
section 4 on complex systems, and section 5 on new  
light sources for new challenges.

In section 1, J. Burgdörfer, C. Lemell and X-M. Tong  
discuss recent advances and give a theoretical  
perspective of attosecond physics, bridging this  
relatively new field to atomic collision physics, which is  
a core ICPEAC topic. Th. Pfeifer describes attosecond  
physics and laser control of electron dynamics in atomic  
systems, arguing that it serves as a map and compass for  
navigating future terra incognita of biomolecular and  
condensed-phase materials. F. Calegari, A. Palacios and  
F. Martin show the roadmap to attochemistry focusing  
on ultrafast biological processes, describing the current  
status of attosecond time-resolved experiments and  
relevant theory as well as the experimental and  
theoretical challenges that must be faced to bring the  
required paradigm shift to the field. P. Corkum  
discusses a solid-state perspective on attosecond  
science, addressing how to control high-harmonic  
generation (HHG) with synthesized materials and how  
to use HHG to probe the materials. G. Sansone, E.V.  
Gryzlova and A.N. Grum-Grzhimilo describe the  
achievements and future directions of attosecond  
coherent control of electronic dynamics using multi-  
colour, phase controlled, extreme ultraviolet pulses  
generated at FERMI, a FEL facility in Italy, from both  
experimental and theoretical points of view.

In section 2, M.N. Piancastelli describes how to probe  
molecular dynamics with time-resolved photoelectron  
spectroscopy, a technique established at FERMI, and  
with time-resolved electron-multi-ion coincidence  
measurements at high-repetition FELs in the future.  
P.M. Weber details how molecular movies are made,  
taking a series of snapshots of atomic motion in  
chemical reactions, using XFEL-based time-resolved X-  
ray scattering from gaseous molecules, and gives  
perspectives for the planned high-repetition rate hard X-  
ray FEL (LCLSII-HE). T. Steinle, K. Amini and J.  
Biegert describe the current status and future challenges  
for mid-infrared-laser-induced electron diffraction as a  
unique tool for imaging hydrogen atoms in a molecule  
in reaction, something which is not possible using other  
techniques at the moment. N. Berrah and E. Kukk  
describe how to study ultrafast electronic and molecular  
dynamics, specifically local damage at the atomic level,  
caused by intense X-FEL pulses, and discuss future  
challenges presented by high-repetition rate FELs. R.  
Santra discusses the importance of understanding the  
interplay between electronic and nuclear dynamics  
triggered by photoionization of a valence electron, from

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2  
3 a theoretical point of view, and considers future  
4 challenges for theory.  
5

6 In section 3, A. Müller looks back on the history of  
7 photoionization (PI) experiments on positive and  
8 negative ions, describes the latest state of the art  
9 experiment at PETRA III, a low-emittance synchrotron  
10 radiation facility in Germany, and argues for the use of  
11 ion storage rings as one of the most important directions  
12 for future PI experiments on ions. D. Doweck, R.  
13 Lucchese and W. McCurdy review the experimental and  
14 theoretical study of photoelectron angular distributions  
15 in the molecular frame (MFPAD) and discuss time-  
16 resolved MFPAD studies as probes of nuclear and  
17 electronic dynamics in molecules and of attosecond  
18 photoionization time delays. P. Bolognesi and L. Avaldi  
19 look back on the history of photo-double ionization  
20 (PDI) on the two-electron systems of He and H<sub>2</sub>, discuss  
21 PDI as a microscope for electron-electron correlations  
22 in larger systems and consider time-resolved PDI  
23 studies as probes of electron-electron correlations in the  
24 time domain. T. Jahnke, M. Schöffler and R. Dörner  
25 describe the successful history of COLTRIMS reaction  
26 microscope experiments, detail how to get time-  
27 resolved information without pump-probe techniques  
28 and discuss new challenges of employing these  
29 techniques to not only high-repetition XFELs but also  
30 high-energy “orbital angular momentum” photon beams  
31 that may become available at new generation ultralow  
32 emittance synchrotron radiation sources and open  
33 untouched areas of fundamental research. Y. Mairesse,  
34 L. Nahon and O. Smirnova describe the history and  
35 current status of the study of circular dichroism in the  
36 photoemission of chiral molecules. Time-resolved  
37 measurements with various kinds of light sources, how  
38 the dynamical chiral response study of biological  
39 molecules benefits from advances in ultrafast laser  
40 technology, sample delivery, electron-multi-ion  
41 coincidence techniques and theory are considered.  
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45 In section 4, T. Schlathölter looks back at the journey  
46 destined to gain a fundamental understanding of the  
47 molecular mechanisms underlying radiation therapy and  
48 discusses the challenges presented by ion traps and  
49 storage rings, as well as new light sources such as  
50 XFELs and attosecond lasers. E. E. B. Campbell  
51 describes the history of the study of electron emission  
52 from C<sub>60</sub> with various light sources and discusses the use  
53 of cryogenic storage rings in combination with laser  
54 excitation and photon detection, predicting that it will  
55 open a new avenue to fullerene astrochemistry, which  
56 will contribute to confronting the challenge of  
57 understanding the formation and role of fullerenes in the  
58 universe. J.M. Rost discusses the study of clusters with  
59 FELs and details the unique information that can be  
60 retrieved by time resolved experiments, such as non-

equilibrium dynamics and decoherence, and considers  
how advances of theory can inform experiments to  
address such dynamics.

In section 5, M. Meyer describes the status and  
challenges at XEF facilities in general and discusses a  
new horizon promised by super-conducting-technology-  
based high-repetition FELs with ultrashort pulse  
durations, combined with advanced detector technology  
and coincidence techniques. K.A. Tanaka describes the  
status of the Extreme Infrastructure-Nuclear Physics  
(ELI- NP), which produces ultra-intense laser fields  
with intensities up to 10<sup>22</sup>-10<sup>23</sup> W/cm<sup>2</sup> and high  
brilliance gamma-ray beams with photon energies up to  
20 MeV, and details two flagship experiments, the  
challenge of dealing with strong QED effects in laser-  
electron collisions with a 10PW laser pulse and a  
nuclear resonance fluorescence experiment for <sup>7</sup>Li  
photodisintegration related to nucleosynthesis in the  
Big-Bang.

We hope that the contributions to this Roadmap  
introduced above, which are representative of many  
more related scientific activities, bear vivid witness that  
photonic collision science has been thriving in the sixty  
years of ICPEAC history and continues to open new  
science fields with new light sources.

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# 1 ATTOSECOND ELECTRON DYNAMICS

**1.1 Attosecond Theory** – Joachim Burgdörfer and Christoph Lemell, Vienna University of Technology (Austria), Xiao-Min Tong, University of Tsukuba (Japan)

## Status

The natural time scale of electronic dynamics in atoms, molecules, and solids and of the formation and propagation of non-stationary electronic states, i.e. “wavepackets”, is attoseconds ( $1 \text{ as} = 10^{-18} \text{ s}$ ). As recognized since the early days of quantum physics, the time-dependent Coulomb field of a charged particle passing by an atom with high projectile velocity  $v_p$  acts as an extremely short half-cycle pulse of a “virtual” photon field [1][2][3] with a broad frequency distribution. When the typical width of its Fourier spectrum  $\omega \sim 1/\tau_p = v_p/a$  ( $a$ : linear dimension of the atom) matches the excitation energy of the atom,  $\Delta\epsilon \cong \hbar/\tau_p$ , the probability for excitation and ionization by the Coulomb field is largest. This celebrated Massey criterion [4] often referred to as the velocity-matching criterion has been one of the corner stones of atomic collision physics, the overarching theme of the ICPEAC conference series now celebrating its 60<sup>th</sup> anniversary (Sir Harrie Massey was one of the co-organizers of the third edition of the ICPEAC in 1963). Translated into collision times,  $\tau_p$  is in the tens of attoseconds for outer-shell electrons and can be as short as  $\sim 10$  zeptoseconds ( $1 \text{ zs} = 10^{-21} \text{ s}$ ) for deeply bound inner-shell electrons. Ultrafast dynamics and attosecond physics are, thus, the hallmark of atomic collision physics.

As major advances in ultrafast laser technology opened up novel opportunities to drive and to interrogate electronic non-equilibrium dynamics with well-controlled electromagnetic “designer” pulses, the field of attosecond dynamics dramatically expanded its scope well beyond the confines of collision physics [5]. The dream of watching electronic dynamics, charge transfer, bond breaking and making, or electron ejection in real time is becoming reality even though many major challenges still lie ahead to fully realize this dream. These experimental advances also pose major challenges to theory. The accurate description of the time-dependent electronic dynamics driven by strong light fields requires the full non-perturbative solution of the many-electron time-dependent Schrödinger equation (TDSE). On an even more fundamental level, conceptual questions have come to the fore as to which physical observables can be accessed in the time domain that add to and complement information accessible by conventional spectroscopy. Such timing information, often referred to as

attosecond metrology and chronoscopy, holds the promise of unprecedented insights into correlated electronic dynamics in atoms, molecules, and solids.

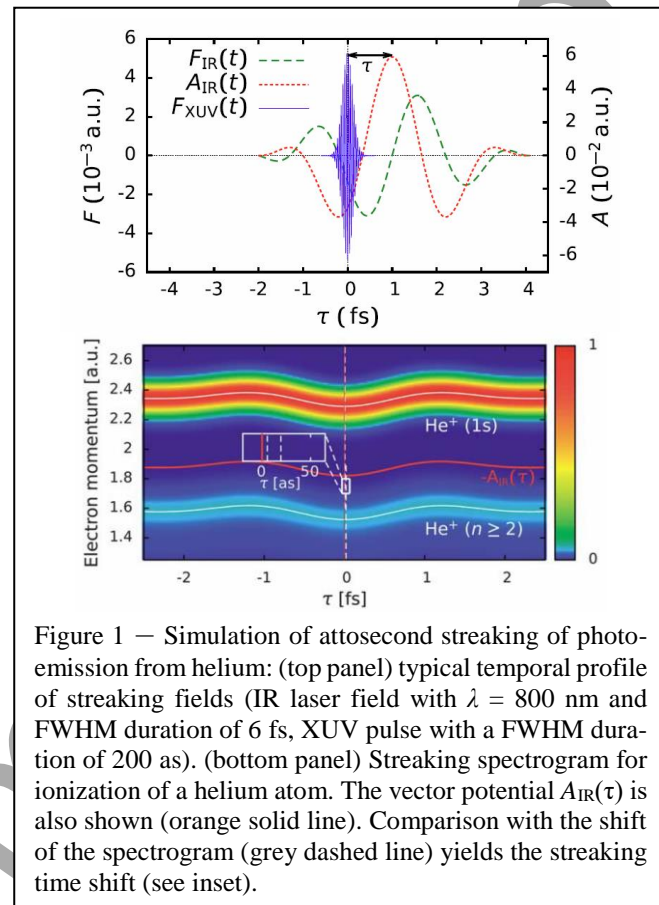


Figure 1 – Simulation of attosecond streaking of photoemission from helium: (top panel) typical temporal profile of streaking fields (IR laser field with  $\lambda = 800 \text{ nm}$  and FWHM duration of 6 fs, XUV pulse with a FWHM duration of 200 as). (bottom panel) Streaking spectrogram for ionization of a helium atom. The vector potential  $A_{\text{IR}}(\tau)$  is also shown (orange solid line). Comparison with the shift of the spectrogram (grey dashed line) yields the streaking time shift (see inset).

The availability of phase-controlled strong few-cycle IR pulses well correlated with bursts of  $\sim 100$  attosecond XUV pulse generated by the same IR pulse through the process of high-harmonic generation (HHG) has enabled the timing of the fundamental photoemission process from atoms, molecules, and solids. In a pump-probe setting the photoemission triggered by absorption by the XUV pulses can be time resolved by the concomitant IR pulse employing either the attosecond streak camera (Fig. 1; [6]), angular streaking (“attoclock”) [7], or the interferometric RABBIT technique [8].

One of the first major breakthroughs of attosecond chronoscopy was the observation of a finite time delay of photoemission relative to the arrival of the ionizing XUV pulse. The Eisenbud-Wigner-Smith (EWS) time delay originally introduced for resonant elastic scattering

$$\tau_{\text{EWS}} = -i\hbar S^+(\mathbf{E}) \frac{\partial}{\partial \mathbf{E}} S(\mathbf{E})$$

governed by the scattering matrix  $S(\mathbf{E})$  became a novel experimentally accessible microscopic observable for photoionization (Fig. 1; [9]). Delays as short as a few attoseconds could be found for photoionization of rare gas atoms. For photoemission from extended solid-state



systems, contributions originating from the primary photoabsorption, the propagation of the excited electron wavepacket to the surface, and scattering during transport could be identified.

Another class of time resolved observables is related to the ubiquitous appearance of quantum beats [10]. A sudden perturbation of the initial state creates a wavepacket composed of a coherent superposition of excited eigenstates of the system

$$|\Psi(t)\rangle = \sum_n a_n(\mathbf{0}) e^{-i\varepsilon_n t} |\varphi_n(\mathbf{0})\rangle.$$

Reduced observables such as the one-particle density  $\mathbf{n}(\vec{r}) = |\langle \vec{r} | \Psi(t) \rangle|^2$ , and, more generally, elements of the one-particle density matrix oscillate as a function of time  $t$  that has elapsed between the pump pulse and the subsequent probe pulse. In the field of atomic collision physics, quantum beats have been observed over a remarkably wide range of time scales extending from nanoseconds for fine-structure quantum beats [10] to the zeptosecond time scale for quasi-molecular x-ray emission in energetic highly-charged ion-atom collisions [11]. The recent availability of carrier-envelope phase controlled few-cycle ir pulses and attosecond xuv pulses has opened the door to direct access to the coherent excitation amplitudes  $a_n(\mathbf{0})$  and beat frequencies. Prominent examples include the beats in coherently excited neutral helium [12], the build-up of a Fano resonance in the time domain (Fig. 2), the fine-structure beats observed by attosecond transient absorption [14], and charge migration in molecules [15], i.e. oscillations of the electron or hole density across the molecules.

A third avenue towards encoding and mapping attosecond-time scale information involves interferences dubbed the “time double slit” or the “time grating”. The highly non-linear ionization response to strong IR laser pulses results in the confinement of emission to a small fraction of an optical cycle. Only near the peak field, tunneling ionization effectively takes place giving rise to narrow “time slits” of sub-femtosecond durations. Repeated emission near different peaks of the same cycle (“intracycle”) or of subsequent cycles (“intercycle”) form a double slit or a grating in the time domain giving rise to electronic wavepackets with a high degree of temporal coherence [16]. Interferences in electron emission can be viewed as Ramsey interference fringes generated by spatially separated excitation in a molecular beam [17] extended to the (sub)femtosecond scale. HHG itself, the most important building block of attosecond pulse technology, is the consequence of the temporal coherence between successive electron emission bursts every half of a period  $T = 2\pi/\omega$  of the optical cycle of the IR pulse. Such time-domain diffraction patterns enable access to a multitude of novel observables ranging from orbital tomography [18] to valence-electron wavepackets [19].

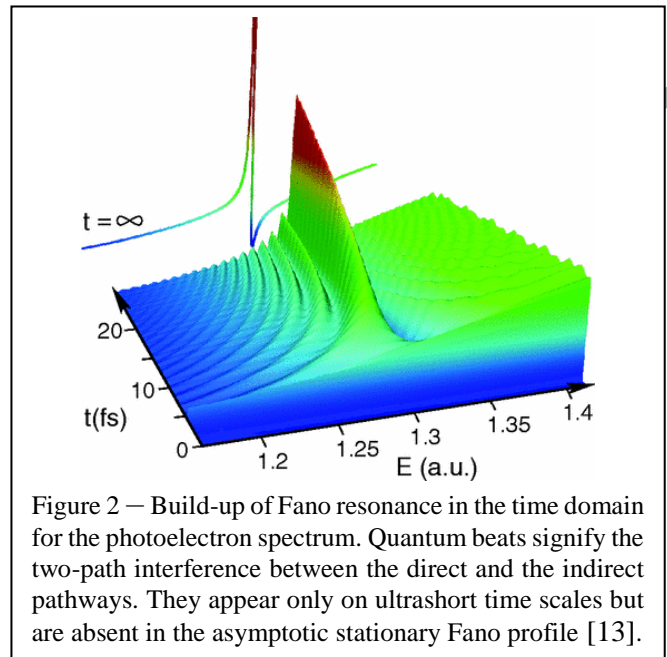


Figure 2 — Build-up of Fano resonance in the time domain for the photoelectron spectrum. Quantum beats signify the two-path interference between the direct and the indirect pathways. They appear only on ultrashort time scales but are absent in the asymptotic stationary Fano profile [13].

Current standard tools for the theoretical description of chronoscopic information include numerical solutions of the TDSE on the single-active electron (SAE) level [20], strong-field distorted-wave Born approximations closely related to the Lewenstein model [21] or the semiclassical 3-step model combining tunneling ionization, classical motion of the free electron driven in the IR field and radiative recombination upon return to the atomic core [22]. For two-electron systems, accurate numerical solutions of the TDSE have become available for attosecond XUV and moderately intense IR fields. In the perturbative limit, multi-configuration Hartree-Fock and random phase approximations have been successfully applied. For the many-electron systems such as heavy atoms, molecules, and solids time-dependent density functional theory (TDDFT) has become the workhorse of simulations of the non-linear attosecond pulse response of matter.

### Current and Future Challenges

For realizing the full potential of attosecond chronoscopic investigations considerable challenges still lie ahead. Present pump-probe settings involve mostly a combination of an ionizing attosecond XUV pulse (or pulse train) and a few-cycle IR pulse. Therefore, coherent excitation of attosecond bound-state wavepackets near the ground state in real time still remains to be realized. The recently developed optical attosecond pulses [23] might overcome this hurdle. Dynamical correlations have so far been mainly accessed in (multiple) ionization processes, most recently for correlated excitation-ionization [24][25]. With further increase in intensity, attosecond pump-attosecond probe investigations of correlated inner-shell processes as well as of multiphoton multi-electron ionization should become available in the near future. Accurately accounting for such

open-shell processes involving multiple continua remains a major challenge to theory.

The development of circularly polarized attosecond pulses provides new opportunities but also new challenges. The trajectory of the ionized electron and, thus, its recollision with the core can be precisely timed by two counter-rotating  $\omega$  and  $2\omega$  driving pulses thereby selectively generating harmonics [26]. Moreover, elliptically polarized high harmonics may allow for inner shell excitation with a net magnetic dipole moment. Probing ultrafast magnetism and relaxation, chirality of molecules, and spin dynamics in materials is just beginning. The additional rotational and spin degrees of freedom accessible pose additional demands on theory as the dynamically relevant Hilbert subspace is considerably enlarged.

Extending concepts of attosecond metrology to more complex targets, most notably to solids, is rapidly becoming one of the key areas of applications. Observing in real time the formation and decay of excitons and plasmons or atomic scale electron and spin transport promises information complementary to what is accessible by conventional spectroscopy. Extensions to the liquid state relevant for applications to biologically relevant materials are underway [27]. Close to single-cycle IR pulses [23] allow for the extreme broad-band coherent illumination and the sensitive detection of vibrational and rotational resonances of bio-molecules [28]. The ubiquitous presence of dephasing and decoherence in the condensed matter environment destroys the strict Fourier reciprocity of spectral and time information encoded in the wavepacket and emphasizes the importance of direct time domain studies of transport and relaxation. Uncovering the microscopic processes underlying line broadening and dephasing in open quantum systems may, quite likely, become one of the key contributions of attosecond physics to complex systems. Attosecond physics is also expected to become relevant for devices: gas-phase streaking set-ups may soon be replaced by sub-cycle resolved measurements of currents in semiconductor or dielectric structures. HHG in bulk solids and two-dimensional materials has already been successfully demonstrated. Irradiation of nano-scale structures such as nanotips opens up the opportunity to combine attosecond time with nanometer scale spatial resolution [29]. Attosecond transient absorption and reflectance give access to the bulk and surface response of valence and conduction band electrons on the (sub)femtosecond scale [30]. A quantitative description requires, however, a multi-scale theory entailing a self-consistent coupling of the non-linear electronic response to the Maxwell field propagating inside the material and, on longer time scales, to lattice degrees of freedom.

One important area of research awaiting progress as higher intensity attosecond sources become accessible

is time-resolved multi-particle break-up. For example, the paradigm case of sequential or non-sequential two-photon double ionization of helium has still not yet been experimentally observed in real time even so time sequencing is at the very heart of this process. Multi-electron emission from larger molecules and solids may shed new light on electron correlations, charge migration, and ultrafast screening of Coulomb holes.

### Advances in Science and Technology to Meet Challenges

Recent theoretical and computational advances are expected to meet some of the challenges posed by the progress in experimental attosecond physics. The time-dependent multi-particle dynamics calls for a theoretical description beyond the highly successful SAE models. While TDDFT allows the treatment of the many-electron system of extended targets such as  $C_{60}$  or dielectrics on a mean-field level, correlated wavefunction methods such as multi-configuration time-dependent Hartree-Fock have opened up novel opportunities to treat many-electron atoms in ultrashort pulses numerically with high accuracy and in full dimensionality. Extending established methods of quantum chemistry to the time domain such as time-dependent complete active space self-consistent field, time-dependent coupled cluster, and the non-equilibrium Greens function methods and method will expand the arsenal of theoretical tools [31]. Similarly, the time-domain extension of quantum Monte Carlo methods may open new possibilities for extended systems. In between accurate correlated  $N$ -electron wavefunction methods that scale rather unfavorably with system size and mean-field methods such as TDDFT which allow for the efficient treatment of large systems but fail to explicitly describe dynamical correlations, the time-dependent two-particle reduced density matrix theory [32] may fill the gap if the challenges of  $N$ -representability can be overcome.

### Concluding Remarks

Attosecond physics is at the core of both atomic collision physics and the development and application of new light sources based on high-harmonic generation. Chronoscopy is increasingly complementing conventional spectroscopy providing information not easily accessible in the frequency domain. Future breakthroughs will likely address relaxation and decoherent dynamics in open and complex quantum systems. Attosecond physics holds the promise to not only observe but to actively control, steer, and manipulate electronic motion with the potential of device applications.

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Accepted manuscript

## 1.2 Attosecond Physics and Laser Control of Atomic Dynamics – Thomas Pfeifer, Max-Planck Institute for Nuclear Physics

### Status

Atomic systems have shaped the development of physics in ways second to none. A key to the discovery of quantum mechanics was the discrete nature of atomic spectral lines. Today, atoms are employed to test quantum electrodynamics, search for the variation of fundamental constants, and to place more and more stringent limits on the existence of new physics beyond the standard model of particle physics.

It thus comes with no surprise that atoms also play a pivotal role in attosecond physics. The reason is obvious: Atoms, at least in their ground electronic state, and near closed electronic shells, are excellently understood by theory (section 1.1). Any departure from their known properties — e.g. in new environments such as in strong fields on attosecond time scales — is thus immediately obvious in experimental observables, such as their spectral features.

Atoms are thus fundamental sensors of dynamical processes on all time scales from attoseconds up to the age of the universe.

High-harmonic generation (HHG), today considered the birth of attosecond physics, was discovered in the 1980s [1, 2], by focusing intense laser pulses into atomic gases while observing surprising plateaus in the extreme-ultraviolet(XUV) spectra.

While providing the key ingredients for attosecond physics — i.e. attosecond pulses coming in trains [3] or isolated bursts [4, 5] — the HHG process itself also gave rise to a whole new field nowadays termed recollision physics [6, 7].

From the first measurement of attosecond pulses in 2001, progress is explosive, leading to the current state of the art with exceptional achievements: Direct measurements of visible light wave electric fields by streaking[8], attosecond photoelectron interferometry [9], clocking Auger-[10] and photoionization[11, 12], time-resolved measurements of bound electron motion in rare-gas valence holes[13] as well as doubly-excited states[14]. Almost any of these achievements has now been transferred to fruitful applications in the fields of physical chemistry and condensed-matter physics.

The growth of this field is fueled by technological developments, such as increasingly controlled driver pulses for HHG, including their carrier-envelope phase, but also on the detection side by employing photoelectron/ion imaging, as well as time-resolved and transient absorption spectroscopy (Fig. 1).

Meanwhile, with yet another technological revolution, namely free-electron lasers (FEL) a new horizon opened: Attosecond Physics in intense XUV/x-ray fields. For approaching this new physics regime of bright light-matter interaction — where a single atom absorbs multiple high-energy photons on short time scales within a single FEL pulse — atomic systems were crucial [15, 16].

The table here summarizes a few key figures of relevance for attosecond physics nowadays and expected in the near future:

Properties\Source	HHG	FEL
Shortest duration	~10 as	~100 as
Pulse energy	~1 nJ -1 $\mu$ J	~0.01-10 mJ
Intensity (1 $\mu$ m focus)	$\sim 10^{16}$ - $10^{19}$ W/cm <sup>2</sup>	$\sim 10^{19}$ - $10^{22}$ W/cm <sup>2</sup>
Photon energy	~0.01-1 keV	~0.1-10 keV

The complementarity of the two source technologies is clearly visible.

It should be noted that x-ray FEL pulses in the SASE regime intrinsically exhibit an attosecond-pulsed temporal structure. While this structure is statistical and irreproducible between shots, it can be used for enhanced temporal resolution [17]. While stochastically isolated sub-fs FEL bursts have already been measured, soon seeded and thus fully controlled attosecond pulses will be accessible at FELs. As discussed in section 1.5, glimpses of the new physics possible in this regime have also meanwhile become available by observing directional emission of electrons from atoms in attosecond controlled FEL waveforms [18].

### Current and Future Challenges

Due to the fundamental *structural simplicity* of atoms, they ideally allow to study the emergence of *dynamical complexity*. When can an atom be described as a quantum system consisting of just two or few levels? When does the coupling to the continuum become important? Can the response of atoms be described by effectively one active electron, or do we need to consider two- or multi-electron effects? And, in the latter case, is it enough to consider a mean-field interaction of a subset of electrons, or does collective, or even correlated, motion of the electrons become important?

For fundamental light-matter interaction, even the most simple case of single-electron ionization remains an active field of research. What is the time delay of photoemission caused by the atomic structure, and the delay due to the combined action of laser and Coulomb fields? Can we define, in a quasi-classical picture, how long it takes an electron to tunnel through a barrier?



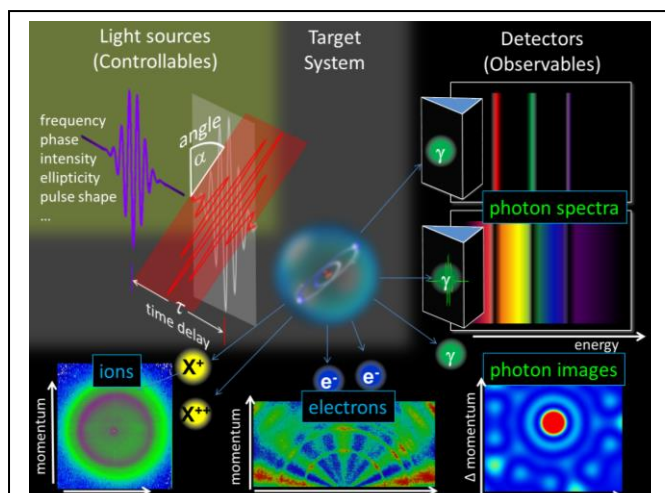


Fig. 1 – Fully-differential measurement of atomic quantum dynamics driven by multi-parameter controlled light: Testing our fundamental understanding of time-domain physics in atoms also sharpens the tools for multi-particle attosecond physics and chemistry.

What we learn from these questions until now is that electron-electron interaction, giving rise to electron correlation, plays a crucial role and introduces complexity already for the case of two-electron atoms. A challenge in the field is thus to measure, understand and to control correlation dynamics among electrons.

Why is this important? One example: The helium atom's ground-state electronic wave function is a perfectly spin-entangled (and thus quantum correlated) state:

$$|\Psi\rangle \propto |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$$

In fact, any atom's ground state provides a fully entangled cluster state. If we can control electron correlation at will, say, with attosecond shaped-pulse sequences, this would allow us to set up a self-learning quantum-computer architecture on an atomic level, with clock speeds near and exceeding PHz frequencies. A far-field vision, perhaps, but so were attosecond light pulses still in 1980, 20 years can mean a lot in science!

On the way there, we need to understand electron correlation dynamics and its manipulation. Two-photon double ionization [19] remains a benchmark experiment to be conducted in Helium and heavier species, and is still only a first step towards multi-photon and strong-field multi-electron ionization and excitation experiments, culminating in the controlled synthesis of atomic wavefunctions, including transiently populated highly-excited and extremely short-lived superpositions of neutral hollow-atom states.

The understanding and full mastery of electron correlation dynamics on an atomic level is also the fundamental key to electron-correlation-driven charge migration [20], which is a major goal in molecular

attosecond science (section 1.3). Full controllability would allow attosecond chemistry to synthesize any stoichiometrically allowed molecular structure or even a macroscopic object from a bunch of atoms in the ultimate realization of a molecular 3D printer.

Required are experiments in intense high-frequency fields, as only these are capable of specifically moving electrons with unit efficiency among inner shells. Again, the intense fields of FELs will play a pivotal role in the near future of attosecond physics. What is the dynamics of electrons driven in strong FEL fields, and what fundamental physical mechanisms can we extract to expand from two-electron systems (Fig. 2) to multi-electron atoms, and from there to clusters (section 4.3) and larger molecules (sections 2.4, 4.2)?

The near-future technological frontier in attosecond physics is the attosecond-pump-attosecond-probe measurement of electron motion. A fertile testing ground of this technique could be attosecond time-resolved photoelectron emission asymmetry in atoms, Auger spectroscopy, as well as attosecond transient-absorption after the coherent excitation of multi-eV-spaced electronic states in atoms.

Applying and understanding these methods in isolated atoms will calibrate these element-specific techniques, in the best case implemented as two-color as-pump-as-probe. This would allow physical chemists and

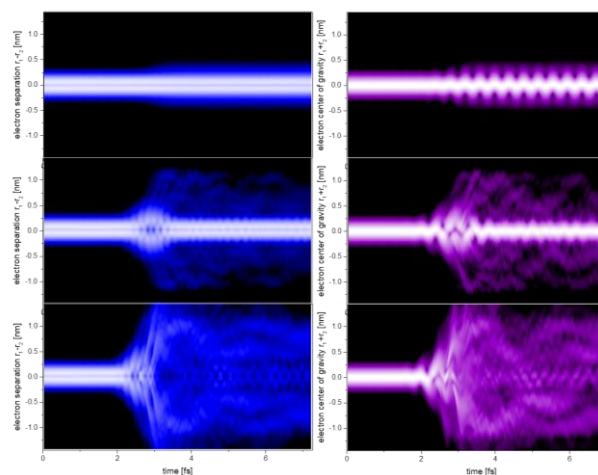


Fig. 2 – Emergence of dynamical complexity in a two-electron atomic model system, resonantly driven by increasingly intense high-frequency (FEL) pulses (0.8 fs duration): Top to bottom: 0.01, 0.1, 0.3 a.u. peak field strength. The squared modulus of the two-electron wavefunction as a function of electron-electron separation  $r_1-r_2$  (left) as well as proportional to their center of mass  $r_1+r_2$  (right) is shown after integrating over the other coordinate, respectively. With increasing intensity, a complex time-dependent evolution is observed, illustrating the origin of correlated charge migration dynamics on an atomic level.

condensed-matter physicists to specifically address and temporally resolve time-dependent processes in these complex systems (section 4) on the electronic time scale [21].

### Advances in Science and Technology to Meet Challenges

To approach these visions, a continuation of the (r)evolution of light-sources as well as detectors is required.

On the light-source side, intense optical lasers and high-harmonic generation have pioneered the field and have meanwhile matured, allowing the control over pulse shape in space and time, as well as the polarization state of light down to the attosecond domain. The next technological frontiers are more intense drivers for HHG (to enable as-pump—as-probe) and higher repetition rates (for multi-parameter controlled coincidence experiments, see Fig. 1). In recent years, FELs have entered the scene, achieving more and more controllability over pulse properties and now come with controlled attosecond durations (section 5.1). While optical laser and harmonic sources still stand out in terms of their shortest pulse durations, coherence and synchronization properties, FEL pulses routinely allow XUV/x-ray-only pump-probe experiments due to their large photon numbers in excess of  $10^{13}$  per pulse, already supplying intensities on the order of  $10^{20}$  W/cm<sup>2</sup> to explore new regimes of nonlinear optics [22].

Detection systems for atomic quantum dynamics basically come in two major categories: Electron/ion detection, and optical spectroscopy. The momenta of several electrons and the remaining ion can be measured in coincidence with Reaction Microscopes/COLTRIMS [23] (section 3.4). Optical spectroscopy techniques observe absorption, also fluorescence, of targets upon transmission of the light through the target.

These two methods complement each other in a fundamental way: Electron detection allows to observe the final (continuum) state of the samples, while photon detection can reveal the initial (bound) excitation step in the sample and how it is driven by intense light fields.

By making use of both existing detection techniques, and combining them with a multi-parameter control approach for the interacting light pulses, big data sets emerge that capture a major part of the full dynamical information. For example, the impinging light pulses can be varied in frequency, intensity, phase, polarization, and come in sequences with several time delays (Fig. 1). Two-dimensional spectroscopy [24] was just the beginning. The mechanistic dynamical information on quantum-state coherences (section 2.5) and populations extracted from such multi-dimensional measurements on atoms in strong fields will serve as

map and compass for navigating the future *terra incognita* of (bio-)molecular (sections 1.3, 2-4) and condensed-phase materials (section 1.4), and for science and technology applications that are yet unconsidered.

### Concluding Remarks

Atomic Physics is the beating heart of quantum dynamics in general, and attosecond science in particular. It provides both a mining and a testing ground for new physics mechanisms, time-domain metrology and control methods. It will play a pivotal role for the future of attosecond physics as the field is pushed forward by the joint thrust of lab-based (HHG) and large-scale (FEL) light sources, and multi-dimensional detection techniques, allowing to capture the full set of dynamical information. It will enable precision tests of our understanding of laser-driven quantum dynamics, and give rise to control techniques and applications which are still beyond our wildest dreams. Molecular-scale petahertz-clocked computing could be one of them, as well as laser-directed chemistry on the atomic level by light-field control of multi-electron correlation dynamics in the core and valence shells of any atom.

**Acknowledgments** – Attosecond physics develops rapidly due the cooperative work of many outstanding junior and senior researchers and their groups, which all deserve recognition. The few references allowed in this section can only reflect a small and non-representative subset of the massive body of researchers and their work in this field.

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### 1.3 A roadmap to attochemistry

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#### Status

Attosecond laser technology has demonstrated its suitability to trace and manipulate electron dynamics in atoms and small molecules. However, application of this technology to more complex molecules, as those involved in biological processes (photosynthesis, radiation damage, etc), is still in its infancy. A common approach to visualize this electron dynamics with few-femtosecond and sub-femtosecond time resolution is attosecond pump-probe spectroscopy [1, 2]. In this method, an attosecond extreme ultraviolet (XUV) pulse is used to excite/ionize the molecule, followed by a time delayed probe pulse to trace the dynamics generated in the first step. Typical probe pulses have frequencies ranging from the near infrared (NIR) to the X-Ray domain, with durations going from few tens of femtoseconds to few hundreds of attoseconds. These experimental capabilities are thus ready to explore the possibility to control chemical reactions by acting at the very early stages of the process, e.g., by guiding electronic motion in such a way that bonds break and form following a desired path.

In pursuing this goal, recent attosecond XUV-pump/NIR-probe gas phase experiments (set up shown in Fig. 1) performed in aromatic amino acids [3] have captured the occurrence of ultrafast charge dynamics after ionization. In these experiments, the molecule is ionized by an attosecond pulse, creating a delocalized electronic hole that moves all over the molecule on a sub-femtosecond time scale. The interplay between the electronic and nuclear motions dictates the evolution of the hole, which can lead to a final localization of the charge on one specific molecular site or specific molecular fragment. One can thus expect that, depending on the initially created hole, fragmentation will occur leading to different chemical species [4,5], hence to distinct reactivity. The chemical selectivity that could be attained by using attosecond technology remains yet to be demonstrated.

The experimental approach used to track charge migration in amino acids is based on the measurement of the photo-fragmentation yield as a function of the XUV pump NIR probe delay [3]. As shown in Fig. 1, a clean plume of intact and neutral biomolecules is produced by evaporation in the interaction region located in the middle of a mass time-of-flight (TOF) spectrometer. A sequence of isolated XUV attosecond pulses and few-femtosecond NIR pulses is used to

initiate and probe the charge migration process. The doubly charged immonium ion (resulting from the loss of the carboxyl group) has been found to show clear oscillatory dynamics, which has been assigned to charge flow between different functional groups of the molecules.

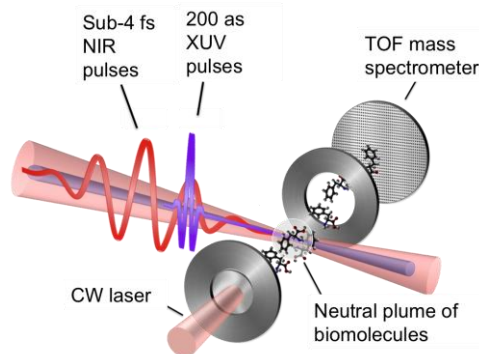


Figure 1. Scheme for the experimental detection of charge migration in amino acids.

An alternative to pump-probe methods is High-order Harmonic Spectroscopy (HHS), which has been demonstrated to be an effective tool to track electron dynamics in linear molecules, as e.g. iodoacetylene [6].

Reliable theoretical tools are mandatory to interpret attosecond spectroscopy experiments. However, a full dimensional quantum mechanical treatment of molecular ionization dynamics, accounting for all electronic and nuclear degrees of freedom, is still out of reach by using current computational approaches, except for the hydrogen molecule. Recent theoretical approaches in large molecules have incorporated the effect of nuclear dynamics by using semiclassical treatments in which the motion of electrons is described quantum mechanically and that of the nuclei by using a classical description. Application of these methods to amino acids indicate that electronic coherences can survive for at least 10 fs (see Fig. 2 and [7]), thus being experimentally accessible. In other molecules, however, electronic coherences have been found to survive 1 or 2 fs [8]. The general consensus is that lifetimes of electronic coherences are strongly dependent on the species under study and the pulse characteristics. Therefore, accounting for nuclear motion seems to be mandatory to correctly describe the effect of electronic coherences on the fragmentation that follows the sudden excitation or ionization of the molecule, even qualitatively.

Besides the above-mentioned limitations, there is the additional complication of dealing with the electronic continuum of the molecule, which cannot be represented using standard quantum chemistry methods. For this reason, most theoretical approaches for large molecules have ignored the ionization step and have only described

the cation dynamics that arises from arbitrarily chosen initial wave packets. On top of that, XUV and X-ray photons can also produce inner-shell vacancies that can subsequently be refilled by outer-shell electrons whose excess energy is then used to ionize other outer-shell electrons. This process, known as Auger decay, is mostly due to electron correlation. While this limitation can be somehow circumvented, e.g., by choosing pumping frequencies such that Auger decay barely contributes, it should be considered in future work, as production of Auger electrons has been recognized as a very important source of damage in biomolecules.

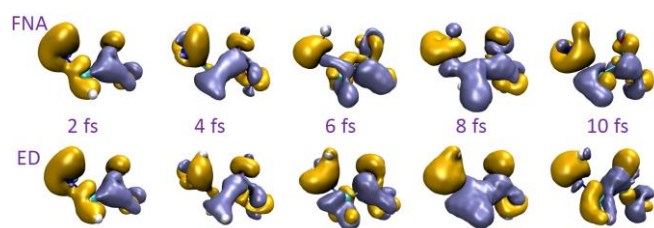


Figure 2. Evolution of the hole density in glycine after interaction with an attosecond pulse (bandwidth  $\sim 18\text{-}35$  eV). FNA: fixed nuclei approximation. ED: Ehrenfest dynamics obtained within a TD-DFT approach.

### Current and Future Challenges

One of the most ambitious challenges of attosecond science is to design control strategies that can efficiently be used to induce ultrafast electron dynamics in molecules, in gas-phase or in solution, so that one can selectively break and form chemical bonds, thus leading to new chemical reactions.

On the experimental side, there are still a lot of challenges to extend attosecond technology to complex molecules. First, complex molecules are produced with relatively low densities and the typical photon flux of the attosecond sources is at the limit to provide a reasonable signal-to-noise ratio. Second, direct imaging methods for the electron dynamics still need to be developed and third, control over the electronic excitation is required. This last goal could only be reached by improving the current capability of tailoring attosecond pulses.

On the theoretical side, many efforts are currently made to combine advanced quantum chemistry methods, as multi-reference configuration interaction methods especially designed to accurately treat electron correlation, with scattering theory formalisms appropriate to deal with the ionization continuum [2]. In addition to this, methods going beyond standard classical descriptions of nuclear dynamics, such as Ehrenfest MD or Trajectory Surface Hopping, should be developed. Fully quantum mechanical approaches, which are currently limited to small molecules, may

provide further insights on the coupled electron-nuclear dynamics and, more importantly, on establishing if quantum interferences and coherences may indeed have an impact in reactions that are relevant in chemistry or biology and occur at a much longer time scale. The latter also implies going beyond gas-phase chemistry and face the additional challenge of considering the effect of the environment [9].

### Advances in Science and Technology to Meet Challenges

Recent and future advances in attosecond technology will allow most of the experimental challenges listed in the previous section to be met. In particular, the tremendous progress in the development of ultrafast lasers (OPCPAs and fiber lasers) - already delivering hundreds of watts - allows for driving the attosecond pulse generation process with higher peak intensities and higher repetition rates. At the same time the pulse duration of the free-electron-laser (FEL) sources is approaching the attosecond limit [10]. With the photon flux provided by the FELs, it will be soon possible to perform diffraction imaging of complex molecules with attosecond time resolution, thus providing an effective method to directly imaging and control the electron dynamics.

To meet the above theoretical challenges and establish the mechanisms governing XUV irradiation of medium-size and large molecules one has to develop tools that allow one to describe (i) ionization, including electron correlation terms that determine processes such as Auger decay, and (ii) the ensuing coupled electron and nuclear dynamics. Also, as molecules responsible for most biological functions are either in solution or in a hydrophobic environment, not in the gas phase as in most attosecond experiments performed to date, theory should go a step further and consider the effect of such environment.

### Concluding Remarks

Is it possible to (efficiently) achieve chemical selectivity by using coherent ultrashort light sources? How electronic coherences can affect fragmentation after molecular excitation and/or ionization? Attosecond pump-probe and HH spectroscopies could help in answering these questions by extending the existing experimental methods -which have proven to be successful to attain sub-femtosecond electronic control in atoms and small molecules- in combination with new theoretical approaches able to describe the coupled motion of electrons and nuclei that dictate the chemistry of these systems. If successfully accomplished, this would inevitably pave the way to attochemistry.

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## 1.4 A Solid State Perspective on Attosecond and High-Harmonic Generation – Paul

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*High-harmonic radiation from solids has potential applications as a source of XUV radiation and as a means to probe solids.*

**Status:** In general, harmonic radiation is created by the nonlinear interaction of light and matter. When light is a perturbation, harmonics arise as the bound electrons explore their near surroundings, thereby formally making virtual transitions to all excited bound states. This is the condition for conventional harmonic generation.

In the other limit, when light is very intense, near-surface electrons are pulled free of the material, enabling them to move over wavelength-scale distances in response to the light's electric field. This non-sinusoidal electron motion can produce very high harmonics, especially when the light intensity is relativistic.

At intermediate intensity, if the material is irradiated by below-bandgap light, the bound electrons can make a valence-to-conduction-band transition, again creating an oscillating current that will yield harmonics of the fundamental beam.

However, inside a material, these internal electrons are confined to the bands of the solid. Electron motion still creates a time-dependent current that produces the emission, but the harmonic-spectrum results from Bloch Oscillations which are, of course, a signature of the periodic structure of the material. In such cases, the harmonic spectrum encodes the potential structure that the electron feels as it oscillates in the material.

That said, we know from studies of atomic and molecular gases that there is a second, very important mechanism for this intermediate field-driven electron to generate high harmonics. (One might think of this mechanism as a hybrid between the extremes.) Here, the newly ionized electron begins its oscillation *in the field*, but the oscillation brings the electron back to its origin, where it re-collides with its parent, recombining and thereby releasing its kinetic energy plus the ionization potential as XUV radiation [1]. This source, very important in gases, should also be present in solids where an electron and its correlated hole plays the role of an electron and ion in gases as illustrated in Fig. 1.

In fact, we now know that “re-collision” is the source of high-harmonic radiation that was discovered in ZnO by Ghimire *et al* in 2011 and analyzed in the theoretical papers of Vampa [3]. Re-collision, (identified by the time-dependent frequency of emission known as the “atto-chirp”) plays a dominant role in many, but not all, experiments [4] where high harmonics

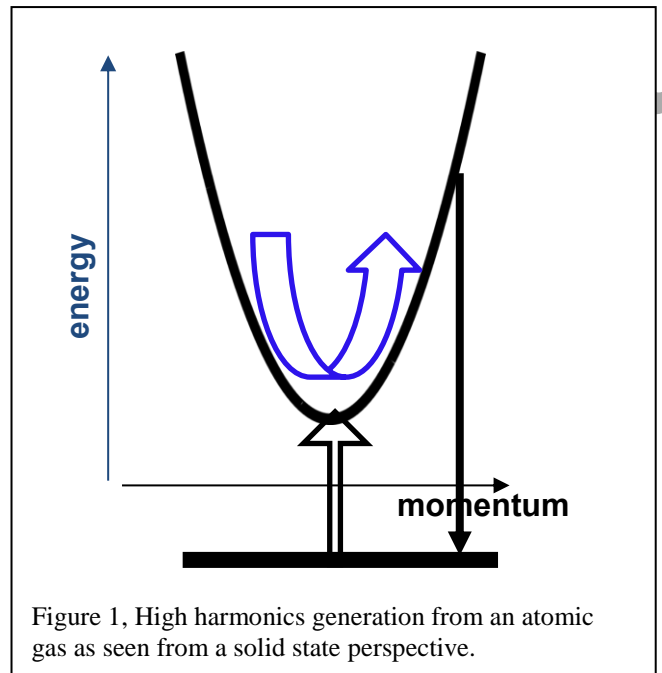


Figure 1, High harmonics generation from an atomic gas as seen from a solid state perspective.

are produced from intense mid-infrared radiation that excites transparent solids.

While I have described four interrelated mechanisms for generating harmonics of a fundamental beam, I believe that it will be possible to create *hybrids* between them – between surface harmonics and bulk harmonic radiation, for example -- but as yet there are no reported observation.

### Can high-harmonic generation become an important source of VUV radiation?

Solids have the advantage of simplicity and flexibility over gases, and the same material can be used multiple times. In addition, attosecond duration pulses can be generated. However, so far we cannot reach nearly as short a wavelength without damaging the solid as we can when we use gases; furthermore, it is not clear if the efficiency of harmonic emission from solids can approach that of gases. That acknowledged, I still predict that attosecond (or high-harmonic) pulses will be commercially generated from solid media at least for photon energy greater than 30 eV.

In gases, high-harmonic photon energy can greatly exceed ionization potential, which implies that there is a maximum medium length beyond which the generated XUV radiation is absorbed by re-ionization.

Equally, in solids, the highest harmonics can have photon energies that are far above the band gap of the material. However, in solids, the density is much higher than in gases and therefore, the absorption length is very short. In this case, phase-matching equally favours forward- or backward-propagating radiation. Therefore, it is often advantageous to irradiate from the front surface, because high-intensity radiation does not pass through the sample before it creates harmonics.

### Solid-state physics has spawned important technical developments that we can exploit to control harmonic radiation.

For example, technology allows us to dope materials, and doping will modify the amplitude and phase of harmonic emission [5] — just as we control the amplitude and phase of high-harmonic emission with mixed gases. Doping can be performed with high-resolution ion injection [5], with the solid serving as a fixed template. *This unique characteristic of solids may be very important, because it allows us to use the material to create light that naturally focusses.* Solids, used as templates on which to write, may free us from the constraints of conventional XUV optics.

But solids also give us an ability to shape materials [5]. Quantum wells, quantum wires, or structured surfaces will allow us to control the electron trajectory, and therefore, to control the efficiency of harmonic conversion [6]. It is even possible to add nano-plasmonic antennas to a solid [7] to significantly decrease the required illuminating power, thereby bringing high-harmonic generation from high bandgap materials within reach of mode-locked oscillators, while also linking high-harmonic generation to modern research on meta-materials. These nano-scale resonant structures will complement build-up cavities for creating XUV frequency combs.

### Can high-harmonic generation become an important probe of materials?

Re-collision harmonics report the difference between the energies of the two (or more) bands that are involved in high-harmonic emission. In fact, we already know that the band structure of solids can be determined from the two-dimensional (two-color) high-harmonic map [8]. These maps are similar to those used to establish the mechanism of harmonic-generation in solids [3]. Reader might be concerned with disentangling the many bands of a solid, but the requirement of phase-matching appears to select only two (or a very few) bands.

Of course, the band structure of most materials is well known from photo-electron spectroscopy, but photo-electron spectroscopy is not accessible in high-pressure diamond anvil cells, nor can it be used to follow ultrafast dynamics during laser material processing, or materials in hostile environments. It would be a major accomplishment for science if we could time-resolve phase changes driven by lattice re-arrangements or by quantum correlations.

Band reconstruction, was inspired by orbital tomography of gas phase molecules. It would be very important if we could transfer the concept of

tomography directly from gases to solids. Such a transfer now appears to be possible [9].

While the penetration depth of XUV light in materials is small, it is many monolayers deep. However, at least in carbon, a single atomic layer of material is sufficient to produce measurable high-harmonic radiation [10]. (We know this because high harmonics can be produced from graphene and other 2-D materials [10].) If a single layer of material is sufficient, then time-dependent surface chemistry can be probed. I predict that high-harmonic spectroscopy

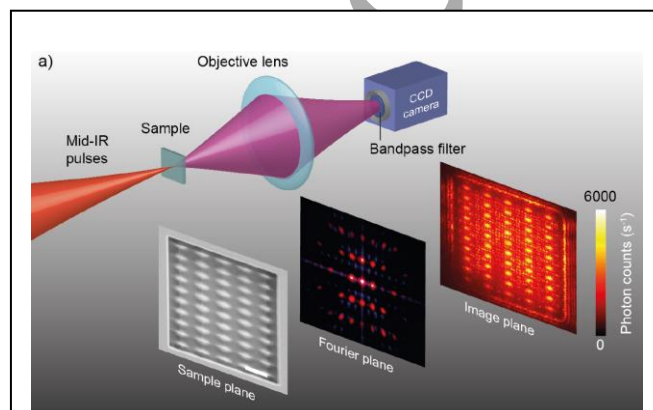


Figure 2. Schematic of an experiment where ZnO with a surface of nano-pyramids is irradiated with 2 micron light from behind. We see a photograph of the diffracted third and fifth harmonic radiation and an image of the emitting surface. (taken from Ref. 5)

will open a new diagnostic for surface chemistry.

Biological material is often highly ordered, and this order will influence the emission spectrum – as we discussed above. Thus, if thin slices of biological material are used as a substrate, then the structure and composition will be indelibly written on the harmonic spectrum. Furthermore, even amorphous and polycrystalline materials are already used to produce high harmonics [4].

Extreme nonlinear optics offers a new chapter in which, I predict, the short wavelength of the harmonic emission will be harnessed for high spatial resolution and the resulting image will accompany 2 or 3 photon imaging or CARS microscopy.

In closing, we consider electronic circuits. Measurements show that the harmonic spectrum is sensitive to small fields – fields that are on the scale of those used for switching in conventional electronic circuits. It should be possible to obtain 2-D images of an operating circuit using real-space optics for measurement up to its limits, and then to use diffractive methods such as ptychography up to the high-harmonic cut-off. It is not immediately clear what the cut-off will be. When MgO is irradiated from the front surface, in my lab we see a cut-off of 45 eV.

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### 1.5 Attosecond coherent control of electronic dynamics using multi-color, phase controlled, extreme ultraviolet pulses –

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#### Status

Coherent control in the micro world is based on the phenomenon of interference of amplitudes of two or several paths from initial state of a quantum system to its final state. Changing phases between the amplitudes controls the final result of the process [1]. Usually coherent control is performed by an optical laser, which is easy to manipulate. Studies of coherent control in the optical domain, both theoretical and experimental, were initiated about 30 years ago. Creation and manipulation of phase synchronized pulses of vacuum ultraviolet and X-ray radiation is a highly complex task because of the lack of the necessary optical elements. Only recently it was first implemented at the seeded Free Electron Laser (FEL) FERMI [2]. Coherent control was realized for the photoelectron angular distributions (PAD) from neon atoms. The two ionization paths were produced by the fundamental frequency (two-photon “process  $\omega$ ”) and by its second harmonic (one-photon “process  $2\omega$ ”). The control of the phase between the harmonics, and, consequently, between  $\omega$  and  $2\omega$  ionization amplitudes, was carried out without optical elements by adjustable time delay with a temporal resolution of 3 as of electrons bunch generating the second harmonic. For two parallel linearly polarized  $\omega$  and  $2\omega$  pulses, the asymmetry of the PAD with respect to a plane perpendicular to the polarization direction was controlled by the relative phase. For two circular polarizations, the axial symmetry with respect to the collinear radiation beams would be violated, as predicted by theory (Fig. 1).

The realization of coherent control in the extreme ultraviolet (EUV) and X-ray range represents a fundamental milestone in the development of coherent spectroscopy in the EUV and soft X-ray range, which is one of the largely unexplored areas of research in ultrafast phenomena. Advances in this field will be boosted by the implementation in the X-ray regime of spectroscopy techniques already established in the optical domain such as Raman and photon echo spectroscopy. Their demonstration will give new insight into ultrafast phenomena including relaxation processes such as Auger decay and migration of electrons in chemical reactions. One of the main advantage of the soft and X-ray spectral range is the chemical selectivity in the excitation and probing of ultrafast electronic dynamics, a characteristics which is absent for interaction with photons in the optical range. Thus, soft

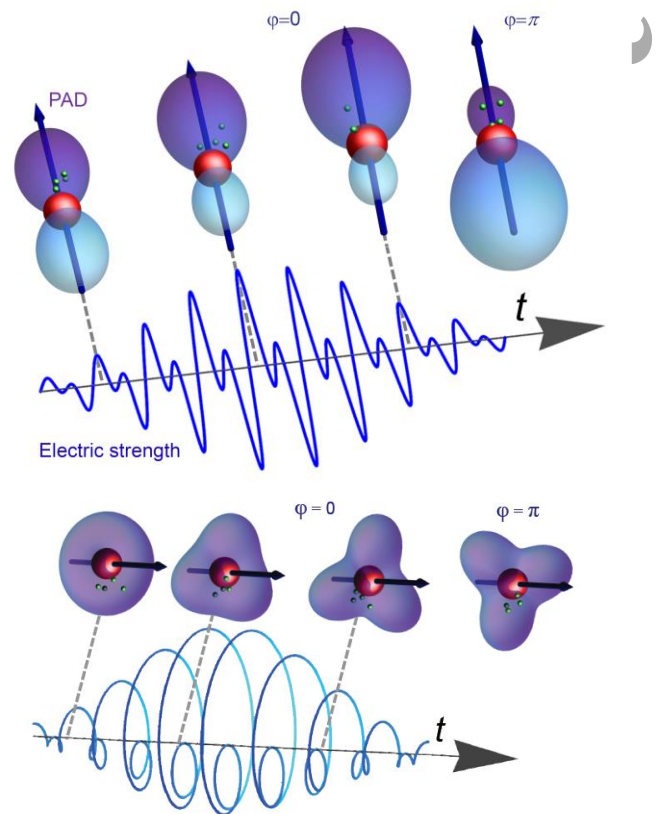


Figure 1 – Formation of typical PADs with time in the Ne(2p) ionization by coherent  $\omega+2\omega$  linearly (a) and circularly with opposite helicity (b) polarized pulses. For the incoherent pulses the PADs would be constant and symmetric, like in the beginning of the pulse. The evolution is shown for the relative phase  $\varphi = 0$  between the harmonics. The final PADs for  $\varphi = \pi$  illustrate their control by the relative phase of the harmonics. For details of the theoretical model see [3].

and X-ray femto- and attosecond pulses offer the unique possibility to resolve the ultrafast electronic dynamics in time and in space.

A first demonstration of coherent spectroscopy in the EUV spectral range, based on a four-wave mixing (FWM) process, was reported at the FERMI in 2015 [4]. Two coherent EUV pulses created a grating inside  $\text{SiO}_2$ , which diffracted a third infrared pulse. The FWM signal presented oscillation, which were attributed to the acoustic modes launched in the sample. Recently the technique has been extended to a pure four-wave mixing in the EUV range by the combination of three EUV pulses [5]. The possibility to introduce a delay between the first two pulses will offer an additional control tool towards the implementation of more complex multidimensional spectroscopy schemes [6].

#### Current and Future Challenges

The coherent control on the attosecond timescale of electronic processes calls for techniques enabling the

manipulation of the amplitude and phase of the spectral components composing the attosecond waveform. This is a formidable challenge as it requires the development of techniques suitable for bandwidth extending over more than tens of eV in the EUV and soft X-ray.

Complementary experimental approaches are currently on going using both table-top high-order harmonic generation (HHG) sources and FELs. They are strongly supported by theoretical works indicating possible implementation schemes and the advantages offered by time-resolved experiments in this spectral range for the observation of ultrafast electronic dynamics [7,8]. Due to the low conversion efficiency of the HHG process, attosecond pulse generated by table-top sources are still rather weak for the application of schemes based on the combination of different EUV pulses. Phase control of attosecond pulses either by metallic filter [9] or chirped mirrors [10] has been suggested, but both strategies suffer from a limited flexibility and the coupling between amplitude and phase shaping. Traditional approaches, well-established in the femtosecond domain such as pulse shapers based on liquid crystals or acoustic-optical modulators cannot be extended to the EUV domain, due to the lack of suitable materials for modifying the phase in this spectral range. However, alternative approaches such as phase-modulation spectroscopy combined with HHG offer a viable strategy for overcoming these limitations [11].

Rather than controlling the spectral components of the EUV pulses, a less direct, but more promising approach for the control of the relative phase is represented by the steering of the motion of the electron bunch (in the case of FELs) or of the electronic wave packet (in the case of HHG) leading to the generation of the EUV radiation. This was adopted in the experiment described above [2] for controlling the relative phase between the two coherent harmonics. The scaling of this approach to multiple-harmonics with complete control of the relative phase between the harmonics should be within reach for seeded FELs like FERMI. The combination of multiple harmonics would also open new perspectives in the synthesis of trains of attosecond pulses at FELs. The temporal characterisation of the phase differences could be achieved by measuring the photoionisation spectra by the combination of the EUV harmonics and a collinearly propagating infrared field. This approach is well established in the attosecond community for the temporal reconstruction of the attosecond pulse train generated by HHG in gases [12]. The extension to FEL-based harmonics should be possible, even though alternative approaches for overcoming the lack of sub-cycle synchronisation between the EUV harmonics and the IR field must be developed.

In HHG in gases, the electronic dynamics after the tunnelling ionisation process can be controlled by sculpting the driving field. Simulations for determining the optimal wave for HHG were reported [13], as well as the first experimental results based on the combination of different multi-color pulses in the visible and near-infrared spectral range [14]. By shaping the electric field of few-cycle, multi-color pulses (also with perpendicular polarisation) it is possible to precisely control the motion of the electronic wave packet, giving the opportunity to modify the recollision probability and the phase accumulated in the continuum. To which extent this control could be translated into a precise and independent manipulation of the amplitude and phase of the EUV components generated upon recombination remains an open question.

The description of coherent control in EUV poses serious challenges for theory even for isolated atoms. Involving the inner shells, relativistic effects, intermediate and autoionizing resonances, competition of photoionization with the Auger decay call for primordial many-electron approaches for description of the target. Contribution of higher multipoles of the electromagnetic field may be noticeable in the EUV and possible nondipole effects in the coherent control need to be considered. At the same time, a broad interval of field intensities, pulse durations, electron energies, – prohibit application of a universal theoretical toolkit and imply using various theoretical approaches for treating the time evolution. It ranges from different versions of the time-dependent perturbation theory and strong field approximation (the latter rather for EUV + optical pulses) to direct solution of the time-dependent Schrödinger and quantum Liouville equations. Since it is not possible to solve the many-electron equations explicitly, several approximate methods are under development (see also Section 1.1). A variety of time-propagation schemes have been considered too, although most of the examples are still restricted by a model of single-electron moving in a local effective potential.

The direct comparison between experimental and theoretical results is very difficult, because the former are obtained in a reaction volume, with not well defined peak intensities and profiles of the beams and pulse durations. Cases when it is necessary to apply different theoretical approaches for different sites and spectral intervals (for example, near resonances) are not exceptional. Probably this kind of difficulty is the main reason why currently only a very limited number of theoretical groups take on the tasks of direct comparison of their calculations with the experiment.

The coherent control of PADs potentially allows the extraction of new information on the ionization amplitudes and atomic continuum in the spirit of

‘complete experiments’ [15], not accessible so far by other means, in particular, with incoherent beams. The field would benefit from realization of coherent control on spin polarization and vector correlations of the reaction products. The former may be useful for the possible usage in an ultrafast spin switch.

### Advances in Science and Technology to Meet Challenges

For the extension of coherent spectroscopy techniques to the soft X-ray at FELs, it will be necessary to further investigate the generation of coherent multi-color schemes already adopted at seeded FEL. Self-seeding approaches deliver pulses with a well-defined spectral bandwidth even though large shot-to-shot fluctuations are observed [16]. Sub-cycle control would also offer new challenges, due to the shorter wavelength of the radiation. The method applied so far for the control of the relative delay with a precision of 3 attoseconds [2] appears to be scalable up to the soft X-ray range, with typical oscillation periods of a few and few tens of attoseconds.

For table-top HHG sources, the development of multi-color, synchronized few-cycle pulses is a prerequisite for sculpting the electric field waveform. Advances in optical parametric amplifiers operating in complementary spectral ranges represent a viable strategy for the synthesis of coherent fields extending from the visible to the near infrared. These pulses offer not only the possibility to generate more efficiently isolated attosecond pulses, but also to shape the characteristics of the EUV spectrum [17]. The independent manipulation of different spectral range of a supercontinuum generated in a hollow-fiber capillary has already been demonstrated and applied for the generation of isolated attosecond pulses [18]. Control of the EUV properties would require strategies similar to those adopted and well-investigated in the femtosecond domain such as the implementation of genetic algorithms and feedback loops [19].

### Concluding Remarks

The field of coherent control and coherent spectroscopy of electronic processes with a resolution from the complementarities between EUV pulses generated by HHG sources and FELs down to the attosecond timescale will strongly benefit in the upcoming years. Cross-fertilization between these two research fields is expected to give rise to dramatic improvements in our capability to initiate, steer and observe electronic processes in the core and valence shells with unprecedented temporal resolution.

### Acknowledgments –

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## 2 FEMTOSECOND MOLECULAR DYNAMICS

### 2.1 Time-resolved molecular photoelectron spectroscopy at a seeded FEL

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#### Status

The investigation of molecular reactions on their “natural” timescale (femtoseconds to picoseconds) is a sought-after goal not only in chemistry, but also in physics and material science, with possible extensions to biology.

Fundamental questions in ultrafast photodynamics of molecular systems, in particular in photoinduced chemical reactions, are: what is the timescale for the reaction? Is it possible to “catch” some intermediate species? What is exactly the mechanism? How many steps does it imply? Is it possible to manipulate and control the system by varying some instrumental parameters?

The exciting possibility of observing the evolution of molecules during chemical reactions is often called “the making of a molecular movie”, implying the goal of preparing a movie showing molecular motions with femtosecond up to picosecond time resolution and possibly electronic and/or geometrical structural resolution for all species involved.

The new free-electron laser sources (FELs), with the availability of very intense and very short pulses, have brought the above dream close to reality.

The main avenue in the investigation of photochemical reactions by time-resolved spectroscopic techniques at FELs is a pump-probe sequence, namely laser pumping the ground state of a molecule of interest to a suitable excited state, and probing the intermediate states thus formed with several different methods.

Previous investigation aiming at characterizing ultrafast molecular dynamical processes have been conducted at FELs, primarily with structural techniques such as photoelectron diffraction [1], and X-ray diffraction/scattering [2]. Although structural techniques such as electron or x-ray diffraction are very informative, sometimes a significant change in electronic structure results in a small or even negligible change in geometry. Furthermore, for such techniques lengthy and complex data analysis is required.

A much more direct way to characterize molecular dynamics even in a very short timescale would be to look at valence photoelectron spectra in combination with ion yield spectra, to follow photochemical

reaction pathways. Valence photoelectron spectra yield direct information of the electronic structure of the species under investigation and are able to reveal even subtle changes in response to external parameters such as e.g. temperature change or photoexcitation. However, to observe the first stages of a chemical process using photoelectron spectra, picosecond (ps) or femtosecond (fs) time resolution is needed, together with sufficient photon and electron kinetic energy resolution. The most suitable light sources are therefore FELs or high-harmonic generation (HHG). However, up to very recently this relatively simple approach has been hindered by the facts that at FELs the photon energy jitter is generally too large to give sufficiently energy-resolved spectra without heavy data treatment and/or averaging over several time delays, and with simple HHG sources the resolution is masked by the simultaneous presence of multiple harmonics. If harmonic filtering is employed, the resulting intensity is too low to follow subtle changes.

#### Current and Future Challenges

A recent very significant step forward in this direction is represented by the newly built FERMI FEL at the Elettra facility, Trieste, Italy [3]. This seeded source provides pulses with negligible photon energy jitter. In this way, photoemission spectra can be measured with sufficient resolution to identify very clearly ionization processes stemming from different electronic states. The seeding scheme used at FERMI also provides an IR pulse which is very accurately synchronized to the FEL pulse, because it is derived from the same Ti:Sa laser that drives the seed chain.

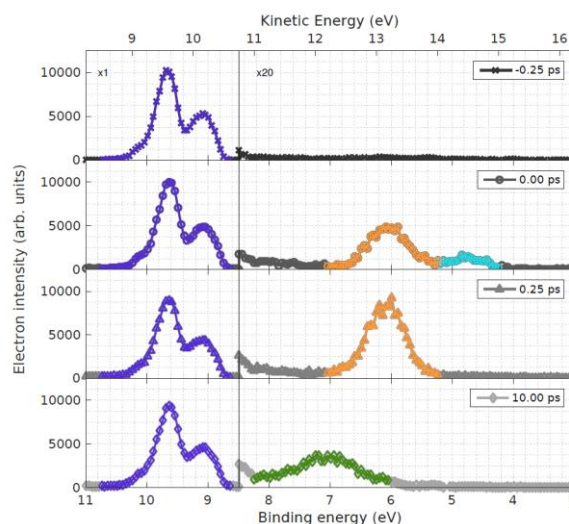


Fig.1. Valence photoelectron spectra of acetylacetone as a function of pump-probe delay (adapted from [4]).

Using the FERMI seeded source as probe in combination with the seed laser as pump, we have



performed time-resolved experiments on a prototypical system, acetylacetone [4].

In Figure 1 we show valence photoelectron spectra measured for different pump-probe delays. The blue curves in the higher binding energy region correspond to ground-state final states, while in the lower binding energy region spectral features are shown which correspond to the ionization of three different excited states:  $S_2$  (light blue),  $S_1$  (yellow), and  $T_1$  (green).

The time evolution was followed up to 1.5 ps pump-probe delay with 50-fs steps and to 200 ps with larger steps. A clear picture of the evolution of the system is reached, showing that the photoexcitation to the  $S_2$  ( $\pi\pi^*$ ) (bright) state is followed by a conical intersection (CI) connecting with the  $S_1$  ( $n\pi^*$ ) (dark) state, and then the  $T_1$  ( $\pi\pi^*$ ) state is reached through ultrafast  $S_1$  ( $n\pi^*$ )/ $T_2$  ( $n\pi^*$ ) crossing which is immediately followed by internal conversion to  $T_1$  ( $\pi\pi^*$ ). A minor pathway leading back to the ground state is also identified. Observed fragmentation yielding  $CH_x$  species is related to the onset of the  $T_1$  ( $\pi\pi^*$ ) state formation.

A schematic view of all intermediate states, their crossing paths and possible fragmentation pathways is shown in Fig.2.

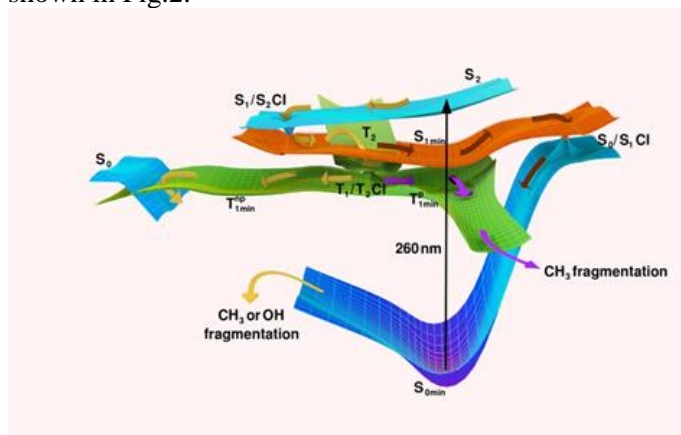


Fig.2. Schematic view of the photoexcited species which are formed and their time evolution (adapted from [4]).

A second experiment was performed under similar condition, allowing a real-time study of a ring-opening chemical reaction [5]. In general, it would be very useful to perform electron-ion coincidence measurements, to link a specific electron emission process to a specific ionic species. However, at FERMI the repetition rate available (50 Hz) does not allow this kind of data collection. Therefore a statistical method was used, namely partial covariance mapping, which allowed us to relate specific electronic final states to specific fragmentation processes.

In our opinion, this approach based on high-resolution valence spectra backed by high-level calculations and ion yield spectra is the ultimate way to shed light on fundamental, basic photo processes such as

photosynthesis, photovoltaic energy production, and vision.

## Advances in Science and Technology to Meet Challenges

While FERMI as the only operational seeded source allows obtaining unprecedented detailed information on molecular dynamics by measuring valence or core photoelectron spectra with a time resolution of 50 fs, a suitable FEL source with negligible photon energy jitter and very short pulses (few fs) does not yet exist. Furthermore, the photon energy range available at FERMI is at the moment limited to reach the carbon K-edge.

To overcome the problem of photon energy jitter, an obvious way would be to use coincidence techniques, to link one particular photoelectron emission to one or more specific ion species production. Coincidence techniques are nowadays available at FLASH, where a repetition rate of hundreds of Hz is possible [6]. Future plans at LCLS, Stanford, California, include the upgrade to LCLS II, with high repetition rate and self-seeding to reduce the photon energy jitter [7]. Also at the European X-FEL the repetition rate will be 27 kHz, sufficient to perform electron-multiion coincidence experiments [8].

Another planned evolution for the field is the possibility of going down to the attosecond regime in pulse length. Such opportunity will provide access at FELs not only to nuclear motion, but even to electron motion on its “natural” timescale. While the attosecond regime is already available at HHG sources, at FELs it will be possible to explore electron dynamics in the x-ray range, with high single-pulse power. Developments in this direction are announced at LCLS II [7] and European X-FEL [9].

## Concluding Remarks

The possibility of exploring in great detail photochemical processes on their natural timescale is an exciting development nowadays in place at FELs. In particular, the seeded source FERMI provides a negligible photon energy jitter which can be exploited to obtain time-resolved valence and more recently core level spectra with sufficient resolution to characterize reaction intermediate species. Future developments at several FELs include high repetition rate, to allow for electron-multiion coincidence measurements, and very short pulses, in the attosecond range, to have access not only to nuclear dynamics but even to electron dynamics.

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## 2.2 Chemical reaction dynamics from time-resolved X-ray scattering of free molecules -

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### Status

Chemical reactions are at the heart of chemistry and form the basis of a wide range of molecular science, engineering and medicine. Yet chemical reactions are difficult to understand on the most fundamental level, because they are inherently time-dependent processes of highly coupled, many-body systems. An ideal observation would determine the time dependence of atomic positions within reacting molecules, and the changes in the electron charge densities as chemical bonds break or form. Additionally, for reactions induced by coherent light pulses, important information could be gleaned from the shape of the wave packets in the multi-dimensional phase space. As a graphical representation of the time-evolving molecular structures, 'molecular movies' could aid the intuitive understanding of chemical dynamics.

Ultrafast time-resolved investigations aiming to explore chemical reactions while they unfold have traditionally employed spectroscopic methods. For all but the smallest molecules, the inversion of spectroscopic data to yield the desired chemical reaction movie is not possible. The aim of ultrafast time-resolved X-ray scattering experiments, as well as the complementary electron scattering, is to derive experimental, time-dependent maps of molecular structures. Beyond the immediate visualization of chemical dynamics useful to synthetic chemists, the molecular movie could reveal previously unseen structural motions associated with reactions, and the excited state structures could serve as benchmarks for the development of computational codes.

While in crystalline phases the molecular structure and the chemical dynamics are constrained by their crystal surroundings, reactions of free molecules in the gas phase can progress uninhibited, offering a direct comparison to calculated reaction pathways. Experimentally, gaseous samples are advantageous because the flowing gas avoids the degradation commonly seen with crystalline samples. Each laser shot hits a refreshed sample, eliminating the restriction to reversible processes. Drawbacks of studies on free molecules include the rotational averaging over the random, or only partially aligned, molecular samples, and their exceedingly low density compared to crystalline samples. The latter is countered by the immense X-ray intensity offered by modern X-ray Free Electron Lasers (XFEL).

Figure 1 shows the experimental concept. A molecular sample is intersected by an optical laser pulse that

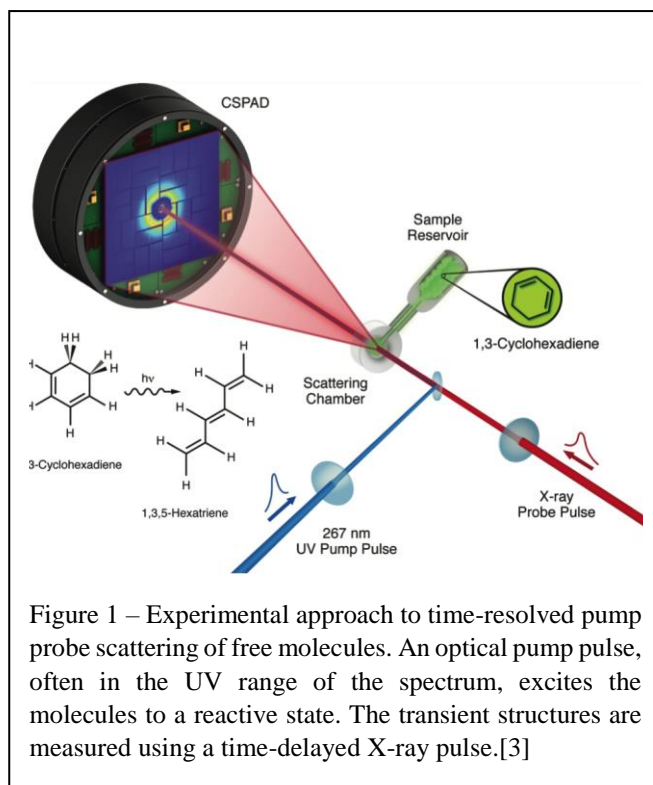


Figure 1 – Experimental approach to time-resolved pump probe scattering of free molecules. An optical pump pulse, often in the UV range of the spectrum, excites the molecules to a reactive state. The transient structures are measured using a time-delayed X-ray pulse.[3]

induces the chemical reaction. The time-evolving system is probed by a subsequent X-ray pulse. Scattering patterns are recorded by a large area detector and read out on a shot-by-shot basis. By binning the recorded scattering patterns according to the delay time between the laser pump and the X-ray probe pulses, a time resolution in the sub-100 fs regime can currently be achieved.[1,2]

### Current and Future Challenges

Ultrafast time-resolved X-ray scattering experiments have made possible the recording of a movie of the molecular structures while a molecule undergoes an electrocyclic ring opening reaction by bypassing a conical intersection.[3] Another example shows the wave packet motions of a diatomic molecule.[4] Further studies are in progress: movies of molecular vibrations, dissociation reactions, charge delocalization processes and many more. The success of these experiments depends critically on the level of optical excitation: if the pump laser pulse intensity is too high, multi-photon processes may lead the molecular system to a mix of poorly characterized, highly excited or ionized states. To place the molecule with confidence on a single reacting surface the pump intensity has to remain low, which leads to small changes in the scattering signal. Consequently, the dynamic range of the experiment must be such that changes in the scattering pattern of less than 0.1% can be recorded over the entire range of scattering vectors.

Photon energies up to about 10 keV can be conveniently obtained at the Linac Coherent Light Source (LCLS) at SLAC National Accelerator Laboratory. Using a large

area detector scattering vectors up to about  $5 \text{ \AA}^{-1}$  can then be accessed, allowing for the determination of excited-state structures with precision in the milli-Ångstrom range. Given the excellent signal to noise ratios obtained with pump-probe measurements, this is sufficient for detailed comparison to computed patterns, and thus for the extraction of the structural dynamics. But to characterize the chemical dynamics, one would ideally determine both the center and the shape of the time-evolving wave packets. Simulations show that to image the density distributions associated with traveling wave packets, scattering vectors up to  $15 \text{ \AA}^{-1}$  need to be accessed, see Figure 2.[5]

Theoretical models are required to accurately describe the scattering of molecules reacting on excited state surfaces. The total scattering signal of an ensemble of free molecules is described by the Fourier transform of electron pair correlation functions in molecules, which are constructed from two-particle density matrices.[6] Thus, a complete description of the total scattering signal consists both of elastic and inelastic components. In addition to the electronic components of the wave function, the nuclear rotational and vibrational wave functions must be included in the simulation of scattering signals.[7] For polyatomic molecules in excited vibrational and/or electronic states as appropriate for reacting molecules, the simulation of the scattering patterns therefore entails high-level calculations of the respective time-dependent nuclear structures and electronic states. Such a complete description has not yet been obtained.

### Advances in Science and Technology to Meet Challenges

Next generation XFELs, in particular LCLS-II, will provide higher photon energies and brightness. Photon energies in the 20 keV range would not only allow for more accurate structure determination, but also reveal details of the wave packet motions. Dedicated detectors that capture scattered radiation over a larger range of scattering angles would be important as well.

The higher brightness of LCLS-II is largely a consequence of the increased repetition rate. To take full advantage of the brightness, the experimental scheme must match the detector parameters to the repetition rate so as to enable read-out of individual pump-probe frames together with the shot-to-shot timing information. The high repetition rate coupled with the high photon energy, as may be possible with LCLS-II-HE, would then enable tremendously improved scattering measurements that could truly determine the time evolution of structure and chemical bonding during molecular reactions.

Limits to the temporal resolution at this time result mostly from the duration of the optical laser pulses.

Shorter pulses would enable the experiment to resolve higher frequency motions. Moreover, the extraction of spatial density distributions of traveling wave packets requires optical excitation with laser pulses of well-defined coherence parameters. Pulse shaping technologies exist, but have not yet been implemented at the CXI beam line at LCLS where the X-ray scattering experiments are performed. Typically UV photons are required to break chemical bonds, so that the pulse shaping technologies need to be implemented at the UV wavelengths.

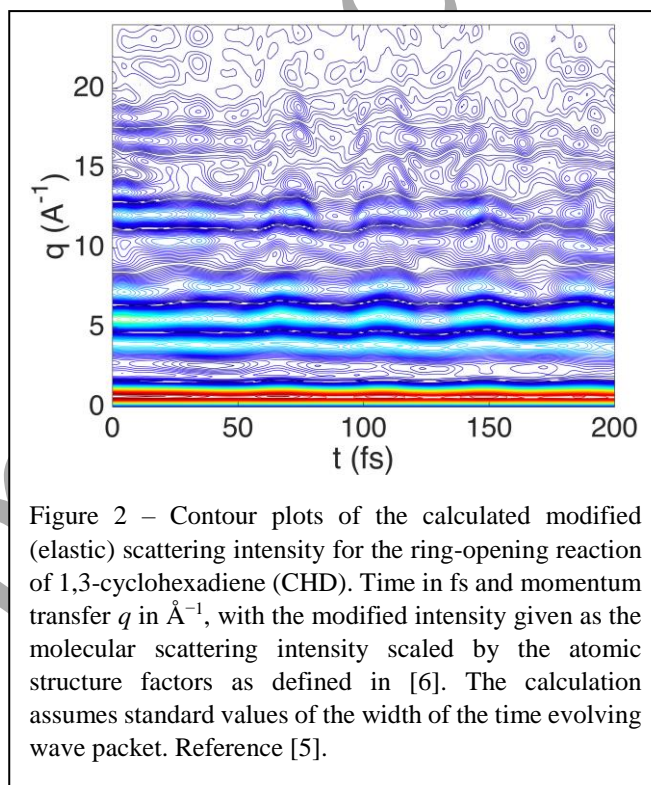


Figure 2 – Contour plots of the calculated modified (elastic) scattering intensity for the ring-opening reaction of 1,3-cyclohexadiene (CHD). Time in fs and momentum transfer  $q$  in  $\text{Å}^{-1}$ , with the modified intensity given as the molecular scattering intensity scaled by the atomic structure factors as defined in [6]. The calculation assumes standard values of the width of the time evolving wave packet. Reference [5].

Much information is lost in the averaging over the random orientation of free molecules. Alignment of the molecules, using any of the well-established methods, could reduce the information loss.[8] However, given the wide amplitude motions associated with pendular states, the common limitation of alignment schemes to one of the three rotational axes, and the additional experimental uncertainties associated with the alignment parameters, it seems doubtful that the benefit merits the added difficulty. A very different approach would instead record scattering patterns of individual molecules and sort the frames by orientation, similar to existing methods on microcrystals.[9] The high repetition rate of LCLS-II may enable such experiments, opening up an additional exciting frontier of single-molecule, time-resolved chemical dynamics X-ray scattering.

### Concluding Remarks

Experiments with ultrafast time resolution can probe processes on the inherent time scales of molecular systems. While spectroscopic experiments are

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2  
3 fundamentally limited by the time-energy uncertainty,  
4 scattering experiments are not: it is possible to  
5 simultaneously measure atomic positions and time  
6 without fundamental limitation. While technical  
7 challenges are substantial, many can be overcome so  
8 that much finer details of reaction dynamics could be  
9 uncovered. For example, it would be exciting to observe  
10 the complicated interplay of electron densities with  
11 geometrical structure as a molecule passes a conical  
12 intersection. Ultrafast time-resolved X-ray scattering is  
13 poised to image, in real time, the atomic rearrangements  
14 of molecules during chemical reactions. Foreseeable  
15 advances will add the imaging of the electron density  
16 distributions, so that it will be possible to watch the time  
17 evolution of chemical bonds – a chemist's dream come  
18 true.  
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### 2.3 Molecular self-imaging with attosecond electron wave packets.

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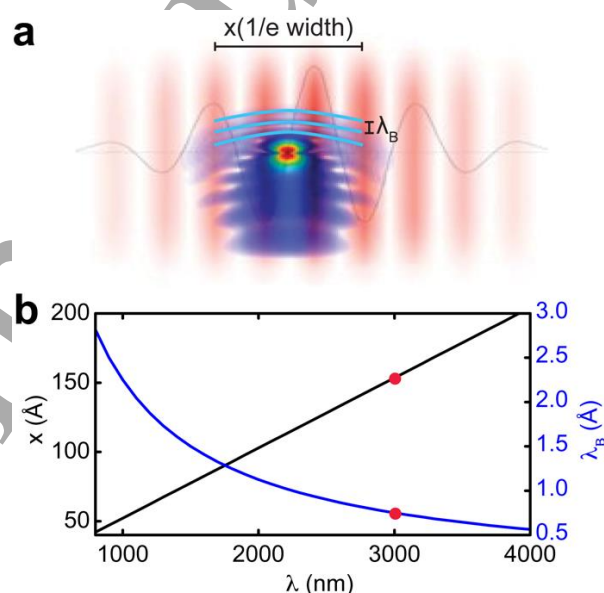
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#### Status and Background

Unravelling the mechanisms of chemical reactions requires imaging techniques that capture snapshots of the molecular structure on the nuclear (*i.e.* few femtosecond;  $1 \text{ fs} = 10^{-15} \text{ s}$ ) timescale with atomic (*i.e.* Ångstrom;  $1 \text{ Å} = 10^{-10} \text{ m}$ ) spatial resolution. Electron scattering techniques are preferable over photon scattering due to the small de Broglie wavelength ( $\lambda_B$ ) of electrons, providing picometre spatial resolution, and the orders-of-magnitude higher scattering cross section, which permit the imaging of single molecules[1]. Consequently, a wide range of scattering methods are utilized in molecular imaging studies, such as: ultrafast electron diffraction (UED)[2], relativistic electron diffraction (RED)[3], ultrafast electron microscopy (UEM)[4], laser-assisted electron diffraction (LAED)[5], and laser-induced electron diffraction (LIED)[6][7][8][9]. Achieving the desired *combined* space and time resolution has proven difficult as electrons interact with each other, thus limiting the achievable temporal resolution to time-scales on the hundreds of femtosecond range[10]. This problem is elegantly solved in LIED[11][12], which achieves the ultimate time-limit on the attosecond scale, by utilizing only a single attosecond electron wave packet (EWP). The quantum trajectory of the EWP is entirely determined by the laser field which, after roughly two thirds of a laser cycle[13][14], drives the EWP back to its origin where it self-interrogates the molecule's geometric structure with attosecond ( $1 \text{ as} = 10^{-18} \text{ s}$ ) and picometre ( $1 \text{ pm} = 10^{-12} \text{ m}$ ) resolution. It is worth noting that the entire process of LIED is well-described in the framework of laser-driven electron recollision[13][14]. Contrary to high harmonic generation, phase-matching is irrelevant for LIED, *i.e.* the rescattered electron signal is entirely dominated by the long quantum trajectory. An important aspect of LIED is its ability to extract the field-free doubly differential elastic scattering cross section based on the classical recollision model[7][15], despite measuring inside a strong laser field. A consequence of the strong field and launching the recollision EWP is the triggering of molecular dynamics that can be investigated using an intra-pulse

pump-probe scheme with mid-IR (MIR) LIED's intrinsic attosecond timing resolution[16].

An important aspect for the development of molecular imaging is the fact that many different molecular reaction pathways exist which complicate the image analysis. We have shown that this issue can be mitigated[8][17] through kinematic coincidence detection in combination with MIR high repetition rate sources[18][19]. MIR recollision[20] is crucial for LIED especially for larger molecules with a lower ionization potential, making it necessary to generate a recollision EWP in the quasi-static regime, and the required "hard" collision to pinpoint atomic locations. Mid-IR wavelengths are also key to image larger molecular structures with excellent spatial resolution.



**Fig. 1 | Molecular imaging with an attosecond electron wave packet.**

(a) An attosecond electron wave packet (EWP; blue) is generated from a molecular target by the electric field of a short laser pulse (red). After propagation during less than an optical cycle, the returning EWP (wave fronts highlighted as blue lines) images its target ion. The lateral spread ( $1/e$  width),  $x$ , and the de Broglie wavelength,  $\lambda_B$ , of the EWP's wave fronts are illustrated in panel (a) and plotted as a function of the driving laser wavelength,  $\lambda$ , in panel (b) with respective values highlighted at a driving laser wavelength of 3000 nm (red circles).

#### Current and Future Challenges

Fig. 1 shows the calculated scaling of the recollision wave packet's transverse extent with driving wavelength together with the de Broglie wavelength of the EWP. The graph indicates that MIR-LIED is in principle able to image molecular structures with large spatial extent; for comparison, the size of the C60 fullerene is roughly  $7 \text{ Å}$ . The large spatial extent of the recollision EWP occurs, however, with the penalty of decreasing scattering cross section ( $\sigma_s \propto \lambda^{-4}$ ) with

increasing wavelength. To overcome this challenge, the repetition rate of the MIR laser must be sufficiently large to acquire enough statistics in a reasonable measurement timeframe. An interesting feature of MIR-LIED is the intermediate kinetic energy regime in which LIED can operate, despite achieving sufficient momentum transfer for structural reconstruction. LIED electron impact energies range roughly between 50 to 200 eV for which hydrogen exhibits appreciable scattering cross section values. Consequently, we have been able to image hydrogen atoms[15][17] to which UED is typically blind due to the much higher impact energies used in UED. The enticing prospect in LIED is therefore imaging the ubiquitous proton motion, as demonstrated by Wolter *et al.*[8].

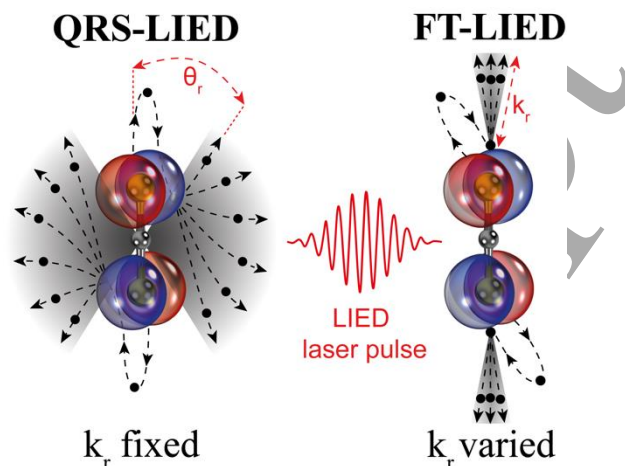
With respect to target complexity and size, the gas-phase measurement of larger molecules presents the difficulty of preparing these in a gas-phase molecular jet since larger structures are typically either in liquid or solid form. More importantly, the ionization potential decreases with larger molecules, and an adequate balance must be struck between photoionization to launch the recollision electron and saturating ionization for a given target species.

Future challenges will arise in the LIED of more complex biomolecules, such as azo complexes or peptides. For example, retrieval algorithms are typically based on the quantitative rescattering (QRS) theory with the assumption of atoms being independent from each other (IAM). With the much increased structural complexity and larger molecular systems, the QRS approach may not be tractable since it requires the extrema to be found in a multi-dimensional solution space. An alternative approach to retrieve structural information is through the Fourier transform (FT) variant of LIED, called FT-LIED, which does not rely on prior structural information or retrieval algorithms. These two straightforward variants of LIED are shown in Fig. 2.

An important further prospect of LIED is the clear possibility to make a full attosecond resolved molecular movie. An obvious choice would be the change of LIED wavelength, but it would only change the time range by a few femtoseconds whilst dramatically changing the recollision dynamics. A better solution could be pump-probe time-resolved LIED in which a UV pulse prepares the molecular state and an MIR pulse images its structure at various different pump-probe delays.

### Concluding Remarks

MIR-LIED has shown its potential to image isolated molecules with a single electron, thus achieving a combined attosecond temporal and picometre spatial resolution. The strength of the LIED method is access to



**Fig. 2 | Variants of LIED.** In QRS-LIED (FT-LIED), the attosecond recollision EWP is investigated at fixed (varied) rescattering momentum,  $k_r$ , and varied (fixed) rescattering angles,  $\theta_r$ . The grey shaded area indicates the rescattering region of interest that the LIED method utilizes in the extraction of structural information. In the case of FT-LIED, the Fourier transform of the back-scattered ( $\theta_r \approx 180^\circ$ ) attosecond EWP provides an image of the object.

the first femtoseconds during which the course of a chemical reaction is determined, and we firmly believe that many future investigations will be able to utilize these strengths in pump—probe studies of complex molecular systems.

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021034.



## 2.4 Molecular dynamics in intense FEL fields - Nora Berrah<sup>1</sup> and Edwin Kukk<sup>2</sup>

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### Status

The first x-ray free electron laser (FEL) to reach the inner atomic shells was the Linac Coherent Light Source (LCLS) that has since been used to investigate femtosecond photo-induced molecular dynamics such as H-migration [1], non-linear physics [2,3], or ring opening [4] which are important physical and chemical processes [5]. The next x-ray FEL to become available was SACLA in Japan, where owing to its unique ability to reach deep atomic inner shells, the investigations concentrated on probing the photoinduced dynamics of highly charged systems containing heavy elements [6,7]. Both of these FELs can deposit, during the few femtoseconds of an FEL pulse, a very large amount of energy into the system; as much as 87 keV into a single fullerene molecule, for example [2]. Such unique experimental conditions trigger also unique molecular dynamics (MD).

Several noteworthy x-ray FEL-based pump-probe investigations of MD have made direct observations of structural changes on femtosecond time-scale. A showcase experiment is the isomerization of acetylene into vinylidene, studied following valence ionization at FLASH FEL. At the LCLS, research with C 1s core ionization focused on determining the onset of the ultrafast deuterium migration in deuterated acetylene (DCCD), converting it into vinylidene (CCDD). That FEL-based work employed x-ray pump – x-ray probe technique with an x-ray split and delay tool. In that experiment, absorption of the first x-ray pulse is followed by Auger decay, initiating isomerization from DCCD<sup>2+</sup> to CCDD<sup>2+</sup>. The motion of D from one C site to the other was captured by the delayed x-ray probe pulse that caused a Coulomb explosion into four ions: C<sup>+</sup>, C<sup>+</sup>, D<sup>+</sup>, and D<sup>+</sup>. The time evolution of the ion momenta was found to be in the first 12 fs following core ionization [1], demonstrating that non-adiabatic interactions can couple nuclear motion to the electronic relaxation.

An experimental and theoretical investigation of the fs ionization dynamics of C<sub>60</sub> at high pulse intensity was carried out to understand the response of a relatively large, strongly-bonded molecule. C<sub>60</sub> was ionized above the C K-edge at 485 eV photon energy to reach conditions in which each atom in a C<sub>60</sub> molecule, in the x-ray focus, absorbs multiple photons. The core ionization with high-fluence x-ray induces the Auger process, resulting in many

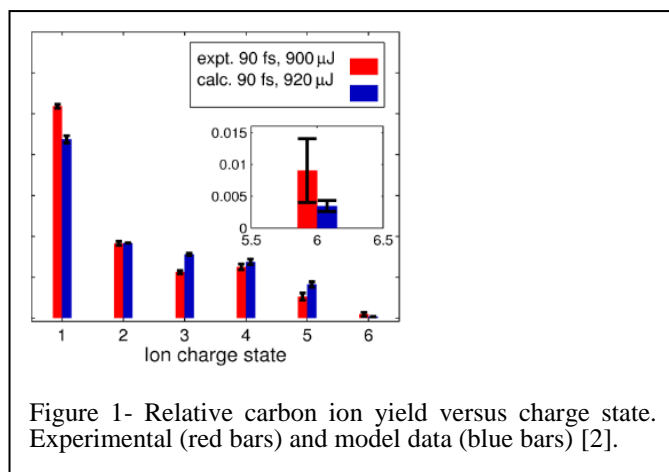


Figure 1- Relative carbon ion yield versus charge state. Experimental (red bars) and model data (blue bars) [2].

photo- and Auger electrons due to the cyclic multiphoton ionization. This process leads to secondary ionization of C<sub>60</sub> as well as its fragments ions by the photo- and Auger electrons. The C<sub>60</sub> molecule charges up to C<sub>60</sub><sup>8+</sup>, and because of the short C–C bond lengths, it fragments via Coulomb repulsion into molecular and C ion charge states distribution from C<sup>+</sup> to C<sup>6+</sup> at 90 fs pulse duration as shown in Figure 1 [2]. Theoretical MD simulations were carried out, as shown by the blue bars in Figure 1, revealing that strong secondary effects were responsible for the suppression of the ion yield for the highly charged C ions [2]. Namely, electrons were being recaptured by the highly charged ions. A recent time-resolved experiment was conducted in C<sub>60</sub> to investigate further this ionization and fragmentation mechanism.

These MD processes were triggered by soft x-ray core ionization in light elements such as carbon. In other MD experiments carried out at SACLA, multiphoton deep core ionization of iodine and xenon atoms led to generation of highly charged states by cascade Auger processes. The following two examples illustrate the ensuing dynamics in a small (di-iodomethane) and a much larger (Xe clusters) quantum system.

In iodine-containing organic molecules such as di-iodomethane, I is an absorption hotspot at x-ray energies above ~5 keV, leading to positive charge flow from iodine(s) to the entire molecule, concurrent with Coulomb explosion [8]. MD was studied using multi-ion coincident time-of-flight spectroscopy with full 3D momentum resolution. A tailored SCC-DFTB theoretical approach was developed [8] in order to successfully capture the dynamical aspect of the induced Coulomb explosion, reproducing the experimental kinetic energy distributions of fragment ion (Figure 2) and various other experimental fingerprints. Similar investigations of larger, iodine-substituted biomolecules such as 5-iodouracil [10] have a link to important applications in medicine, namely radiotherapy, since the very same “soup” of



molecular fragments and emitted electrons can be responsible for the destruction of a cancerous cell, via introducing and irradiating such a radio-sensitizer molecule.

In the next example, Xe clusters were studied in an FEL-pump near-infrared (NIR)-probe experiment at SACLA, again using ion time-of-flight spectroscopy to detect and analyze the atomic Xe fragments from Coulomb explosion of clusters [9]. In this much larger quantum system, the focus was on the collective phenomena such as the formation, expansion and evaporation of electronic nanoplasma. By probing the transient excited states of the atomic Xe fragments, the experiment was able to capture the ultrafast population ( $\sim 12$  fs) and depopulation ( $\sim 250$  fs) of these excited states during nanoplasma formation in Xe clusters. Here, not only photoionization and Auger processes, but also inelastic scattering and interatomic Coulomb

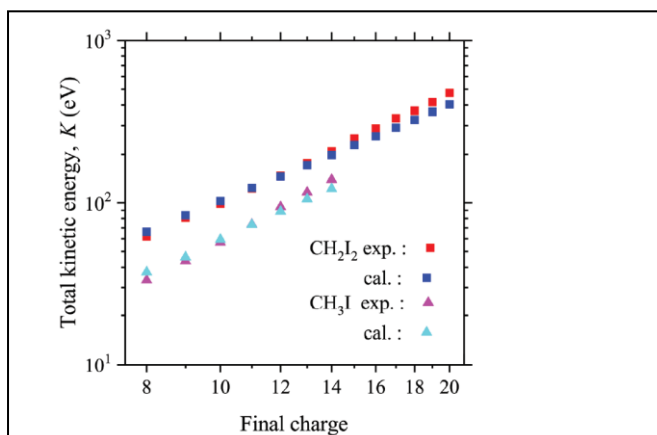


Figure 2 - Charge dependence of the total kinetic energy of fragment ions for CH<sub>3</sub>I and CH<sub>2</sub>I<sub>2</sub> molecules[8].

decay (ICD) come into play. The nanoplasma formation and MD are the two processes that cannot be treated in isolation, but in a combined description of the system. Such type of experiments have a direct relevance for the use of XFEL pulses in determining the structure of nanosized objects.

### Current and Future Challenges

X-ray FELs have already proven to be excellent tools for the investigation of MD, but are not without technical challenges. The self-amplified spontaneous emission (SASE) that produces FEL radiation at the LCLS and SACLA facilities showcased above, is noisy and presents photon energy jitter that makes it difficult to fine-tune to specific atomic resonances, for example, or to perform electron spectroscopy with high resolution. The most complete information about the photodissociating sample can in principle be gathered by multiparticle coincidence measurements, recording all created charged particles, and their momenta. In reality, this is impossible for all but the smallest

systems. However, combining electron and multi-ion detection is very beneficial also for much larger systems. Here, some practical challenges relate to the ability of the detection system to cope with the specific pattern of FEL operation – large number of particles can be produced and must be detected in a short interval followed by long periods of inactivity. On the one hand, this can make successful recording and disentanglement of coincident datasets overwhelmingly difficult; on the other hand, low repetition rates have made the collection times of more complex datasets prohibitively long.

Time-resolved pump-probe experiments that offer the most direct access to MD have their own additional challenges. For example, obtaining full and stable spatial and temporal overlap of the pump and the probe pulses has proven often to be time-consuming and failure-prone processes in the earlier experiments at LCLS and SACLA, partly due to the challenges in developing suitable and accurate calibration and diagnostics tools; a goal often requiring close collaboration between the facility and the scientists.

One of the next scientific challenges is to move beyond the simple showcase molecules and investigate complex photo-induced processes in a broad category of systems, that occur in the sub-fs and attosecond regimes.

### Advances in Science and Technology to Meet Challenges

A key trend in these complex experiments is that no deeper interpretation is possible without close theoretical support. Theoretical models, novel approaches and practical simulation tools such as XMDYN [2,5] are being developed at a fast pace. Experiment-wise, a variety of data analysis tools have been harnessed. Particularly in multiparticle detection experiments, even gradual and small improvements in technology can make a crucial difference – for example employing novel “funnel” type microchannel plate detectors could improve four-ion coincident detection efficiency more than five-fold.

The advances in FEL technology are being made at a very fast pace. Perhaps most crucial for molecular dynamics experiments are the dramatically increased repetition rates, new, highly flexible x-ray pump-x-ray probe creation schemes and improvements on the pulse jitter, such as using self-seeding and shortening the FEL pulse duration to under femtosecond range for accessing even faster dynamics.

### Concluding Remarks

The first decade of experiments on molecular dynamics using FELs in the soft- and hard x-ray regimes has already yielded a much more detailed insight into photoinduced ultrafast processes than has been possible with conventional light sources.

The next generation x-FELs such as the LCLS-II, the European X-FEL and others that provide high-frequency attosecond x-ray pulses and versatile pump-probe schemes will ensure that the FEL-based science meets the next challenges in the investigations of molecular dynamics.

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## 2.5 Interplay between electron dynamics and nuclear dynamics

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### Status

When an  $N_{\text{el}}$ -electron system is singly ionized, the remaining  $(N_{\text{el}}-1)$ -electron system may or may not be in a nonstationary state. This depends on the degree of coherence present in the  $(N_{\text{el}}-1)$ -electron system. The concept of coherence is defined through the reduced density matrix obtained by tracing the density matrix of the entire  $N_{\text{el}}$ -system over the degrees of freedom of the ionized electron. Thus, from the perspective of system-bath models, the  $(N_{\text{el}}-1)$ -electron system is the “system” and the ionized electron is the “bath”.

The topic of valence-electron dynamics following ionization has risen in importance because of the development of light sources that enable electron removal on a time scale that is comparable to the natural time scale for valence-hole dynamics. For example, the observation of valence-hole dynamics in Kr atoms was made possible via optical tunnel ionization using a strong quasi-single-cycle optical pulse [1]. These measurements demonstrated the need for a reduced-density-matrix description and revealed a degree of hole coherence of about 60% – limited by the finite duration of the ionizing pulse.

An alternative to optical tunnel ionization is one-photon ionization using an ultrashort XUV pulse from a high-harmonic-generation source or from a free-electron laser. Theory has demonstrated that in ultrafast XUV photoionization of atomic Xe, the maximum degree of coherence in the  $(N_{\text{el}}-1)$ -electron system is limited by the electron-electron-interaction-mediated entanglement between photoelectron and hole [2].

In polyatomic systems, an additional effect appears that limits electronic coherence: the motion of the atomic nuclei. In this situation, not only the photoelectron, but also the nuclei act as a bath for the  $(N_{\text{el}}-1)$ -electron system. In other words, in order to characterize the coherence properties of the hole, the density matrix of the entire system must be traced over the nuclear degrees of freedom as well.

In theoretical work [3, 4], it has been shown that, for a number of representative molecular systems, decoherence of an initially perfectly coherent hole wave packet [a wave packet of the  $(N_{\text{el}}-1)$ -electron system] is typically completed within just a few femtoseconds (see Fig. 1). The main reason for this effect is the (zero-point) nuclear motion in the ground-electronic state of the target molecule: Over the range of molecular geometries sampled by the nuclear

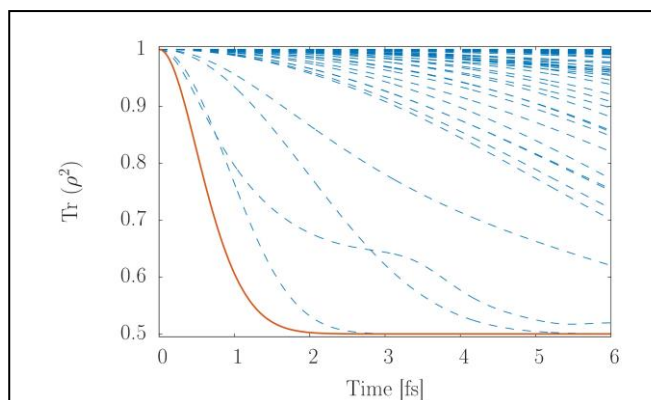


Figure 1 — A measure of electronic coherence (the so-called purity) is plotted as a function of time for a hole wave packet in phenylalanine involving two  $(N_{\text{el}}-1)$ -electron eigenstates. The dashed lines show the individual contributions to electronic decoherence from each of the 63 normal modes. The solid line shows the combined effect of all normal modes. A purity of 0.5 corresponds to a completely incoherent superposition of the two electronic states. This figure is taken from Ref. [4]. Copyright 2017, by the American Physical Society.

motion, the Born-Oppenheimer potential energy surfaces (BO PESs) of the  $(N_{\text{el}}-1)$ -electron eigenstates involved in the electronic wave packet are not perfectly parallel to each other. As a consequence, averaging over these nuclear geometries gives rise to a characteristic time scale for electronic dephasing.

In regions of nuclear configuration space where BO PESs intersect and so-called conical intersections are formed, the BO approximation breaks down and the question of electronic coherence becomes more subtle. If the electronic wave packet is formed at nuclear geometries sufficiently close to a conical intersection, the electronic phase information is not entirely lost before the molecule passes through the conical intersection. Then, the electronic coherence imprinted in the photoexcitation process can affect the time evolution of the molecular geometry and, hence, offers a potential route towards chemical control [5].

### Current and Future Challenges

Depending on whether the valence-electron removal is triggered by optical tunnel ionization or XUV one-photon absorption, different hurdles must be overcome when attempting a theoretical description of the associated dynamics. In the following, I will focus on XUV photoionization, under the assumption that the sudden approximation is valid. Under these conditions, it is possible to concentrate on the properties of the residual  $(N_{\text{el}}-1)$ -electron system.

Using an XUV photon, in principle any valence-hole state can be accessed. It is known that particularly in the inner-valence regime, strong electron-correlation effects predominate [6]. Therefore, computationally



expensive electronic many-body theory is required. So far, such calculations are available primarily near the equilibrium geometry of the target molecule in its ground state. However, in order to predict XUV-induced chemical dynamics, it is necessary to construct BO PESs over a wide range of nuclear geometries.

This is a challenge for the following reason. Typically, XUV photoionization causes molecular fragmentation. At the corresponding nuclear geometries, the Hartree-Fock reference state generally is no longer a good first-order approximation to the electronic ground state of the target molecule. The reason why this is significant is because standard many-body methods such as coupled-cluster theory or Green's functions typically rest on the assumption that the  $N_{\text{el}}$ -electron ground state can be approximated using a single Slater determinant, based on a mean-field framework such as ground-state Hartree-Fock.

In addition to the failure of the single-determinant approximation, there is the challenge that a rather large number of  $(N_{\text{el}}-1)$ -electron BO PESs must be computed. Consider a molecule consisting of  $N_{\text{at}}$  atoms. At the electronic mean-field level, we expect that the number of initially occupied molecular orbitals – and, therefore, the number of BO PESs reached by electron removal – scales as  $N_{\text{at}}$ . Thus, for a molecule containing ten second-row atoms, we would already need to calculate tens of BO PESs.

If one allows for configuration mixing of the one-hole configurations with the two-hole one-particle configurations, the number of potentially relevant BO PESs no longer scales merely as  $N_{\text{at}}$ , but, in the worst case, scales as  $N_{\text{at}}^3$ . Thus, in the ten-atom system alluded to above, it may be necessary to compute thousands of BO PESs. Note in this context that standard molecular configuration-interaction codes do not allow the user to compute more than a few electronic eigenenergies.

The high number of  $(N_{\text{el}}-1)$ -electron eigenstates generally required for describing XUV-induced dynamics has another consequence: The energy splittings between neighboring BO PESs tend to be small. This is illustrated in Fig. 2. Hence, conical intersections are no longer well-isolated entities, but there is a high density of conical intersections as a function of nuclear displacements. In other words, not only must one deal with severe electron-correlation effects, one must also cope with an increased incidence of non-Born-Oppenheimer processes [7, 8].

Further note that some, or even many, of the BO PESs may be associated with metastable states that can decay by spontaneous electron emission. This challenge is difficult to handle within standard Hermitian quantum mechanics, but may be treated, in

principle, using methods developed within the framework of non-Hermitian quantum mechanics [9]. However, such methods are still not routinely available in molecular electronic-structure codes.

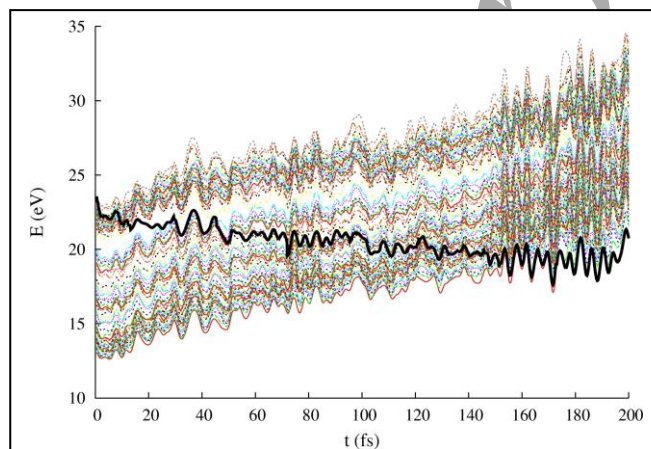


Figure 2 — Time evolution of the 63 one-hole outer-valence BO PESs of photoionized  $\text{H}^+(\text{H}_2\text{O})_{21}$ . The black curve shows the instantaneous electronic state, starting from the energetically highest outer-valence state. In order to reach the electronic ground state, the system has to move through a large number of conical intersections. Reproduced from Ref. [7] with permission from The Royal Society of Chemistry.

### Advances in Science and Technology to Meet Challenges

The challenges must be addressed using a two-pronged strategy. First, dedicated advances in electronic-structure theory are required. These must take into account strong electron-correlation effects and must provide access to a large number of electronic eigenstates characterized by a high spectral density. Methods based on the one-particle Green's function are a natural choice for the context described here. However, among other things, the limitations imposed by the Gell-Mann and Low theorem must be overcome.

Second, advances are required to handle the nuclear-dynamics problem. Mixed quantum-classical methods, which traditionally work in the adiabatic representation, lose their effectiveness when dealing with a quasi-continuum of electronic states. Pure quantum methods that rely on a local, diabatic representation (such as the vibronic-coupling Hamiltonian) are unsuitable for situations where a wide range of nuclear configuration space must be covered and where the density of conical intersections is high.

### Concluding Remarks

In a contribution to a roadmap of ultrafast x-ray atomic and molecular physics [10], I discussed challenges for theory in connection with high-intensity x-rays. In the

present contribution, I have been focusing on femtosecond dynamics triggered by photoionization of only one valence electron (a process that may be accompanied by the ejection of additional electrons). In spite of many past investigations in this area, there are still important unsolved problems, which offer exciting opportunities for future research.

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### 3 PHOTOIONIZATION

#### 3.1 Photoionization of ions –

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##### Status

Photoionization (PI) is the removal of one (or several) electron(s) from an atom or an ion by a single photon either directly or in a multi-step process involving intermediate autoionizing states. Indirect PI mechanisms include multiple excitation of outer-shell electrons with subsequent electron ejection and the production of an inner-shell vacancy followed by Auger decay(s).

PI and the related spectroscopy are behind many of the ground-breaking advances in modern physics. Besides fundamental atomic-physics insights, studies of PI provide key information for applications in many fields of research such as plasma physics, the physics of planetary atmospheres and astrophysics including research on interstellar and intergalactic media.

Experiments on PI of ions are impeded by the difficulty to provide dense pure targets of ions preferentially at low temperatures to avoid Doppler broadening. Relatively high densities can be provided by laser-produced plasmas or in ion traps. The conceptually cleanest experiments use ion targets in the form of accelerated beams. Advantages of photon-ion interacting-beams arrangements are the relatively straightforward measurement of absolute PI cross sections, the selectivity in providing pure ensembles of mass-over-charge analyzed ions and the low ion temperature that allows for high-precision spectroscopy. This contribution focuses on the photon-ion merged-beams technique [1].

The first interacting-beams PI experiments were conducted with negative ions looking into the release of the least-bound electron for which filtered optical light from a bright projection lamp was sufficient. A boost of PI experiments with negative ions came with the development of optical and UV lasers. PI of positive ions requires higher photon energies and, hence, the first merged-beams experiments with positive ions had to await the development of bright synchrotron radiation sources. A review of work on PI of positive and negative ions until 2005 has been provided by Kjeldsen [2]. Presently, only two experimental facilities using merged beams are regularly producing PI results, MAIA [3] at SOLEIL, and PIPE [4] at PETRA III.

The bulk of previous experiments with ions has yielded total single ionization cross sections for valence-shell PI. Due to limitations in photon energy and photon flux not many measurements on deep inner shells could be reported. The development of new high-brilliance synchrotron radiation sources, particularly SOLEIL and PETRA III, provides possibilities to extend PI experiments to higher photon energies and to processes with

smaller cross sections. In addition, spectroscopy of electrons released by PI of ions is now feasible [3]. PETRA III is presently the most brilliant synchrotron radiation source. Small bandwidth photon beams at energies  $E_\gamma$  of 250 to 3000 eV are available for PIPE experiments while the upper limit of  $E_\gamma$  at MAIA is 1000 eV with much smaller photon flux.

##### Current and Future Challenges

Presently, great efforts are devoted to the detailed understanding of multi-electron processes such as multiple ionization by a single photon accompanied and followed by multiple-electron excitations and decays. Partial absolute cross sections for multiple ionization associated with deep inner shells are being determined with the number  $\nu$  of electrons removed from the target ion by a single photon reaching up to  $\nu=6$  [5].

As an example for multi-electron processes, non-sequential triple-Auger decay in a four-electron interaction was unambiguously observed with  $C^+$  ions [6] which provide the minimum number of electrons in the  $L$  shell to facilitate the process. The outstanding sensitivity of the PIPE setup with its strong suppression of detector background has made it possible to measure cross sections as small as a few hundred barns. This sensitivity facilitated the observation of the direct triple-Auger decay.

Direct removal of two electrons by a single photon was recently observed with  $F^-$  ions. Simultaneous ionization of the  $K$  and  $L$  shells were detected in the triple ionization channel on top of a contribution arising from  $K$ -shell ionization with subsequent emission of two electrons. The cross section for direct double ionization was measured over a wide energy range [7].

Such studies on multi-electron processes can be supported by the coincident detection and spectroscopy of electrons and ions subsequent to PI. An example is the investigation of the Auger decay of the  $4d \rightarrow nf$  ( $n=4,5$ ) resonances in the  $Xe^{5+}$  ion [8].

The theoretical description of multi-electron processes in atoms and ions is challenged by the new experiments (see, e.g., [5]). New approaches have to be developed and tested against measured data. Changing the ion charge state along isoelectronic sequences provides a "knob" in such experiments by which the relative importance of electron-electron versus electron-nucleus interactions can be tuned for better insight into the fundamental mechanisms.

Resonances are usually the dominant features in PI cross sections. High-precision spectroscopy and accurate determination of resonance energies is essential for the detailed understanding of atomic structures and decay dynamics. Applications, for example in astrophysics need very accurate level energies. Theory can provide such energies with small uncertainties for two- and three-electron systems which can serve as refer-



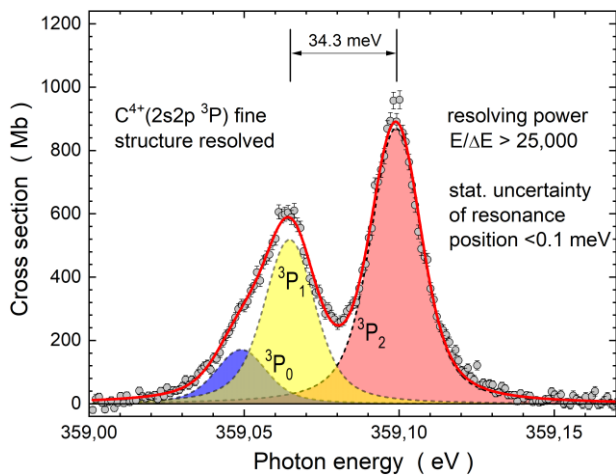


Figure 1 — High-resolution observation of fine structure within the  $2s2p\ ^3P$  resonance term in the absolute PI cross section of metastable  $C^{4+}(1s2s\ ^3S_1)$  ions. The solid line is the result of theory. The  $^3P_J$  ( $J=0,1,2$ ) contributions are shown by the dashed lines. The experiment is shifted by 1.4 meV to match the accurate calculation [9].

ence standards. An example for a high precision experiment is shown in Fig. 1 where measured and calculated cross sections for PI of metastable  $C^{4+}(1s2s\ ^3S)$  ions via  $2s2p$  resonances are compared to one another [9]. The photon energy bandwidth was less than 14 meV and the statistical precision of the energy determination was better than 0.1 meV. Such numbers provide a promising basis for meeting a great challenge for the future, namely, the excitation of  $U^{89+}(1s^22s_{1/2} \rightarrow 1s^22p_{1/2})$  at about 280 eV with sub-meV accuracy which would allow for a critical test of quantum electrodynamics and could eventually lead to new physics beyond the standard model.

Core-level excitation of atoms and ions has become increasingly more accessible with the new light sources providing the basis for interpreting site-specific excitation of molecules and molecular ions. X-ray absorption spectroscopy of molecular ions is opening new avenues to study isolated radicals which are produced in an ion source, accelerated and made to form a fast beam. Such beams of electrically charged molecular fragments can be mass and charge analyzed and can be transported to the photon-ion interaction region before they decay. Flight times between the production of ions and their arrival at the interaction region are typically 1 to 10  $\mu s$ . First experiments with molecular ions including (endohedral) fullerenes have been successfully conducted [1] and point the way to further exploration of this immense field of research. PIPE is equipped for measurements on heavy molecules and nanoparticles with masses up to 100,000 u at 1.5 keV beam energy.

PI of neutral atoms such as C and O is extremely important in astrophysics. Accurate level energies and cross sections are required in order to fully exploit modern satellite-based observations. Targets of neutral C and O are not readily available and call for the appli-

cation of the merged-beams technique. Intense neutral beams can be made available by photodetachment of negative ions within intense laser fields. Thus the present experiments with ions can be extended to atoms which cannot directly be accessed by other methods.

### Advances in Science and Technology to Meet Challenges

The most obvious requirement for meeting future challenges is the continuation of light-source development. High brilliance at high energies with very narrow bandwidths together with high stability in position, photon energy and beam handling are the keys to new improved PI experiments.

For extending the range of feasible measurements, efforts in target development are necessary. Several high-performance sources of charged particles are available for installation on the ion-source platform at PIPE. State-of-the-art sources of cluster ions and highly charged atomic ions should be made available for PI experiments. While the existing magnetic spectrometer provides pure beams as far as mass and charge of the ions are concerned also the control and manipulation of the electronic state of the target ions is highly desirable. This will require combinations of radiation fields (lasers), evaporation ovens, charge-exchange gas cells and trapping techniques. Ions and particularly molecular ions can be de-excited and cooled by storage in a suitable ring, preferentially at cryogenic temperatures. This requirement calls for a combination of a synchrotron radiation source with an ion storage ring, both being major pieces of equipment. An alternative is emerging from the laser-based development of portable short wavelength radiation sources which can be integrated in existing heavy-ion storage rings. A proof of principle experiment on PI of  $C^+$  is already envisaged at CRYRING in Darmstadt [10]. Very high charge states and ultimately even  $U^{89+}$  can be made available for photo-absorption experiments by employing high-energy heavy-ion storage rings.

Advances in detection capabilities are as instrumental for meeting future challenges as is the enhanced performance of particle and light sources. Electron detection in photon-ion experiments should be further improved. Particularly with respect to PI of electrically charged molecules (radicals), clusters and nanoparticles the detector should be capable of separating (fragmentation) products emerging from the photon-ion interaction region under different angles and at different kinetic energies. Suitable combinations of magnetic, electrical, and time-of-flight separation fields together with appropriate detectors have to be devised to reach this goal.

### Concluding Remarks

Given the numerous applications and their increasing demands for a better understanding of PI processes along with the ongoing developments of light sources, target arrangements, detection devices, and experimental schemes the field will remain very active for decades and provide new fundamental insights as well as information required for the numerous fields of PI applications. Even new physics beyond the standard model may emerge from high precision photo-absorption experiments.

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### 3.2 Molecular-frame photoelectron angular distribution from static to time resolved –

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#### Status

Photoionization (PI) of molecules is at the heart of many fundamental processes induced by absorption of VUV to X-ray photons, as well the dominant response to intense laser fields. The molecular-frame photoelectron angular distribution (MFPAD), i.e., the photoelectron emission pattern in the non-spherical potential of the molecule, for each orientation of the molecule relative to the propagation or polarization axis of the ionizing radiation, stands out as the most detailed observable characterizing molecular photoionization stereodynamics. Measuring MFPADs for single photon ionization in the dipole limit is equivalent to resolving the interference pattern of the partial waves from which the electron wave function is composed at each energy: it allows for the most complete determination of the complex ionization matrix elements, i.e. their magnitude and phase for each asymptotic partial wave [1]. A number of recent reviews [2-4 and references therein] give a broad summary of the progress in MFPAD research achieved through 2012, and describe the success of MFPAD studies in probing the electronic dynamics of photoionization and subsequent dissociation, as well as intramolecular dynamics and the relaxation of excited molecular states.

The principal challenge of experimental MFPAD studies is the establishment of the molecular frame. Taking advantage of dissociative photoionization (DPI) in an assembly of randomly oriented molecules, powerful electron-ion coincidence techniques that provide the 3D momentum vector of each charged particle have yielded a rich harvest of results since the mid 1990's (See also Sec. 3.4). For small molecules the molecular orientation relative to the light polarization axis is inferred from the recoil ion fragment momenta. Then the photoelectron momenta measured in coincidence for each DPI event provide MFPADs for any polarization geometry. The molecular frame (MF) is easily established when the dissociation time is short relative to other motions of the molecule. These methods have strongly benefited from the properties of third generation synchrotron radiation including wavelength tunability up to the X-ray range, spectral resolution, exotic linear to circular

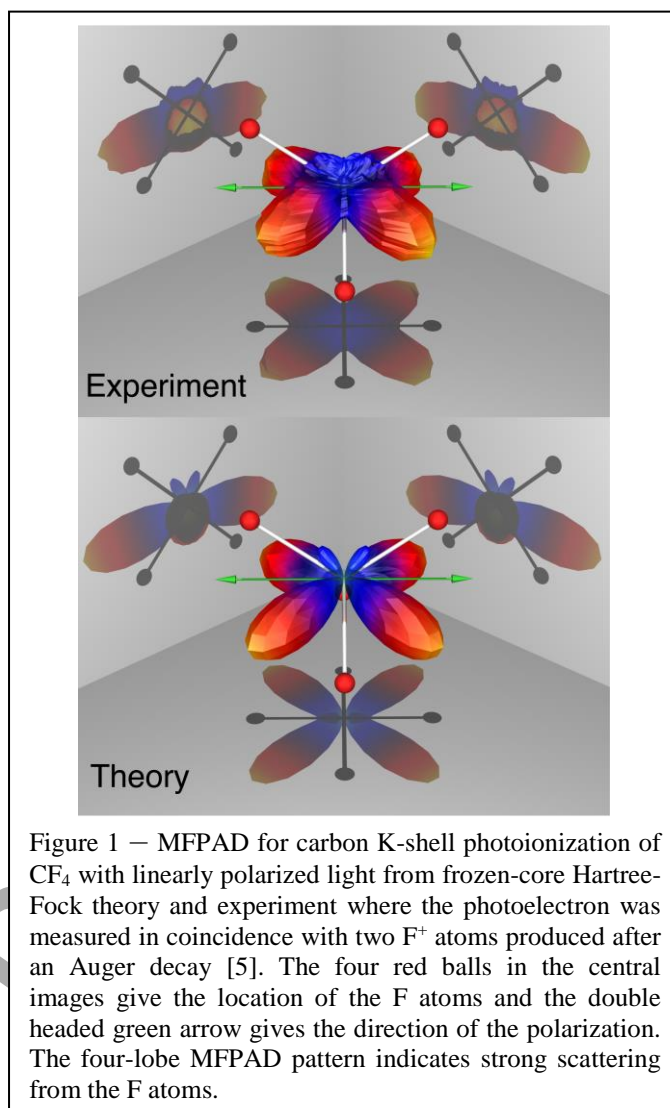


Figure 1 – MFPAD for carbon K-shell photoionization of  $\text{CF}_4$  with linearly polarized light from frozen-core Hartree-Fock theory and experiment where the photoelectron was measured in coincidence with two  $\text{F}^+$  atoms produced after an Auger decay [5]. The four red balls in the central images give the location of the F atoms and the double headed green arrow gives the direction of the polarization. The four-lobe MFPAD pattern indicates strong scattering from the F atoms.

polarization and a repetition rate in the MHz range, which is a good match for coincidence techniques. MFPAD studies have revealed remarkable features of photoionization dynamics [2-4], such as orbital angular momentum properties of the photoelectron temporarily trapped in a shape resonance, quantum interferences between autoionization and direct ionization channels, gerade/ungerade symmetry breaking, circular dichroism in the angular distributions (CDAD) or the multi-electron nature of the ionization process due to the breakdown of Koopmans' theorem. Core-shell photoionization is particularly suited for probing MFPADs in small polyatomics with the use of coincidence techniques since core-shell ionization is followed by rapid electronic relaxation or Auger decay producing dissociation into multiple charged ionic fragments. Their momentum-resolved detection in coincidence with the photoelectron give direct access to the MFPADs, as shown in Fig. 1 for  $\text{CF}_4$ , or to recoil-frame photoelectron angular distributions when only two fragments are produced. Site and element specific core-shell photoionization processes have provided the basis for the study of fundamental

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2  
3 quantum processes including two-source interference,  
4 coherence, and core-hole localization. Furthermore, far  
5 from the ionization threshold, core-shell ionization  
6 MFPADs map the diffraction patterns arising from the  
7 superposition of the partial waves emitted from a  
8 single center within the molecule and scattered by the  
9 molecular potential, therefore reflecting the  
10 geometrical structure of the molecule [6]. The  
11 sensitivity of MFPADs even to very small geometry  
12 changes has been demonstrated through vibrationally  
13 resolved measurements. For future applications, it is  
14 stressed that measured MFPADs also encode the  
15 complete polarization state of the ionizing radiation  
16 [7].

17  
18 An alternative route to measure MFPADs, required in  
19 particular to investigate non-dissociative ionization  
20 channels, or ionization from complex molecules, relies  
21 on the laser control of the spatial alignment of  
22 molecules in the laboratory frame (LF). Different  
23 strategies have been developed [2-4], from multiphoton  
24 ionization of small molecules through rotationally-  
25 resolved states by nanosecond lasers or via resonant  
26 electronic states by femtosecond lasers, extended to  
27 adiabatic or impulsive strong field alignment relying  
28 on short nonionizing laser pulses [8]. The latter  
29 methods allow the confinement of the molecular axes  
30 about a well-defined direction in the LF, in the  
31 presence of the laser field or in field-free conditions  
32 e.g., at recurrence times of the rotational wavepacket.  
33 Photoionization induced by a second femtosecond  
34 pulse (optical laser, high-order harmonic generation  
35 (HHG) or free electron laser (FEL)) at a controlled  
36 delay therefore provides PADs which converge on the  
37 MF distributions when the degree of alignment  
38 approaches unity. These PADs are mostly recorded  
39 using velocity-map imaging (VMI), where the velocity  
40 distribution of the photoelectrons is projected onto the  
41 2D plane of the detector. The progress in molecular  
42 alignment and orientation has boosted the investigation  
43 of MFPADs induced by intense near infrared  
44 femtosecond laser pulses [9].

45  
46 MFPAD observables allow a comparison between  
47 experiment and theory at the most detailed level,  
48 probing the theoretical descriptions of molecular  
49 photoionization [2-4]. Such comparisons have been  
50 facilitated by the development of the functional forms  
51 of MFPADs which capture the range of possible  
52 MFPAD angular dependence and allow for filtering out  
53 unphysical behavior in measured data [4]. The  
54 essential challenge in predicting MFPADs is the  
55 calculation of the continuum electronic states, which  
56 are not described by the methods of bound-state  
57 quantum chemistry. A review of well-established  
58 methods that have overcome these difficulties has been  
59 given in [4]. If the process is dominated by a single

ion target state then relatively simple Hartree-Fock like  
methods can be used, as seen in Fig. 1. Otherwise a  
coupled-channel formalism is needed such as the  
multi-channel frozen-core Hartree-Fock or time-  
dependent density functional theory. If a more  
complete treatment of correlation including target  
response is needed then R-matrix or complex Kohn  
methods should be used. Almost all calculations have  
been done within the fixed nuclei approximation for  
the photoelectron dynamics. Thus even when nuclear  
motion is included, the continuum electron dynamics  
are assumed to respond adiabatically to any change in  
molecular geometry.

### Current and Future Challenges

A number of challenges in MF-resolved  
photoionization involve time-resolved experiments. A  
general method for MFPAD determination was  
demonstrated by measuring LFPADs for XUV  
photoionization of impulsively aligned  $N_2$  molecules,  
as a function of time during the evolution of the  
launched rotational wave packet [10].

#### *Nuclear and electronic dynamics from the picosecond to the attosecond time scale*

Time-resolved MFPADs (TR-MFPADs) at the femto-  
to-picosecond time-scale provide a powerful tool for  
studying relaxation dynamics of molecular excited  
states using pump-probe techniques, where a short  
laser pulse prepares a vibronic wavepacket and a time-  
delayed pulse photoionizes the molecule. Taking  
advantage of field-free alignment, combined  
experimental and theoretical studies of the relaxation  
of the  $CS_2^*(C)$  excited state, chosen as a prototype  
example of non-adiabatic couplings at conical  
intersections, spotlighted that TR-MFPADs capture the  
pure valence-electron dynamics of the evolving  
wavepacket [11]. Current challenges address such  
processes in molecules of increasing complexity.  
Nowadays, the advent of bright and tunable FELs  
delivering intense femtosecond XUV to X-ray pulses  
[12] opens new avenues for MF- resolved studies at the  
femtosecond timescale, using e.g., XUV-pump XUV-  
probe schemes. Moreover, the availability of optical  
lasers at XFEL facilities, combined with the  
demonstration of femtosecond photoelectron  
diffraction on laser-aligned molecules [13], paves the  
way towards time-resolved studies of molecular  
structure during ultrafast dissociation, with angstrom  
spatial resolution.

Extending further the pump-probe approach,  
computations have shown that TR-MFPADs can probe  
electron dynamics or charge migration initiated by a  
pump pulse at the atto-femtosecond time-scale, i.e.,  
before the onset of the nuclear motion, and reflect the  
spatial localization of the electron density [14]. In the

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3 model presented in Ref. [14], excitation by a few  
4 femtosecond UV pump prepares a non-stationary  
5 coherent superposition of molecular states while the  
6 system is probed by sudden ionization using a delayed  
7 attosecond XUV pulse: such a scheme poses another  
8 challenge for experiments.  
9

#### 10 *Photoemission time delays in the molecular frame*

11 Most recent developments, calling for the  
12 characterization of the ionization dynamics at its  
13 relevant attosecond time-scale, focus on the  
14 determination of the photoemission time-delays (or  
15 Wigner delays), as demonstrated in recent years for  
16 atoms [15]. For molecules, the challenge is to measure  
17 the spatial dependence of photoionization delays  
18 probing the electron wavepacket scattered in the non-  
19 spherical molecular potential. Beyond the  
20 determination of the magnitude and the relative phases  
21 of the ionization matrix elements provided by the  
22 MFPADs at fixed photon energies, accessing the  
23 Wigner time-delays characterizing XUV ionization of  
24 a molecule requires that the energy dependence of  
25 those phases is obtained across the explored energy  
26 range. This can be achieved by performing angle-  
27 resolved RABBITT [16] (reconstruction of attosecond  
28 beating by interference of two-photon transitions),  
29 where ionization by an attosecond XUV pulse train  
30 probed by a coherent IR pulse ( $\omega$ ) leads to the  
31 production of sidebands in the photoelectron spectra  
32 resulting from the interference of two ( $XUV_{2n-1}+IR$  and  
33  $XUV_{2n+1}-IR$ ) paths. A first achievement in MF  
34 angularly resolved studies relying on the determination  
35 of a molecular Stereo Wigner Time Delay (SWTD),  
36 i.e. the relative delay between electrons emitted in the  
37 direction of, or opposite to, one end of an asymmetric  
38 diatomic molecule, was demonstrated in recent  
39 measurements combining a RABBITT scheme and a  
40 COLTRIMS reaction microscope set-up [17].  
41 Extracting the spatially resolved Wigner delays for  
42 each emission direction ( $\theta, \phi$ ) in the MF, and well  
43 defined orientation of the molecule in the LF, stands as  
44 a current challenge.  
45  
46  
47

#### 48 *Photoionization of molecules exposed to strong fields*

49 Within the flourishing developments of attosecond  
50 science, coincidence and/or molecular alignment  
51 approaches to molecular frame photoemission are  
52 being gathered to face the challenge of providing  
53 detailed insights into the multiple electronic nature of  
54 strong-field ionization of molecules and related  
55 ultrafast phenomena, including tunnel [18] and above  
56 threshold [19] ionization, electron recollision [20],  
57 high-order harmonic generation [3], or laser induced  
58 electron diffraction [21].  
59

#### 60 *One-photon synchrotron based studies*

One anticipated development in synchrotron based  
MFPAD studies is the highly differential angular  
distribution of two electrons in the continuum,  
resulting from one-photon double ionization, where the  
emitted photoelectrons are strongly correlated (See  
also Sec. 3.3). Beyond double ionization of  $H_2$ , results  
on one-photon double ionization MFPADs for  
diatomic and polyatomic molecules are scarce. From a  
theoretical perspective, the angular distributions in the  
double continuum are sensitive to subtle effects in  
initial state correlation and to strong continuum  
correlation effects when the two electrons have  
relatively low kinetic energy. The extension of  
methods developed earlier for atomic systems and  $H_2$   
to molecules such as  $H_2O$  requires inclusion of the  
interaction of the continuum electrons with those left in  
the ionic core, involving several low lying doubly-  
charged states. Additionally, many more partial waves  
are needed to obtain converged scattering results. Such  
a study will provide unparalleled details of the nature  
of correlation in the ground state electronic wave  
function of the target.

Another challenge will be the measurement of  
MFPADs in the tender to hard X-ray range [12]  
including the study of non-dipole effects, or ultrafast  
relaxation dynamics at femto or sub-femtosecond scale  
based on core-hole lifetime internal clocks, extending  
e.g., soft X-ray studies of temporally resolved resonant  
Auger MFPADs [22] (See also Sec. 3.4).

#### **Advances in Science and Technology to Meet Challenges**

Combining progress in electron-ion coincidence 3D  
momentum spectrometry, velocity map imaging and  
tomography, as well as expanding techniques for  
alignment and orientation of complex molecules, with  
the unprecedented opportunities offered by the  
development of free electron lasers delivering VUV to  
X-ray intense femtosecond pulses and the availability  
of laser-based HHG sources delivering attosecond  
XUV pulses, is at the core of the technological  
advances needed to fulfill the challenges posed by the  
MFPAD control of time-resolved electronic and  
vibronic wavepacket dynamics in gas phase molecules.  
The substantial ongoing increase of their repetition rate  
reaching few tens of kHz up to the 100 kHz range is  
crucial for future challenges. MFPAD studies should  
also benefit from the availability of variable  
polarizations.

Although it is possible to understand multiphoton and  
strong-field dynamics in terms of time-independent  
energy eigenfunctions, there is a large class of  
problems involving short light pulses for which this  
approach is not adequate and a direct solution of the  
time-dependent Schrodinger equation is needed.



Although there have been a number of theoretical methods developed to treat this problem, there have generally been compromises simplifying either the dynamics of bound electrons or the continuum or both. The combination of an adequate treatment of multichannel and correlation effects with an accurate time-dependent wave packet treatment of the continuum remains a challenge and is an active area of research.

### Concluding Remarks

The measurement of MFPADs led to a breakthrough in the understanding of molecular photoionization dynamics through spectrally resolved studies. MFPADs are now becoming significant probes of time-resolved electron and nuclear dynamics at the shortest time-scales highlighting links between adjacent fields involving ionization in different regimes of light-matter interactions.

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### 3.3 Photo-double ionization of atoms and molecules – Paola Bolognesi and Lorenzo Avaldi, CNR-Istituto di Struttura della Materia

#### Status

Double and multiple electron emissions from an atom or a molecule by the absorption of a single photon represent minor processes in the radiation-matter interaction. However, they are of extreme interest, because they are completely determined by electron correlation, which plays a crucial role in several processes and in determining physical and chemical properties in a wide class of materials. Furthermore, a proper treatment of two-particle correlation is a central problem in many-body physics [1,2]. This has motivated the large experimental and theoretical efforts devoted in the last thirty years to the understanding of the photo-double ionization, PDI, of two electron systems, the He atom and the  $H_2$  molecule.

In the case of the He atom a single photon with energy larger than 79 eV, the double ionization potential, produces a bare nucleus and two free electrons. With the polarization axis of the incident radiation defining a natural quantization axis the symmetry of the free electron pair, defined by the dipole selection rules, and the Coulomb interaction determine the differential cross section which describes the correlated emission of the pair. The experimental challenge represented by the low cross section (about 1kbarn) of the process has been overcome on one hand by the advent of the high intensity third generation synchrotron radiation sources and on the other by the development of set-ups that allow the simultaneous detection of the two electrons with several energies over a large solid angle or one/two electrons in coincidence with the recoil ion. The combined experimental and theoretical efforts have led to a clear identification of three basic mechanisms of PDI. In two of them the photon couples with the system formed by one electron and the nucleus and the second one is emitted either in a knock-out or a shake-off process. The relevance of the former (latter) decreases (increases) with the incident energy. The energy sharing between the two ejected electrons evolves from a flat to a U-shape distribution from threshold to higher energy (Figure 1). At the same time the correlated angular distribution changes from a “butterfly” shapes, determined by the  $^1P^0$  symmetry of the wave function and the Coulomb interaction in the continuum, to a pure dipole/isotropic shape for the fast/slow electron, respectively. In these mechanisms, the nucleus actively participates and gains a large recoil momentum. A third predicted mechanism, which involves the photon absorption by two electrons spatially close and the nucleus acting as spectator, has been observed only recently[3]. The main feature of the process is the emission of the two electrons with the same energy and

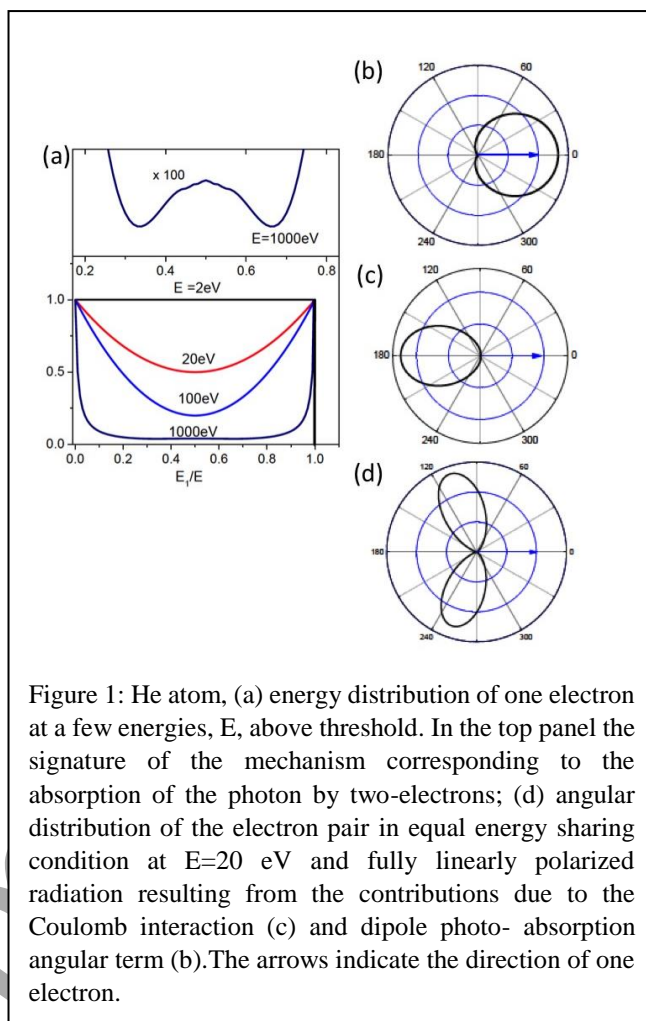


Figure 1: He atom, (a) energy distribution of one electron at a few energies,  $E$ , above threshold. In the top panel the signature of the mechanism corresponding to the absorption of the photon by two-electrons; (d) angular distribution of the electron pair in equal energy sharing condition at  $E=20$  eV and fully linearly polarized radiation resulting from the contributions due to the Coulomb interaction (c) and dipole photo-absorption angular term (b). The arrows indicate the direction of one electron.

back-to-back configuration, a geometry strictly forbidden in a dipole transition but allowed in quadrupole transitions, which contribute to about 0.1% of the total double photoionization cross section. In the case of the  $H_2$  molecule, the other two electron system, a few factors make the understanding of PDI even more challenging. The ground state electronic configuration is determined by the two-center nuclear potential, the dissociative final state makes the double ionisation threshold not unique and results in a Coulomb explosion of the two protons. Thus, a kinematically complete description of the break-up process calls for the detection of at least one of the protons. Such an approach allows the reconstruction of the molecular geometry when the PDI event occurred, i.e. to measure the correlated angular distribution of the electron pair in the molecular frame. The results of the few realised experiments [4,5] and their interpretation by theories based on fully numerical approaches show that the angular distribution of the pair is strongly affected by the molecular orientation as well as by the internuclear distance, i.e. by the electron-ion interaction.

#### Current and Future Challenges

In the case of He the dynamic of PDI is fully determined by the final state effects. Nevertheless, an effect of the

target electronic structure, due to radial extent and oscillations of the target orbital of the singly charged ion, has been predicted and observed in the angular correlation pattern of the PDI of Mg to  $\text{Mg}^{2+}$  ( $3s^{-2}$ ) [6]. It has been shown that the angular distribution bears the information for electron pairs “originated” in different regions of the charge density. This has been the first prove that PDI experiments can access initial state effects. The possibility to investigate electron correlation in the molecular target wavefunction and therefore to distinguish among molecular wavefunctions on the base of the different treatment of electron-electron correlation was the motivation for the proposal of PDI experiments in the eighties of the last century.

In the case of molecules however the correlated emission of electron pairs from a multi-center Coulomb potential is a severe challenge for theory as well as for experiments. In the case of diatomic molecules different from  $\text{H}_2$  as well as several polyatomic molecules the lowest dication states are metastable with barriers to dissociation. This prevents complete experiments and only experiments on randomly oriented samples can be done [7]. Exploratory studies with momentum imaging techniques have been performed on higher dissociative states of  $\text{H}_2\text{O}$  and interpreted via classical trajectories on computed potential energy surfaces [8]. These exploratory experiments have shown the richness of the process in more complex systems and are expected to stimulate new activity.

## Advances in Science and Technology to Meet Challenges

In many-electron systems, the effective electron-electron interaction is not known a-priori, but builds up during the photoexcitation process in a dynamic behavior determined by the active surrounding of the emitting atom/molecule. In these cases, an interesting scenario for the emission of an electron pair is that the photon excites one electron that senses its environment for accessible scattering channels (elastic, plasmonic, phononic, etc). By a proper choice of the energies and angles of the electron pair, PDI enables us to zoom into those channels where electron-electron interaction is operational. An example is the electron-electron interaction mediated by charge density fluctuations in a confined geometry, which has been studied in the case of  $\text{C}_{60}$  [9] (Figure 2). The dominant pathway identified for the process is the following: a valence electron absorbs the photon and rescatters inelastically from multipolar collective modes (plasmons) that mediate the coherent emission of a second electron. These pioneering results show that the measurement of energy and angle resolved electron pair opens the opportunity to explore different regimes of electron correlations,

including Coulombic scattering, local field effects and dynamical screening.

Another aspect only partially explored is the dynamics of the electron correlation in the time domain. While the dynamics of the interaction of the ejected electron with the residual ion in single ionization has been investigated with attosecond pulses, only one attempt exists for double ionization [10]. In that work two weak XUV and IR pulses have been used to double ionize the Xe atom via both sequential and non-sequential processes. In equal energy sharing conditions the non-sequential process dominates and the observed quantum interference, due to two paths to the same final dication state, indicates that the two-electron wave packet remains coherent during the measurement. The

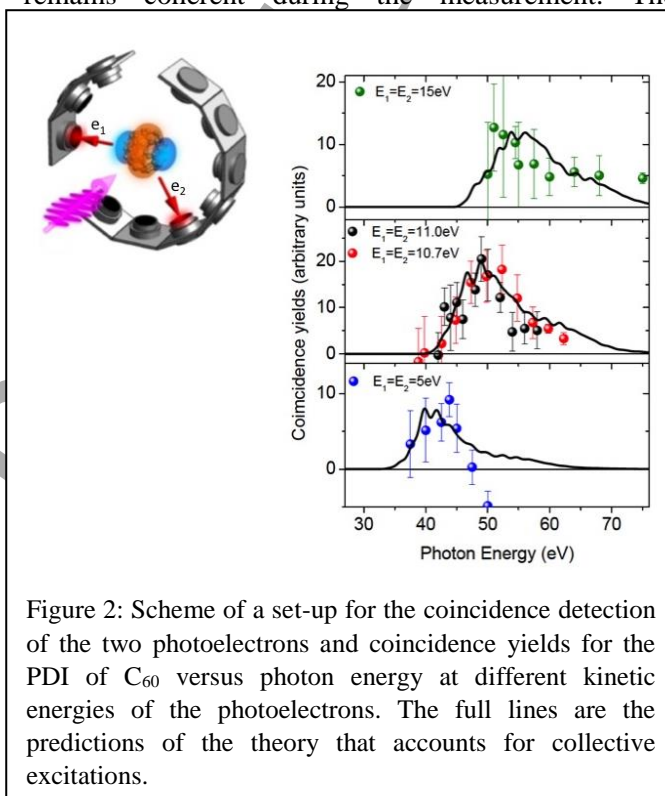


Figure 2: Scheme of a set-up for the coincidence detection of the two photoelectrons and coincidence yields for the PDI of  $\text{C}_{60}$  versus photon energy at different kinetic energies of the photoelectrons. The full lines are the predictions of the theory that accounts for collective excitations.

signature of the electron pair dynamics is encoded in the interference signal variable with the delay between the XUV and IR pulses. The results and their interpretation by lowest order many-body perturbation-theory have provided insights in the different mechanisms (sequential/non-sequential, shake off/knock-out) of PDI. The actual repetition rate of ultrafast lasers hampers these studies, but further development in the sources and more efficient coincidence techniques may allow full energy and angle resolution.

## Concluding Remarks

Electron correlation plays a crucial role in determining physical and chemical properties in a wide class of materials that exhibit fascinating properties (e.g. high temperature superconductivity, colossal magneto resistance, metal insulator or ferromagnetic anti

ferromagnetic phase transitions, self assembly and quantum size effects as well as in biology). PDI, with its unique property to be fully dependent on electron correlations, has provided and will provide the most suited tool to unveil their different facets. The advances in imaging techniques will allow the extension of the range of many-electron systems where kinematically complete experiments are achievable, while high repetition rate and high intensity FEL and HHG sources will fully disclose the investigation of electron correlation in the time domain.

**Acknowledgments** – The support of the FOE-Eurofel project to the PDI activity at Elettra is acknowledged

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### 3.4 Photo-induced reaction dynamics probed by COLTRIMS

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#### Status

The experimental toolbox for investigating light-matter interaction is well equipped nowadays. Ultrashort, ultrastrong laser pulses, free electron laser light and synchrotron radiation provide light of different characteristic features – with each radiation source covering unique aspects to drive quantum dynamics in small quantum systems in a fully controlled way. The dynamic response of atoms, molecules and clusters to such irradiation is extremely rich and often leads to many reaction products, such as several electrons and ionic fragments. COLTRIMS reaction microscopes [1] allow the measurement of the momentum vectors of all these products in coincidence independent of their initial emission direction.

The first experiments on photon interactions pioneering this technique dealt with atomic double ionization and Compton scattering (see references in [1]) investigating electron-electron correlation. Soon after, these instruments were used to examine photoionization dynamics of small molecules [2,3]: After molecular photoionization, the resulting molecular ions often fragment rapidly along steeply repulsive potential energy curves. In these case the molecules do not rotate before fragmentation, and thus the emission directions of the ionic fragments coincide with the spatial orientation of molecular bonds at the instant of photoabsorption (see. Fig.1). This opened a new era of molecular photoionization studies, as such experiments accessed the photoelectron and Auger electron angular distributions in the molecular frame of reference. Such experiments elucidated, for example proton migration in small carbohydrates [5] and pioneering work by Davies et al. [6] demonstrated the temporal evolution of such photoelectron emission patterns in  $\text{NO}_2$  employing a laser pump-probe approach. The unprecedented detail and completeness of these studies unveiled exciting new physical effects, which are often hidden to non-coincident spectroscopy techniques [7]. Photoelectron waves emitted from a well localized K-shell have been used to illuminate molecules from within [8]. This photoelectron diffraction is an extremely sensitive probe of molecular geometry. Even small asymmetries between otherwise equal bonds in  $\text{CO}_2$  lead to observable asymmetries in the electron angular distributions [9] and the angular distributions are sensitive to the vibrational motion of the molecule. Molecular frame angular distributions obtained upon Sulphur 1s-ionization of  $\text{CS}_2$  demonstrated that the effect of corehole localization can depend explicitly on

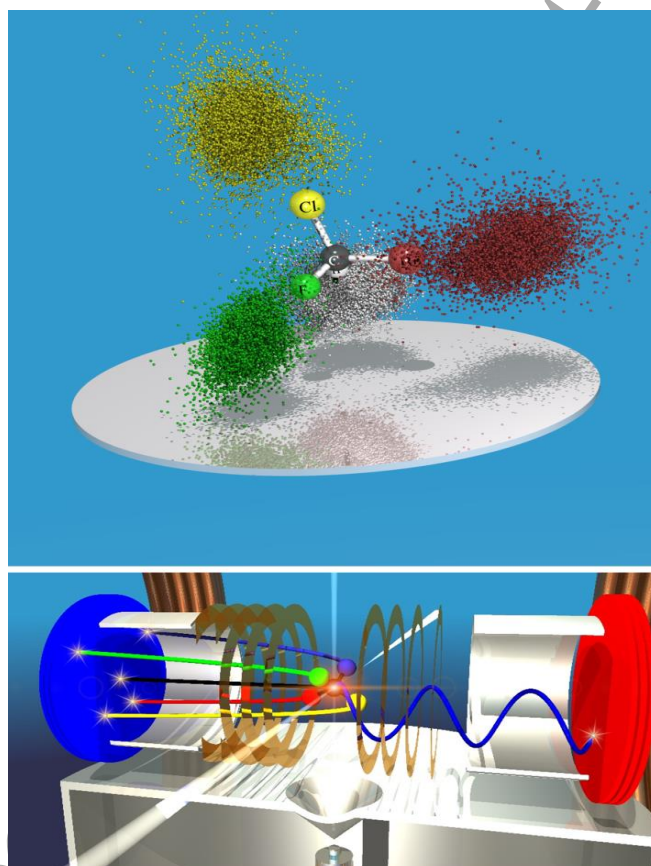


Figure 1: Upper panel: The molecular geometry ( $\text{CHBrClF}$ ) can be inferred from measuring the emission directions of the ions after a Coulomb explosion triggered by K-ionization and multiple Auger decay (adapted from [4]). Lower panel: scheme of a COLTRIMS Reaction Microscope.

the reaction dynamics [10]. In weakly bound systems particularly rich nuclear dynamics takes place during interatomic Coulombic decay (ICD). ICD is a non-local electronic decay mechanism with decay times in the regime of several 10 femtoseconds to picoseconds, i.e. on the same time scale as typical nuclear motion. Imaging studies inferred the mean decay width from these nuclear dynamics [11] and showed the vibrational motion of the decaying state [12].

#### Current and Future Challenges

One of the current frontiers of this field is to extend the studies of photon driven reactions to the time domain. Accessing the time domain is traditionally done by employing pump-probe approaches. The advent of higher repetition rate lasers and free electron lasers bears the promise that elaborate multi-coincidence studies with COLTRIMS can be performed, but the technical challenges are still daunting. Up to today, for example, only a few true ion-electron coincidence imaging experiments performed at a free electron laser have been reported despite a decade of efforts [13]. Furthermore, while ultrashort light pulses are obviously a key tool to investigate dynamics taking place on the shortest time scales, they have an intrinsic energy uncertainty. While such (coherent) broadband features

may even generate dynamical behavior, they can in many cases prohibit the preparation of distinct initial states or impede triggering definite processes.

The pump-probe approach is the most mature but by far not the only technique to access the atomic or molecular time domain. Recently, new schemes have been developed which fully exploit the power of coincidence detection. These do not require a probe step, which is often disruptive or at least perturbing to the process under investigation. Instead, internal clocks provided by the ionization process or the

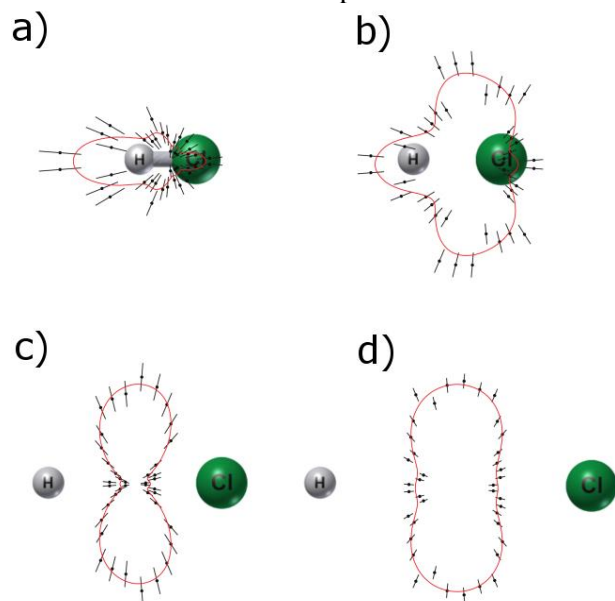


Figure 2: Auger electron angular distribution in the molecular frame as the HCl molecule undergoes the transition from a) molecule to d) separated atoms (adapted from [14])

molecule itself are exploited. One such ultrafast intrinsic clock, for example, is nuclear dynamics occurring on a known repulsive potential energy curve of an excited molecule. The prime example here is ultrafast dissociation of molecules, where narrow bandwidth excitation initiates a rapid dissociation and the electronic decay dynamics can be inferred from the kinetic energies of the final state ions and electrons. According studies [14] have depicted the temporal evolution of a molecular orbital during the dissociation process (see Fig. 2). A second example utilizes a naturally occurring streaking field [12]. If a (comparably) fast secondary electron is emitted shortly after a low energy photoelectron both electrons will interact in the continuum and the photoelectron is typically decelerated due to post collision interaction (PCI). As the strength of this interaction depends on the time between the emission of the two electrons, the decay time is encoded in the energy loss of the primary photoelectron. Possible decay processes, which may drive this internal clock also include the Auger decay. Accordingly, time scales of a few femtoseconds become accessible by using this scheme termed “PCI streaking”.

Such internal electronic or nuclear motion clocks can be exploited using the COLTRIMS coincidence detection. Here one of the detected particles, e.g. a photoelectron in case of PCI streaking or a proton when examining ultrafast dissociation, is used to decipher the decay time on a single event basis. The temporal evolution of properties of other particles detected in coincidence can then be investigated. The PCI streaking approach has been employed, for example, to study ICD in real time in helium dimers, showing the contraction and vibration of the dimer during the decay [12].

A further challenge besides approaching the time domain is exploring the role of quantum entanglement in photon driven processes. The ground state of the atom or molecule under investigation is a single quantum state consisting of many interacting particles. In the gas phase photo break-up of an atom or molecule is a fully coherent process occurring without coupling to any environment. Thus, analogous to parametric down conversion, where entangled photon pairs are created by splitting one photon into two, it is to be expected that the photofragments emitted from one joint initial state can share more than purely classical correlations and form multi-particle Bell-states. These quantum features become accessible only by coincident detection of all the reaction products. Despite the fundamental nature of this many-particle entanglement, so far only few cases demonstrated it directly in an experiment. For example, the entanglement of two electrons emitted from molecular hydrogen has been demonstrated by proving that Cohen/Fano double slit interference does not occur on the individual electron but is observable in the momentum distribution of the electron pair. Already ten years ago, the formation of a Bell-state consisting of the photoelectron and the Auger electron emitted upon K-shell ionization of  $N_2$  has been shown. This unexpected finding finally resolved the controversial debate whether K-shell ionization of equivalent sites of a homonuclear molecule leads to a localized or delocalized hole. It was shown that the question is quantum mechanically ill posed and the localization properties are defined by the quantum entanglement between all emission products of the photoreaction [7]. For larger molecules and more multifaceted decay channels, however, the issue of quantum entanglement is much more complex and calls for in-depth theoretical work along with more refined coincidence experiments. As of today there is not even a clear quantitative measure of entanglement, which is well adapted to the observables occurring in studies of molecules and to continuous variables like emission angles and momenta in the continuum.

### Advances in Science and Technology to Meet Challenges

COLTRIMS coincidence experiments have been driven by major technological advances in the past. They rely on projecting charged particles onto position and time-sensitive detectors (see Fig. 1). Thus, one of the important drivers is detector technology. As low energy charged particles are to be detected, there is currently no alternative to multichannel plates in combination with some sort of position readout. The decisive parameters are detection efficiency (which is more and more crucial as the number of coincidentally detected particles increases), time resolution, position resolution and multiple hit capability. Much progress in terms of efficiency has been made by so-called funneled MCPs or ones with special coating and by implementing imaging spectrometers that have a larger transmission. In terms of time resolution, however, 4<sup>th</sup> generation synchrotron radiation sources counteract the efforts increasing the detectors time resolution, by the increased length of the photon pulses. The second domain of COLTRIMS where progress can be expected (and has partially already started) are next generation imaging spectrometers. Today, the most basic analyzers rely on homogenous electric field regions and parallel magnetic fields. Great potential lies in more sophisticated electron/ion optics schemes. The goal is to spread out particles in time and to focus different starting position in the interaction region to the same time-of-flight and onto the same impact position on the detector.

The tremendous increase of photon flux at synchrotrons and (X)FELs has already opened the route to extend gas phase coincidence studies to much higher photon energies. These were for long inaccessible because of the steeply decreasing cross section. With today's high photon fluxes coincidence experiment become possible for processes with cross section below  $10^{-24}\text{cm}^2$ . Among the opportunities anticipated in high photon energy regime is the possibility to approach the Born limit of photoionization. In this limit photoionization becomes a prime tool for imaging of electronic wave functions, because once the ejected electron can be described by a plane wave, photoionization directly maps the Fourier transform of the bound state onto the detector [15]. Compton scattering is another process, which becomes accessible at high energies. Coincidence experiments on Compton scattering in the gas phase are mostly virgin territory. They promise a rich scientific harvest as Compton scattering gives access to the bound orbital, as well. For the future it will become possible to pair the photon detectors at XFELs, which are used today for photon diffraction with a COLTRIMS. This will allow imaging of the Compton scattered photons in coincidence with molecular fragments and secondary electrons. If this is used in a pump-probe experiment it will finally yield molecular movies of the time evolution

of both, the electronic wave functions and simultaneously the nuclear motion.

A revolutionary new era of studies on fundamental atomic and molecular quantum processes will emerge once "orbital angular momentum" (OAM) photon beams (or "vortex beams", "twisted light") become available. By manipulating the wave front of a photon field a phase modulation is created, which equals that of a photon carrying orbital angular momenta (i.e. in addition to the spin of  $1\hbar$  that a "normal" photon has). Such light beams are routinely generated in the visible employing table-top laser sources. The properties of 4<sup>th</sup> generation synchrotron sources, as they are in the first design stage nowadays, will allow for the first time the routine generation of these novel photon beams at high photon energies. A prerequisite for such photon beams is a high transverse coherence, which will be available at future diffraction-limited radiation sources. The photon-matter interaction using OAM photons is a completely untouched area of fundamental research. Possible vortex photon beams allow, for example, addressing completely different states than beams consisting of "normal" photons.

### Concluding Remarks

COLTRIMS coincidence experiments provide the most detailed and complete account of photon driven dynamics in the gas phase. They are the eyes into the many-body quantum world, which deliver pictures of unseen detail and completeness. The information content provided by these experiments is so rich that already today hypothesis-driven data-mining is as important as the experiment itself. It is the imagination during that process and the ingenuity of questions asked to the data which often drives the progress.

### Acknowledgments

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### 3.5 Circular dichroism in photoemission of chiral molecules: from static to time-resolved

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Yann Mairesse, CNRS-CELIA Talence,  
Olga Smirnova, Max Born Institute & Technical  
University Berlin

#### Status

Holding a sea snail in your hand during your beach vacations you may wonder about the beauty and the function of countless chiral objects, including the snail and the hand. Chiral objects cannot be superimposed onto their mirror-images; what is the purpose of this asymmetry in Nature? Our ability to shake hands is one of them. Chiral molecules, including elementary bricks of life such as amino-acids and nucleic acids sugars form chiral complexes, a “molecular handshake”, when embedded into a chiral medium. Thanks to this chiral recognition we can perceive odors and tastes and metabolize drugs. Discriminating enantiomers and understanding chiral interactions are major challenges in pharmacology, food and fragrance industry, chemistry, life sciences and physics. We spotlight here a recent disruptive step: the “electric dipole revolution” in chiral discrimination.

The chiral response originating from magnetic light-matter interactions is generally markedly weak. It can be boosted by increasing the light frequency – e.g. in X-ray absorption circular dichroism [1] – or the light intensity – in chiral high-order harmonic generation [2]. Both methods provide dichroisms in the few % range, surpassing standard absorption circular dichroism by orders of magnitude and enabling ultrafast time-resolved experiments.

A powerful alternative to increase chiral signals is to find observables producing enantiosensitive signals in the electric-dipole approximation: this is the “electric dipole revolution”. This revolution broke into chiral discrimination when several extremely efficient methods emerged, exciting and probing chiral dynamics in bound electronic, vibrational (photoexcitation circular dichroism PXCD) [3] and rotational states (enantio-sensitive microwave spectroscopy ESMS) [4] of molecules. Its herald, photoelectron circular dichroism (PECD), was predicted by Richie [5], quantified by Powis [6], and then observed by Bowering et al [7]. The revolutionary methods – PXCD and its detection via photoionization (PXECD), ESMS and PECD – exploit the same physical mechanism [8], exciting vectorial observables such as induced polarization (PXCD, ESMS) or photoelectron current (PECD, PXECD), and detecting them using a chiral reference frame fixed by the experimental set-up. The

“electric dipole revolution” deposed the magnetic field and substituted it by a chiral experimental arrangement: the directions fixed by the two detectors together with the plane fixed by the electric field constitute a chiral reference frame needed to measure vectorial observables and distinguish enantiomers (Instead of circularly polarized light, one can also define the plane using two orthogonally polarized electric fields, with different frequencies (ESMS)). Several 10’s of % [9] difference in signals from the two detectors in PECD highlights the extreme efficiency of dipole-based techniques (see Fig. 1).

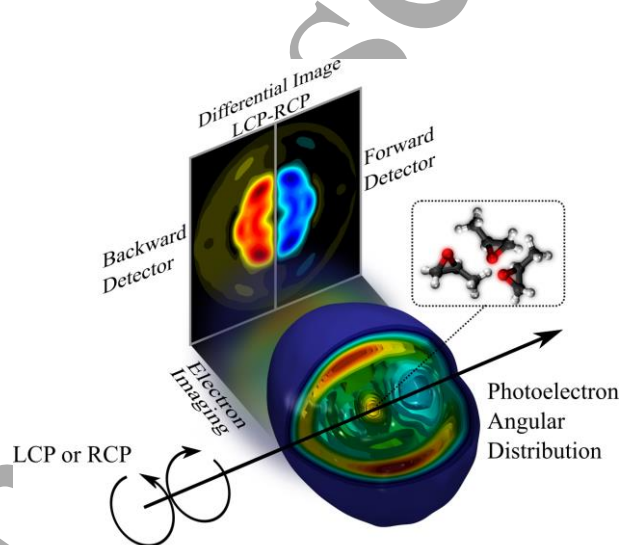


Figure 1: Principle of a PECD experiment. A gaseous sample of randomly oriented chiral molecules is ionized by Left or Right Circularly Polarized light, producing a forward/backward asymmetric electron angular distribution.

The combination of synchrotron-based studies and tremendous advances in modeling one-photon ionization of chiral molecules from valence and core states has underscored the basic features of PECD: its initial-orbital dependence, its complex dependence on photoelectron energy reflecting the scattering dynamics in the asymmetric molecular potential and therefore its sensitivity to static and dynamical molecular structures [10] such as isomers, conformers, clustering and vibrational motion [11]. The advent of table-top laser-based PECD [12,13] exploiting resonant multiphoton ionization coupled sometimes to ion momentum imaging in a coincident scheme enabled accurate determination of enantiomeric excess [14], time-resolving ultra-fast chiral molecular dynamics [3,15,16] and revealed a universal nature of PECD from weak to strong field ionization regimes [17].



### Current and future challenges

The electric dipole revolution introduces highly efficient methods but also poses new challenges. We leave the rich and exciting field of ESMS outside the scope of this paper.

*Enantio-sensitive spectroscopy in complex systems: biomolecules, nanomaterials, liquids.*

Many chiral systems, and especially biomolecules, are floppy *i.e.* can be found in several conformations. While the PECD signal for each conformer may be different and even have an opposite sign, isolating individual signals is challenging. For example, a hard-to-assess distribution of conformers in molecular beams obscures interpretation of measurements. *Designing new schemes for efficient selection or tagging of conformers* is thus an acute challenge. One route is to select conformers using electrostatic selectors. Alternatively, individual conformers can be tagged spectroscopically using high resolution in a ns-REMPI scheme.

So far, most PECD studies focused on easy-to-handle systems (terpenes and oxirane derivatives). *Bringing biomolecules such as amino-acids in the gas phase as intact species is a challenge.* A successful example is detecting PECD in alanine [19], suggesting PECD as a possible origin of the homochirality of life operative in interstellar medium [9]. Can we use PECD to study how elementary bricks of life form larger biopolymers such as peptides? Could supramolecular chirality, such as alpha-helix, be studied by PECD?

Chiral clusters and nanomaterials are a playground for exploring self-organization in chiral systems. The evolution of PECD from monomers to clusters to nanoparticles (NPs) could monitor the competition between order and disorder, e.g. through the decrease of PECD due to inelastic electron scattering vs its enhancement due to micro-crystallization. Functionalized nanoparticles, e.g. metallic NPs capped with chiral species shells, could also be investigated by PECD, which should be sensitive to possible chiral induction from the ligand-shell to the interface, or even to the metallic NP substrate. These individual contributions might be disentangled thanks to the different electron binding energies.

Finally, solvation processes in chiral systems from nanosolvation in a molecular beam towards full solvation in a liquid jet could be efficiently probed by PECD, with direct relevance to *in vivo* conditions.

*Developing ultrafast imaging of chiral dynamics and learning to “read” the images*

Ultrafast imaging of chiral dynamics has begun. Time-dependent PECD [15], PXECD [3], chiral high harmonic generation (cHHGm) [2], phase-resolved PECD giving access to the forward/backward

asymmetry in the photoionization delays [20], could open a view into asymmetric chemical reactions, the coupling between chiral electronic dynamics and torsion modes capable of changing molecular handedness, the role of decoherence in stabilizing transient chirality, etc. Bringing the “electric dipole revolution” to high harmonic generation (cHHGd) [21] and other nonlinear optical spectroscopies will broaden the scope of accessible dynamics. The extension of the first time-resolved measurements (see Fig. 2) to more complex systems is an ultimate goal. As in any nonlinear spectroscopy, the challenge is to correlate measured signals with specific types of dynamics. The short-term goal is to perform such experiments on smaller, well characterized systems, to determine the sensitivity of every observable to different classes of dynamics. The excitation of vibrational wavepackets by stimulated impulsive Raman scattering could reveal the importance of different molecular sites, associated with different vibrational modes, in the probe signal (PECD, cHHG or phase-resolved PECD). Pump-probe measurements of photodissociation dynamics leading to achiral species could provide a dynamical picture of the death of chirality. PECD will be a natural observable of ultrafast isomerization and conformational changes, whose measurement is challenging. Increasing the excitation photon energy to the XUV and soft-X-ray range will trigger multi-electron relaxation dynamics driven by electronic correlation and would as well provide site and chemical-specificity. Time-resolved measurements could reveal the transfer of the forward/backward asymmetry through correlation. Pre-alignment of the molecular samples by laser pulses – challenging for large molecules with low symmetry – could enhance the asymmetries, offering new information [22].

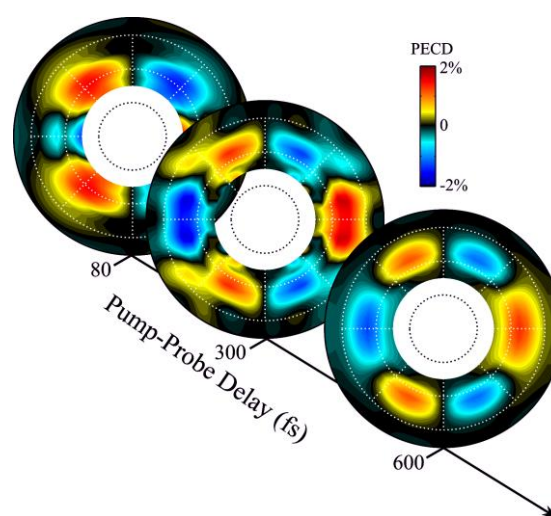


Figure 2: Time-resolved PECD measurement in fenchone, photoexcited by absorption of two linearly polarized photons of 3.1 eV energy and probed by ionization by two circularly polarized 1.55 eV photons [16].

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5 “Reading” ultrafast images requires theory at all levels  
6 of complexity: from simple models [17,18] to  
7 sophisticated simulations [23,24]. These are challenging.  
8 A quantitative description of photoionization may  
9 require inclusion of two-hole excitations, fundamentally  
10 difficult in the computationally efficient density  
11 functional theory while computationally demanding for  
12 bio-molecules in ab-initio methods. Compounding the  
13 problem is the coupling between electronic and nuclear  
14 motions, evolving on the highly multi-dimensional  
15 potential surfaces. Interpreting these dynamics will  
16 necessitate answering fundamental questions: What  
17 determines the sign of PECD? How does it depend on  
18 the photoionization mechanism: one-photon, REMPI, or  
19 strong-field [14], and on the localization of the initial  
20 orbital vs the chiral center(s)? What is the role of the  
21 photoionization propensity rules, different in different  
22 photoionization regimes?

#### 23 *Adapting the “electric dipole revolution” to practical* 24 *applications*

25 Both PECD and cHHG [21] can be used to perform  
26 accurate determination of enantiomeric excess. In  
27 PECD accuracies below 1% can be reached, now with  
28 laser-based setups [14]. Coincidence electron-ion  
29 detection enables analysis of multicomponent samples,  
30 with high potential for industrial applications  
31 (<http://www.massspecpecd.com>). Challenges include  
32 reducing measurement times from several hours to  
33 minutes or seconds to follow the evolution of chemical  
34 samples in real time, and new methods for disentangling  
35 the contribution of multiple species with or without (in  
36 the case of isomers) coincidence detection.

#### 37 **Advances in science and technology to meet the** 38 **challenges**

39 Compact high-power fiber-laser systems are  
40 revolutionizing ultrafast science by enabling  
41 femtosecond and attosecond pulses to be generated with  
42 unprecedented repetition rates and photon fluxes, from  
43 the visible to the XUV. A key challenge is generating  
44 XUV light with circular polarization. One route is  
45 offered by free electron lasers, but their repetition rate  
46 should be increased to allow time-resolved coincidence  
47 imaging. Another possibility is to drive high-harmonic  
48 generation with tailored light fields or to use resonances,  
49 which allow one to tune the harmonic ellipticity. From  
50 the photoexcitation point of view, broadband ultrashort  
51 UV sources will also have to be developed.

52 Advances are needed to bring fragile biomolecules as  
53 neutrals in the gas phase, e.g. by aerosol  
54 thermodesorption or laser-desorption, ensuring  
55 compatibility with imaging techniques and high  
56 repetition rate light sources. Size-selected large clusters

and large biopolymers could be produced as ions by  
electrospray techniques, where high densities should be  
achieved while coupling with photoelectron  
spectrometers, ideally in coincidence with ion detection.

Coincidence electron-ion imaging spectrometers will  
benefit from the repetition rate boost of fiber lasers, but  
the electronics and detection schemes may have to be  
improved to admit highest count rates.

Modeling ultrafast chiral response requires the  
description of coupled electronic-nuclear dynamics  
together with photoionization. The potential energy  
surfaces governing the dynamics could be computed  
“on-the-fly”, as demanded by nuclear trajectories  
spawned into this multi-dimensional space [24].  
Algorithms winnowing out trajectories to speed up  
computation [25] would be a game changer. Coherent  
electronic- nuclear wavepackets are characteristic of  
PXCD [3], with signals sensitive to their coherence.  
Here, fully quantum description of electrons and nuclei  
may benefit from looking at the problem as that of an  
open quantum system, sorting degrees of freedom into  
“system” and “bath”, and treating the bath accordingly.

#### 57 **Concluding remarks**

58 Photoionization processes in chiral systems involve new  
59 chiroptical observables, sensitive to molecular  
60 structures and ideal to template molecular dynamics. A  
joint development of theoretical and experimental  
methods will allow the extension and understanding of  
chiral interactions from model chiral molecules up to  
much more complex chiral systems as encountered in  
biology and nanosciences.

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## 4 COMPLEX SYSTEMS

### 4.1 Fragmentation of gas-phase biomolecules – Thomas Schlathöler, University of Groningen, Zernike Institute for Advanced Materials

#### Status

Almost 20 years ago, a large number of groups from the ICPEAC community embarked on a journey destined to a fundamental understanding of the molecular mechanisms underlying radiation therapy. In particular, interactions of low-energy electrons [1], keV ions [2], highly charged ions (HCI) [3] and photons [4] with gas-phase DNA building blocks or their model systems were studied. It was soon established that nucleobases are much more resilient towards particle-collisions than deoxyribose molecules. Energetic fragments were observed particularly for HCI impact and these ions were found to potentially induce severe secondary DNA damage.

From an experimental perspective, the appeal of studying molecules such as nucleobases and deoxyribose was the straightforwardness of preparing a sufficiently dense gas-phase target by means of evaporation at modest temperatures. From a theoretical point of view, the level of complexity, i.e. the number of constituent atoms in a nucleobase was sufficiently small to allow high-level quantum-mechanical treatment.

The studies on DNA building blocks delivered a wealth of important data such as cross sections (which are input in track structure calculations) and greatly improved our understanding of dynamics in strongly perturbed polyatomic molecules. However, a major fundamental question persisted: Are collision-induced dynamics in small gas-phase DNA building blocks representative for the molecular processes occurring in living cells, where radiation interacts with entire double stranded DNA strands that have complex tertiary and quaternary structures defined by DNA interactions with water and proteins?

To answer this question while at the same time sticking to the strategy of evaporating DNA building blocks, the Caen-Groningen collaboration studied keV HCI collisions with clusters of nucleobases [5]. In these clusters, nucleobases form a hydrogen-bound network. As a consequence of the intermolecular hydrogen-bonds, a nucleobase in a cluster exhibits fragmentation channels absent in a single gas-phase molecule, proving the critical importance of a chemical environment for the molecular response.

For the nucleoside thymidine, Maclot *et al.* [6] have compared fragmentation by HCI collisions with state-selective photoelectron-photoion coincidence data. This

data served as a basis to precisely determine the excitation energy distribution of thymidine induced by HCI collisions.

A limitation of preparing molecular targets by evaporation is the thermal stability of the molecules. It is impossible to evaporate intact DNA strands in useful amounts. Even for DNA building blocks such as deoxyribose, evaporation can lead to coexistence of structurally conformers, which can respond to collisions in very different ways.

A much more gentle alternative technique was explored by the Aarhus group, who used electrospray ionization (ESI) to produce beams of complex biomolecular ions and studied their collisions with neutral atoms. In collaboration with the Stockholm and Caen groups, they were able to study collision induced dissociation of the nanosolvated intact nucleotide adenosine monophosphate (AMP) anion [7]. The isolated AMP anion was found to break preferably between deoxyribose and nucleobase. With increasing number of attached water molecules  $m$ , fragmentation is quenched and AMP is fully protected when  $m > 13$ .

#### Current and Future Challenges

ESI has become a key technique for preparation of biomolecular targets for collision studies, as it allows for bringing biomolecular systems of virtually unlimited complexity into the gas phase. Our group in Groningen and the SOLEIL group in parallel tackled the problem of low target density by collecting the electrosprayed ions in a radiofrequency (RF) ion trap which can be interfaced with synchrotron or keV/MeV ion beamlines.

Pioneering studies on photon and keV ion induced ionization and fragmentation of the gas-phase protonated oligonucleotide (dGCAT+H)<sup>3+</sup> confirmed the important role of glycosidic bond cleavage leading to base loss and also the fragility of the deoxyribose unit. However, the actual fragmentation patterns are entirely different from what is observed in isolated DNA building blocks, stressing the relevance of studying such large systems [8].

With establishment of ESI and related methods, the size of the system became virtually unlimited as with ESI even an entire virus can be brought into the gas-phase. At present, the influence of DNA tertiary structure on radiation damage is explored. Milosavljevic *et al.* were able to determine ionization energies of various gas-phase protonated proteins as a function of protonation state and could link these data to map structural transitions from folded to unfolded states [9]. Applications of the method include proteomics, biomolecular radiation damage and drug-receptor interactions but also astrochemical problems or studies on magnetic molecules.

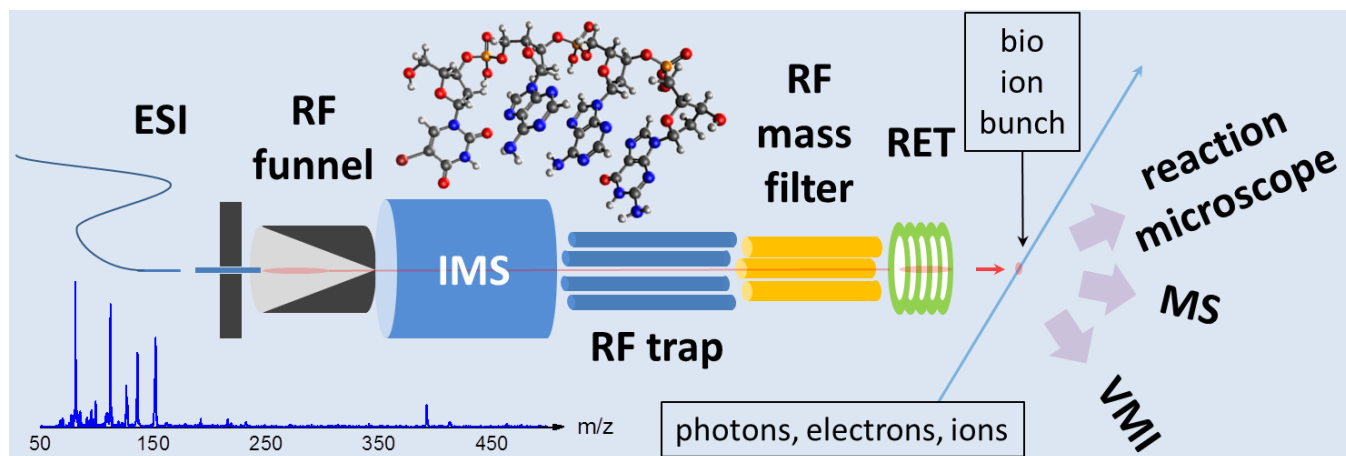


Figure 1: Sketch of a future crossed-beam type experiment.

Combination of ESI with ion mobility spectrometry, a technique sensitive to the molecular geometric cross section, would even allow for production of conformationally pure targets. It is clear that for instance the application of this experimental strategy to biopolymers (e.g. polynucleotides and polypeptides) will open up a virtually unlimited library of tailored target systems, to investigate charge migration, proton migration, isomerization or conformational changes in general, but also specific problems from the field of biomolecular radiation damage. Not only sequence and length of the biopolymers can be arbitrarily chosen, but the polymers can be chemically modified, for instance by incorporating halogenated nucleobases where the halogen atom can be selectively core excited/ionized.

A next breakthrough step will be towards real crossed-beam experiments (a first study on keV ion collisions with a protonated peptide have been published very recently [10]). For collisions in RF ion traps, many photoproduct properties such as momenta and correlations are lost and information on emitted electrons is fully unavailable. However, with sufficiently intense continuous ESI beams or, or with bunched beams extracted from RF traps, it will still be possible to generate “free” targets. Only then, it will be possible to still use state-of-the-art experimental techniques for biomolecular ion manipulation but then in combination with the entire arsenal of molecular physics techniques and in particular advanced coincidence spectroscopies (PEPIPICO-schemes, velocity map imaging, reaction microscopes etc.)

The appeal of gas-phase experiments is not only the possibility of investigating collision induced dynamics in the absence of a disturbing chemical environment but also the possibility to study for instance nanosolvation. Up to now, only very few groups succeeded in attaching water molecules to gas-phase biomolecular systems. An experimental technique that allows for controlled addition of an arbitrary number of H<sub>2</sub>O molecules is in high demand.

Last but not least, the step towards real time-resolved measurements has to be done. We need to find answers to the following questions: How does energy deposited in a complex biomolecular system develops in space and time? How – and how fast – does this electronic excitation couple to nuclear degrees of freedom and ultimately to conformational changes and/or fragmentation? How can we use this knowledge to solve problems of technological and societal relevance? The pioneering work of Calegari et al. [11] has shown the feasibility of tracking attosecond charge migration in the amino acid phenylalanine, a relatively simple biomolecular system. For investigating ultrafast processes in much larger biomolecular systems with sufficient detail, the availability of “free” targets extracted for instance from ESI sources (discussed above) is also a requirement

It is clear that the current experimental progress will need corresponding advances on the theory side. Complex biomolecular systems are far beyond reach for ab-initio quantum calculations and must be tackled by molecular dynamics type approaches. The combination of a quantum treatment of the initial excitation/ionization and a molecular dynamics treatment of the subsequent dynamics (multiscale modelling) is an ongoing challenge for theory.

### Advances in Science and Technology to Meet Challenges

The field is currently benefitting from a large number of technological developments, a few of which are particularly interesting. Electrostatic ion traps and storage rings are ideal tools for studies on dilute targets of extremely heavy molecular systems that are delivered by ESI sources. The double electrostatic storage ring DESIREE [12] allows for investigation of meV ion-ion interactions.

For RF guiding, trapping and manipulation of biomolecular ions, more and more printed-circuit-board (PCB) devices such as ion funnels, RF guides and



traveling-wave ion-mobility spectrometers are being developed PCB-based devices for production of complex biomolecular ions and their manipulation and analysis are compact, versatile and cost-effective and will allow for easy interfacing of with existing experimental setup.

Last but not least, we are currently witnessing enormous developments in the field of ultrafast pulses of energetic photons. With the advent of X-ray free electron laser facilities and attosecond pulse lasers new tools for biomolecular fragmentation studies are currently becoming available and even direct diffraction imaging of biomolecular processes are now feasible.

### Concluding Remarks

Fragmentation studies on gas-phase biomolecules have long focused on small biomolecules such as amino acids and nucleobases. These systems were considered extremely complex from a physics point of view and overly simplified for application of the obtained results in medicine, biology and material sciences. The field has now reached a turning point, where experimental developments already allow for investigation of very large gas-phase biomolecular systems such as DNA complexes or proteins of direct biological relevance. The experimental challenge for the near future is time-resolved experiments on ultrafast electronic and structural dynamics in gas-phase DNA and polypeptides. From a theoretical point of view, it will be most challenging to implement the initial excitation/ionization in a quantum mechanically correct way and to follow the subsequent dynamics of the large biomolecular systems by molecular dynamics type approaches.

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Accepted manuscript

## 4.2 Fullerene Spectroscopy and Dynamics -

Eleanor E.B. Campbell, University of Edinburgh

### Status

Since the discovery of  $C_{60}$  by Kroto, Smalley and co-workers in 1985, the molecule has been shown to be a fruitful and interesting model system for exploring the properties of complex molecules. Due to its high symmetry and ease of handling,  $C_{60}$  has often been the molecule of choice for the atomic and molecular physics community to explore molecular systems with many degrees of freedom. Fullerene studies have shown the existence of a range of phenomena and motivated development of theoretical tools that have subsequently been applied to other complex molecules and clusters.

Early studies of ns laser photo-excitation/ionization showed behavior that can be considered to be highly statistical in nature. This involved thermionic photoelectron emission, metastable  $C_2$  emission from excited cations and radiative decay. The unusual properties of  $C_{60}$  and other fullerenes, i.e. their relatively low ionization energies and high energetic thresholds for fragmentation, made it possible to observe all of these decay processes on the timescale of conventional mass spectrometers. These studies have been well covered in a number of review articles and statistical models have been applied to explain the experimental results (see e.g. [1, 2] and references therein). A statistical interpretation in terms of efficient energy equilibration followed by thermionic electron emission explained the early photoionization studies, involving multiphoton absorption of pulses in the visible and UV regime. This has been validated by more recent electron imaging experiments at the IR free electron laser FELICE where the vibrational degrees of freedom were directly excited [2]. Single-photon ionization experiments in the VUV, using synchrotron radiation, showed excitation of a plasmon resonance in  $C_{60}$  at ca. 20 eV and interesting oscillations in the HOMO/HOMO-1 photoionization cross-sections that could be related to the electron density on the carbon cage [2]. Time-dependent DFT calculations can reproduce the main features of the ionization dynamics [3]. The influence of encapsulated metal atoms on the single-photoionization cross sections has also been investigated with evidence for significant redistribution of metal electron oscillator strength [4]. Photoelectron imaging studies have provided surprising evidence for the occurrence of low energy thermoelectronic emission (thermal emission from a hot electron bath prior to coupling with vibrational degrees of freedom) following single-photon VUV absorption in the 27 – 65 eV energy range [5].

A similar thermoelectronic emission was observed in multiphoton excitation experiments using femtosecond

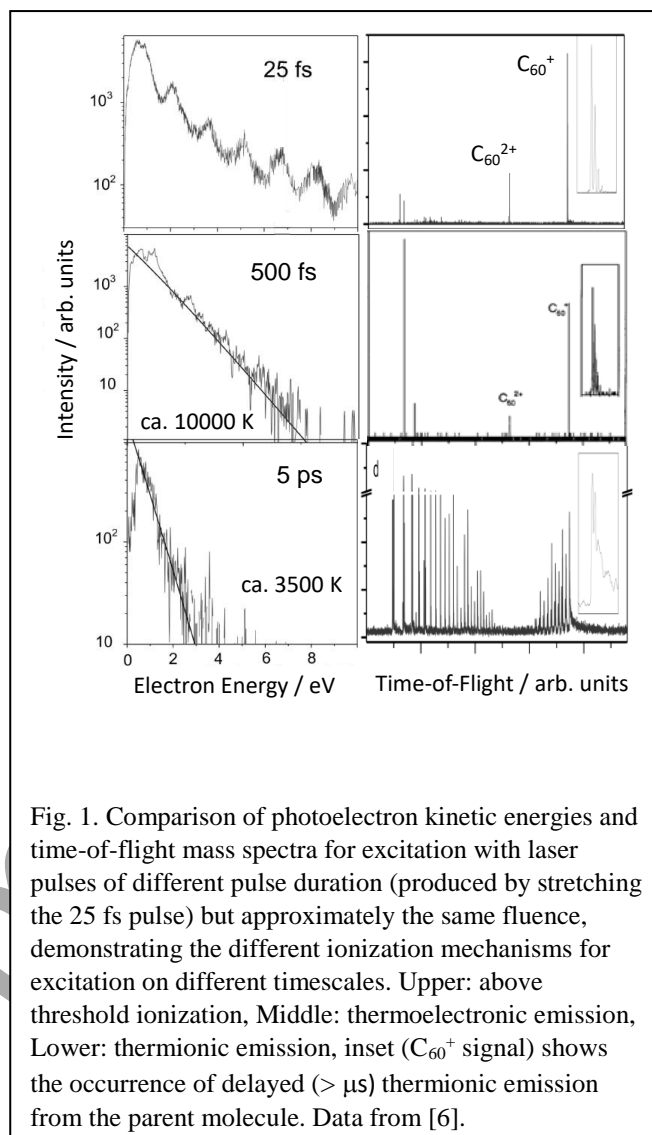


Fig. 1. Comparison of photoelectron kinetic energies and time-of-flight mass spectra for excitation with laser pulses of different pulse duration (produced by stretching the 25 fs pulse) but approximately the same fluence, demonstrating the different ionization mechanisms for excitation on different timescales. Upper: above threshold ionization, Middle: thermoelectronic emission, Lower: thermionic emission, inset ( $C_{60}^+$  signal) shows the occurrence of delayed ( $> \mu s$ ) thermionic emission from the parent molecule. Data from [6].

laser pulses. Three different regimes can be identified for multiphoton ionization of  $C_{60}$ , [6] Fig. 1: above threshold ionization for pulse durations,  $\tau < 30$  fs [2]; thermoelectronic emission from hot electrons for  $50\text{fs} < \tau < \text{ps}$ , both regimes characterized by electron emission from vibrationally cold parent ions and the production of multiply-charged parent species; and thermionic electron emission accompanied by extensive fragmentation and no multiple ionization, for timescales longer than the electronic-vibrational coupling timescale (a few hundred femtoseconds). Photoionization studies on empty and endohedral fullerenes using ultrashort x-ray pulses, accompanied by MD simulations, show a high degree of multiple ionization accompanied by statistical fragmentation, similar to results obtained for intense fs laser pulses in the UV and visible spectral region [7]. Attosecond studies that can probe the photoemission delay are providing detailed insight into the role of the plasmon resonance on the electron dynamics [2].

The efficient coupling between electronic and vibrational degrees of freedom that leads to a rapid

1  
2  
3 equilibration of excitation energy has the consequence  
4 that it is extremely difficult to apply conventional  
5 multiphoton ionization spectroscopy techniques to  
6 obtain high resolution spectroscopic information about  
7 the properties of excited electronic states of gas phase  
8 fullerenes. Only limited resonant two-photon ionization  
9 data on cold gas phase molecules exists and although  
10 single-photon absorption measurements in rare gas  
11 matrices or He droplets have been performed, the  
12 influence of the matrix is uncertain [8]. Angle-resolved  
13 photoelectron spectroscopy using fs laser excitation,  
14 that relies on the efficient energy coupling and does not  
15 require vibrationally cold molecules, has been used to  
16 obtain information about diffuse low-lying Rydberg  
17 (“Super-Atom Molecular Orbital”) states of fullerenes  
18 and endohedral fullerenes [9]. The high resolution  
19 absorption and emission spectroscopy of fullerenes in  
20 the uv-IR has become highly topical with the discovery  
21 of relatively large amounts of  $C_{60}$  and  $C_{60}^+$  in space. The  
22 confirmation of  $C_{60}^+$  as the first unambiguously  
23 identified source of diffuse interstellar bands was  
24 recently made possible by high resolution He-tagging  
25 spectroscopy studies of cryogenically cooled He- $C_{60}^+$   
26 [8].

### 27 28 29 **Current and Future Challenges**

30  
31 A major current challenge is to produce sufficiently  
32 dense molecular beams of isolated fullerenes with well  
33 defined levels of internal excitation energy (and  
34 specifically vibrationally and rotationally cold  
35 molecules). This is important for determining the  
36 absorption spectra in the UV to near-IR range for  
37 comparison with diffuse interstellar bands, for  
38 investigating the excitation mechanisms leading to IR  
39 emission and for determining reaction rate constants of  
40 relevance for astrochemistry. There are currently no  
41 high resolution electronic spectra of neutral fullerenes  
42 that have sufficient accuracy to provide a meaningful  
43 comparison with astronomical observations. Not only  
44  $C_{60}$  is relevant here, also higher fullerenes,  
45 hydrogenated fullerenes and possibly some endohedral  
46 fullerenes are expected to be important constituents of  
47 the interstellar medium and play a role in astrochemistry  
48 [10]. There is also the important challenge of developing  
49 accurate, predictive theories for such large systems that  
50 can provide aid to the spectroscopists in identifying the  
51 most likely candidates for detection in space and also for  
52 providing input for astrochemical models.

53  
54 Sources of gas phase molecules with well-defined (and  
55 controllable) internal energy content will also be  
56 valuable for exploring and understanding the complex  
57 dynamics induced by excitation with intense, ultrashort  
58 laser pulses in experiments that are beginning to become  
59 possible (see e.g. section 2.4). This would enable the  
60 role of laser-induced nuclear dynamics to be more  
clearly decoupled from that of thermally excited

vibrational modes of large, complex molecules and  
provide a more accurate comparison with detailed  
theoretical calculations. Endohedral fullerenes can be a  
particularly fruitful category for studying the effects of  
confinement on the spectral properties of atoms/ions or  
small molecules and the timescales for energy transfer  
and equilibration.

Fullerene experiments are often justified on the basis of  
using them as model systems to understand more  
complex molecules of e.g. biological significance. This  
may be appropriate but advances in the ability to  
produce gas phase targets of such molecules combined  
with complex multi-coincidence measurement  
techniques may soon negate that argument. Instead one  
should not be afraid to justify experiments on the basis  
of fundamental curiosity-driven science. As in the past,  
this is likely to lead to important new insights that may  
be even more far-reaching.

### **Advances in Science and Technology to Meet Challenges**

In the author’s opinion, the challenge of understanding  
the formation and role of fullerenes in space is the most  
exciting area that can be expected to develop in the  
coming years. This requires the availability of well-  
characterised molecular targets (both charged and  
neutral) and the application of new experimental  
facilities that are beginning to become available. In  
particular, the cryogenic storage rings at Stockholm  
University (DESIREE), MPI Heidelberg (CSR) and  
RIKEN (RICE) as well as other facilities to be built or  
upgraded to cryogenic operation in Leuven and Lyon,  
will enable experiments under well-controlled extreme  
conditions that are relevant for astrochemical  
conditions. In combination with laser excitation and  
photon detection, many novel experiments can be  
envisaged that will provide new information of direct  
relevance for exploring and modelling fullerene  
astrochemistry. Laboratory advances will be matched  
by advances in observational possibilities. The James  
Webb Space Telescope, due to be launched in 2021, will  
provide unprecedented sensitivity, spatial and spectral  
resolution in the near to mid-IR part of the spectrum and  
will provide a wealth of data to provide clues to the  
formation of fullerenes in space and the possible role of  
other fullerene species. The laboratory and  
observational advances will also need to be  
accompanied by advances in the efficiency and accuracy  
of predictive theory and advanced data mining  
techniques to extract the relevant signals from the vast  
amounts of observational data that are to be expected.

### **Concluding Remarks**

Fullerenes have been convenient and extremely fruitful  
model systems for understanding the fundamental  
electron dynamics of complex molecules over the past

25 years. With the current advances in extreme light facilities and detection techniques this role is likely to continue into the foreseeable future. In addition, the more conventional absorption/emission spectroscopy and reactions of fullerene species under conditions relevant for astrochemistry is likely to be an important area of research in coming years.

**Acknowledgments** – The author gratefully acknowledges the contributions of many collaborators and students over the past 30 years.

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### 4.3 Clusters interacting with attosecond and FEL pulses - Jan M. Rost, Max Planck Institute for the Physics of Complex Systems, Dresden

#### Status

Clusters consist typically of atoms or small molecules with a density of more than  $10^{22}$  /cm<sup>3</sup>, i.e., close to solids. Yet, clusters can be handled like individual atoms or molecules in experiments since they are available in beams following supersonic expansion from a gas jet. This property has rendered them ideal targets for probing time-gated high-density interaction of light and matter. In other words, clusters absorb or scatter a large number of photons within the fraction of the (short) duration of the light pulse. This leads inevitably to serious radiation damage, and in most cases, Coulomb-explosion, whereby electrons and ions are ejected with partially high velocities. Therefore, fast transport of ever new targets via the cluster beam through the overlap region with the laser pulses is a valuable and actually unique property of clusters among high-density targets to achieve good statistics in experiments.

The first spectacular result with FEL light at FLASH in Hamburg (at that time still called “TESLA test facility”) was the unexpected large number of photons absorbed by Xenon clusters, eventually explained by inverse Bremsstrahlung through ions, higher charged in a transient nanoplasma than concluded from the spectrum of the (recombined) ions [1]. Nanoplasma formation is the result of the pertinent two ionization mechanisms in clusters, namely *outer ionization* for electrons leaving the cluster, and *inner ionization* for electrons which remain trapped by the large positive background charge of the cluster and form a transient plasma.

Apart from rare gas clusters, special cluster forms such as fullerenes [see section by N. Berrah] and helium droplets [2] have become popular targets. The former are extremely simplified models for large (bio-) molecules, while the latter, in particular if doped, exhibit intriguing dynamics and attain superfluid properties if large enough, giving rise to vortex lines which were made visible in an XFEL experiment [3]. Both of them exhibit collective behavior of the transient nanoplasma [4] whose formation is a general signature of impulsive excitation of clusters [5]. Of pertinent interest are also composite clusters. These can be core-shell systems, e.g., with a heavier rare gas species (xenon) embedded into an argon or helium hull. Protonated systems, e.g., clusters formed with methane molecules, get rid of the absorbed high-frequency photon energy by ejecting protons reducing the plasma temperature [6].

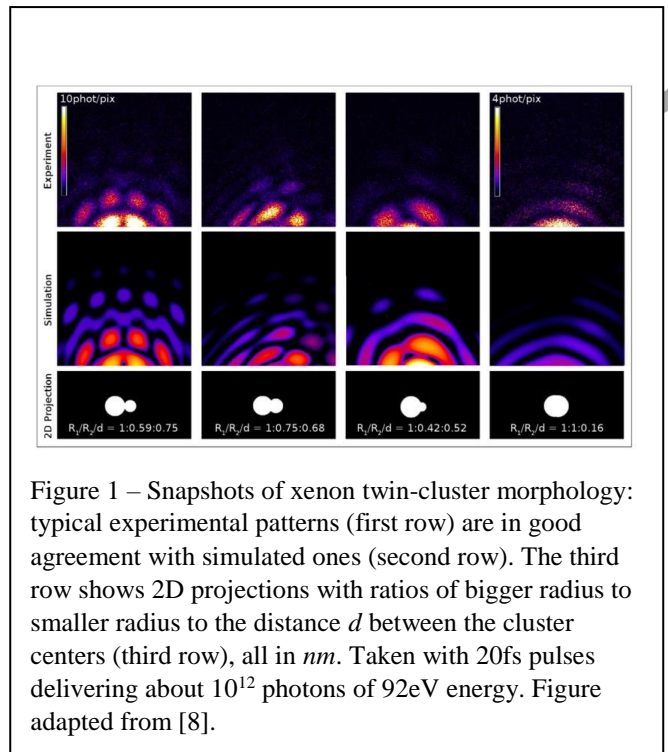


Figure 1 – Snapshots of xenon twin-cluster morphology: typical experimental patterns (first row) are in good agreement with simulated ones (second row). The third row shows 2D projections with ratios of bigger radius to smaller radius to the distance  $d$  between the cluster centers (third row), all in  $nm$ . Taken with 20fs pulses delivering about  $10^{12}$  photons of 92eV energy. Figure adapted from [8].

Clusters also build a bridge to nanoparticles and much larger structures such as viruses [7].

The probably most serious disadvantage of cluster sources is their loose definition as a target since they are not mass selected but are generated in the supersonic beam with a statistical (log-normal) size distribution about a mean size of  $N$  constituents. In addition, traditional intense pulse experiments suffer from a statistical (typically Gaussian) intensity distribution about the laser focus since one does not know where the target, still much smaller than the focus volume, was hit.

Both problems have been overcome simultaneously in recent XFEL experiments by imaging the cluster along with recording traditional parameters such as ions and electrons and eventually fragments [8]. This permits to determine cluster size and light intensity of interaction events in individual laser shots “single-shot experiments” a posteriori. One surprise in such experiments was the finding that clusters often are not “clean” targets but occur as twin systems, with a smaller and a larger cluster forming one system, see Figure 1.

#### Current and Future Challenges

Clearly, imaging objects on the nanoscale or even smaller is an extremely powerful technique which is only at its infancy. Quests go for higher spatial and temporal resolution. Unconventional proposals such as holographic projection with the help of a cluster as a

1  
2  
3 reference particle in the target (in this case a virus) have  
4 already been realized [7]. Also, attention has been given  
5 recently to the recombination dynamics in clusters, that  
6 is, to the long-term fate of inner-ionized electrons [9].  
7 Recombination dynamics is also an important feature on  
8 the way to understand the transition to equilibrium.

9  
10 In general, XFEL or attosecond pulse experiments  
11 follow a time-dependent perspective in a pump-probe  
12 scheme, often with a combination of an infrared and a  
13 high-frequency pulse or of two high-frequency pulses,  
14 slightly detuned in frequency, such that the first and the  
15 second pulse can be distinguished. For small systems  
16 such experiments extract the same information as time-  
17 independent measurements as long as the coherence or  
18 full information of the system is preserved (in this case  
19 time-resolved and energy resolved observables are  
20 roughly speaking related by Fourier transform).

21  
22 Conversely, unique information can be retrieved by  
23 time-resolved experiments if the target exhibits transient  
24 dynamical behavior whose signature is lost once  
25 particles have arrived at the detector, where they provide  
26 the information for measurement without time-  
27 resolution. One example is the transient creation of  
28 highly charged ions which are responsible for the large  
29 photo absorption observed in the early FEL experiments  
30 with Xe clusters [1] as discussed above. Other proposed  
31 experiments include the ultrafast but transient  
32 nanoplasma response to an attosecond pulse excitation  
33 [9,10] or more systematically the monitoring of the non-  
34 equilibrium dynamics of complex systems during the  
35 excitation and in the early times of relaxation,  
36 respectively dissipation, of the excitation energy.

37  
38 Such questions parallel the interest of condensed matter  
39 physics in the fundamental question, to which extent,  
40 how and on which time scale, finite quantum systems  
41 relax to equilibrium.

42  
43 In the context to ultrashort pulses and in particular for  
44 clusters an interesting question is also if relaxation (to  
45 equilibrium) is always governed by an almost  
46 homogenous (and slow) dissipation of energy into many  
47 channels, respectively degrees of freedom. Can it  
48 happen that dissipation of electronic energy happens  
49 suddenly and localized, e.g. through a specific transfer  
50 of energy to particular ions? This would correspond to  
51 true charge migration, eventually on ultrafast time  
52 scales, while electron dynamics following short pulse  
53 excitation without irreversible relaxation is more a  
54 coherent beating of electron density of all eigenstates in  
55 the energy interval covered by the energy width of a  
56 short pulse.

57  
58 In the violent excitation of clusters so far addressed the  
59 energy intake is so large that the systems behave to a  
60 large extent classically. Moreover, rare gas and metallic  
clusters become very similar in producing a

nanoplasma. However, much more gentle activation of  
non-equilibrium dynamics would preserve initially  
coherence and would permit to investigate the genuine  
quantum effects of dephasing and decoherence with  
ultrashort pulses in clusters representing complex but  
finite quantum systems. First steps towards this goal are  
described in [5].

## Advances in Science and Technology to Meet Challenges

Advances to get a handle on non-equilibrium dynamics  
and decoherence rely on progress in the theoretical  
description to propose experimentally accessible  
observables and their interpretation. Since we deal with  
truly time-dependent and excited complex systems,  
where many degrees of freedom are activated through  
the light pulse, this is already a challenge. What makes  
the problem even more difficult is the absence of the  
thermodynamic limit or other small parameter regimes  
in finite systems even if it becomes computationally  
feasible to describe such dynamics one needs to develop  
relevant simplifications and approximations for two  
reasons: First, to really gain insight which always is  
linked to meaningful simplification and therefore  
generalization, and secondly, to have a basis for  
meaningful coarse-grained formulations for even larger,  
also biologically relevant systems, approaching the  
microscale. This makes clusters also for the future  
fascinating but challenging objects.

Experimentally, the control over the cluster parameters  
has to and will be steadily improved. Sometimes  
progress is almost disruptive, as through imaging, e.g.,  
which suddenly has opened the possibility to gain  
information about cluster size and incident light  
intensity in an unexpected fashion [7,8,12]. A better  
characterization of single XFEL pulses and better  
statistics through high repetition rates is also desirable.  
With seeded beams, already available at ELECTRA in  
Triesete for XUV photons, and, e.g., LCLS II geared  
towards high pulse rates, these problems are addressed  
with next generation concepts. Moreover, novel analysis  
tools aided by machine learning certainly promise a  
positive impact [13].

## Concluding Remarks

Investigation of structure and properties of clusters, and  
in particular metal clusters, began some 35 years ago  
often by nuclear physicists. Ever since then clusters  
have served as paradigmatic systems to develop and  
explore experimental as well as theoretical techniques  
to answer pressing questions arising also in other, even  
more complex systems. Today, clusters do not only  
bridge the gap towards smaller systems such as

molecules but also towards larger systems at the nano- and biologically relevant micrometer scale. Ultrashort intense light pulses have added another chapter to this success story enabling a time-dependent perspective. In the future this combination is expected to provide unique insight into non-equilibrium dynamics of finite systems.

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## 5 NEW LIGHT SOURCES FOR NEW CHALLENGES

### 5.1 X-ray free-electron laser facilities – Michael Meyer, European XFEL

#### Status

Free-electron lasers (FELs) operating in the short-wavelength regime, i.e. at photon energies up to several tens of keV, have strongly advanced and revolutionized many research areas related to photon-matter interaction. Built up from long linear accelerators for ramping up the energy of electron bunches to values of several GeVs and long undulator sections to enable self-amplified stimulated emission (SASE), these large scale facilities are characterized by a dramatic increase (factor of  $10^8$ ) in average and peak brightness compared to other sources, such as synchrotron radiation facilities (see figure 1). Highly coherent x-ray pulses with duration in the femtosecond regime and pulse energies up to several mJ can be produced. Intensities up to  $10^{18}$  W/cm<sup>2</sup> are available for scientific applications in dedicated experimental areas. These unique performances open up numerous novel research opportunities. Non-linear multi-photon processes can now be driven at short wavelengths, i.e. favoring the interaction with core electrons; time-resolved studies of ultra-fast processes are enabled with simultaneous site-specific monitoring; and single-particle imaging of clusters, nano-particles and biomolecules can be realized due to the high number (exceeding  $10^{14}$ ) of photons in a single X-ray pulse.

First user operation of a short-wavelength FEL started in 2005 with FLASH at DESY in Hamburg, Germany [1] with lasing in the VUV-XUV regime. The SCSS test facility at Spring-8 in Japan and FERMI at ELETTRA in Italy were following in the softer photon energy range. The first X-ray FEL, the LCLS at SLAC in the USA [2], was going online in 2009 followed by SACLA in Japan, PAL in Korea, European XFEL in Germany and SwissFEL in Switzerland (for more information about the individual sources see e.g. <https://lightsources.org>). Many further FEL projects are presently under consideration all over the world demonstrating the strong impact of X-ray FELs to the present research as well as the large expectations for future research enabled by these powerful facilities.

Beside continuous improvement of the general FEL characteristics, several new features have been developed during the last years. Stable wavelength operation and wavelength tuning are achieved by seeding techniques (e.g. [3]) based on an external sources or on self-seeding schemes. Operation with circularly polarized radiation and variable linear polarization ([3]) has enabled investigations of dichroic

phenomena and two-color lasing at different x-ray energies has been successfully demonstrated at different facilities (e.g. [4]). Experiments combining x-ray pulses with a synchronized optical laser are standard at all facilities and combination with THz radiation produced by the same electron beam allows almost jitter-free operation ([5]). Alternatively, various split-and-delay schemes were developed, which divide optically the FEL pulses and use both parts for x-ray pump – x-ray probe type of investigations. Finally, stability and precision on the attosecond level was realized even for the relative spectral phase between two coherent FEL pulses and used for the demonstration of coherent control in the atomic photoionization [6].

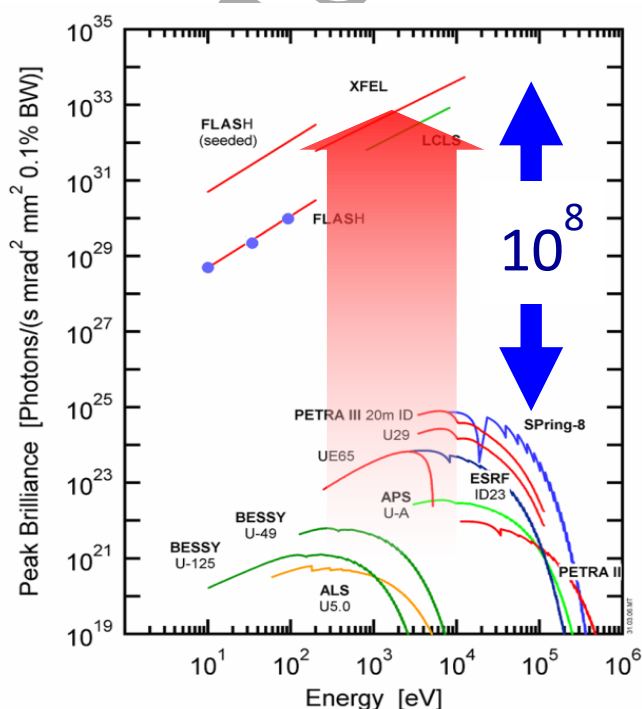


Figure 1: Comparison of peak brilliance for different FEL and synchrotron radiation sources (adapted from [1]).

#### Current and Future Challenges

After the first proof-of-principle experiments demonstrating the outstanding FEL performances, many challenges have still to be met to fulfill the requirements for more sophisticated applications. Since the generation of the FEL pulses is a highly non-linear process, variations in pulse energy, wavelength as well as in the spectral and temporal width are observed between the individual pulses. Especially the SASE process is causing rather non-uniform pulse shapes, which are characterized by a large number of spikes in the spectral and temporal distribution. Stabilization and control of all parameters in combination with their precise online diagnostic for each single pulse is



essential to realize reproducible experimental results and to provide a meaningful basis for the comparison with theoretical models.

Dedicated optical components of highest quality are used for the transport of the x-ray beam to the experimental areas and for providing a micrometer focus to the interaction region with the sample. Pointing stability of the focused beam and preservation of the almost perfect wavefront of the FEL beam has to be assured with the same micrometer precision after transport over distances of several hundreds of meter from the undulator to the experiment. Maintaining these conditions even under the huge head load produced by the irradiation with the intense x-ray pulses is extremely challenging, especially for facilities providing a high repetition rate, such as the European XFEL with 27 000 pulses per second [7].

The generation of ultra-short x-ray pulses has proven to be particularly important for many, not only time-resolved studies, but also imaging and non-linear applications. Few-femtosecond pulses are generated by dedicated bunch compression schemes and single spike lasing can lead even to sub-femtosecond durations. More promising techniques, such as x-ray laser-enhanced attosecond pulse generation (XLEAP), are currently under development and tested at LCLS. As illustration for a possible outcome of XLEAP, the expected performances are given for the SASE3 soft X-ray undulator at the European XFEL [8], providing pulses of 300-500 as width with still remarkable intensities of several hundreds of  $\mu\text{J}$  (see table 1).

Photon energy (eV)	600	870	1000	1500	2300
Undulator cells	6	7	6/7	8	8
Duration (as, FWHM)	500	500	400/550	300	300
Peak power (GW)	600	500	100/650	800	400
Pulse energy ( $\mu\text{J}$ )	350	300	45/350	250	100

Table 1: Photon beam parameters for the SASE3 soft X-ray undulator at the European XFEL when operating in the XLEAP mode [8].

Use of these extremely short x-ray pulses for time-resolved studies in combination with an external optical laser generates strong demands on the temporal stability and reproducibility of the x-ray pulses. Especially for long pulse trains at high repetition rate FELs sophisticated feedback systems were developed providing synchronization of better than 30 fs and promising sub-10 fs in the future [9]. For lower repetition rates, the pulse-to-pulse characterization by THz streaking techniques can be implemented and recent experiments have demonstrated attosecond resolution when additionally angle-resolved electron spectroscopy is applied [10]. Alternative approaches are based on seeding techniques, as demonstrated for

example at FERMI [4], using the same optical laser for the seeding process and for the experiments.

## Advances in Science and Technology to Meet Challenges

Although the technology to generate FEL radiation in the short wavelength regime is relatively new, it is based on a long-standing and highly developed technology of electron beam manipulation in long linear accelerators. Therefore, tremendous advances were achieved during the last years in order to provide for example ultra-short pulses with sub-femtosecond durations via dedicated electron bunch compression schemes, wavelength-stabilized pulses via various seeding schemes, circularly polarized pulses and multi-color lasing as well as the generation of repetition rates in the kHz regime due to the application of super-conducting technology. In addition, techniques which are already well established for table-top lasers are successfully transferred to the X-ray regime such as phase control, chirped-pulse amplification (CPA) and THz streaking techniques.

For the optical transport of the intense radiation, super-polished mirrors have been produced providing shape errors of about 2 nm and corresponding slope errors in the order of 50 nrad for mirrors of 80-100 cm length. The precision to fabricate such mirrors is to a large extent determined by the limitation in measuring the quality of polishing with an adequate high resolution.

The efficient use of higher repetition rates in order to enhance the data acquisition performances and counting statistics is directly related to advances in detector technology, enabling for example recording of several hundreds of high-resolution diffraction patterns during one second (e.g. [7]). Coincidence techniques will strongly benefit from the higher number of X-ray pulses supported by the fact that enormous technical expertise was already acquired during operation at synchrotron radiation sources.

## Concluding Remarks

Science at x-ray free-electron lasers has become a very active and promising research field and many new facilities are presently under development. In addition, all operating facilities have initiated large R&D programs to further improve the performances of the x-ray pulses and to provide them routinely during user experiments. It is clear that this short review can only represent a temporal snapshot of these operations and the selected examples can only illustrate the broad range of activities, but don't claim completeness. There is still a vast potential for further improvements and developments, which can and will be fully exploited in

the future by the fruitful collaboration between experts in accelerator science, X-ray optics and x-ray science.

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## 5.2 Where does ELI-NP stand now?—Two arms of 10PW and 20 MeV gamma beam systems – Kazuo A. Tanaka, Extreme Light Infrastructure: Nuclear Physics.

### Status

ELI-NP has been selected in 2006 by the European Strategy Forum on Research Infrastructures (ESFRI) as making ultra-intense laser fields with intensities reaching up to  $10^{22}$ – $10^{23}$  W cm<sup>-2</sup> and high brilliance gamma ray beam with 20 MeV photon energy possible for nuclear physics. The uniqueness of the facility is to have these laser and gamma beams as stand-alone systems or even to have unusual combination of both laser and gamma beams together for some specific experiments. Since the construction fund has been approved in 2012, the three ELI pillars are being built in Czech Republic, Hungary and Romania. The Romanian pillar is ELI-Nuclear Physics (ELI-NP). The new facility is designed to serve a broad national, European and International science community. Its mission covers scientific research at the frontier of knowledge involving two domains. The first one is high power laser system (HPLS: laser power <10 PW)-driven experiments related to nuclear physics, strong-field quantum electrodynamics and associated vacuum effects [1]. The laser pulse compression has been based on the chirped pulse amplification (CPA) technique for obtaining 25 fsec laser pulse and is to which the Nobel Physics prize has been awarded in 2018 [2]. The second is based on a Compton backscattering high-brilliance and intense low-energy gamma beam system (GBS: gamma beam photon energy <19.5 MeV), a marriage of laser and accelerator technology which will allow us to investigate nuclear structure and reactions as well as nuclear astrophysics with unprecedented resolution and accuracy. In addition to fundamental themes, a large number of applications with significant societal impact are being developed such as medical applications and radio isotope production. The ELI-NP has been located in Măgurele south of Bucharest, Romania. The project is implemented by ‘Horia Hulubei’ National Institute for Physics and Nuclear Engineering (IFIN-HH). The project started in January 2013. The new facility will be fully operational by the end of 2020.

### Current and Future Challenges

The HPLS is being installed on time by Thales French and Thales Romania. A fairly good pulse contrast is expected to have  $10^{13}$  ratio between the main laser pulse to the pre-pulse levels. As of June 2018, 3 PW laser output was demonstrated with a very good Strehl ratio 0.88. This contrast ratio is a key factor to conduct experiments for producing high energy (>200 MeV) protons. The central laser wavelength is 820 nm with the

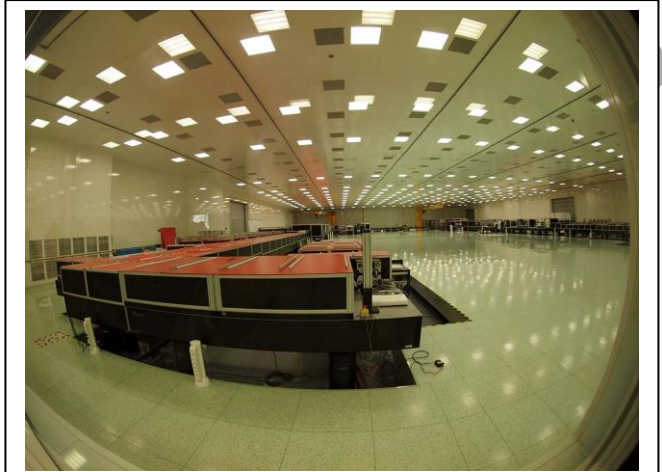


Figure 1 – High Power Laser System (HPLS) of 10 PW 2 arms. The central laser wavelength is 820 nm with the laser bandwidth of 10 nm. The pulse width is <25 fsec with 250 J laser energy output. The focused laser intensity may reach  $10^{23}$  W/cm<sup>2</sup>. The final laser beam size is 50 x 50 cm<sup>2</sup>. HPLS sits on the vibration-dumped floor in the 70 x 40 m<sup>2</sup> clean room. 100 TW or 1 PW laser outputs can be extracted for separate experiments simultaneously with 10 PW output.

laser bandwidth of 10 nm. The pulse width is <25 fsec with 250 J laser energy output. The focused laser intensity may reach  $10^{23}$  W/cm<sup>2</sup>. The final laser beam size is 50 x 50 cm<sup>2</sup>. HPLS sits on the vibration-dumped floor in the 70 x 40 m<sup>2</sup> clean room.

The components of the GBS is being produced by various companies that are organized by EuroGammaS (EGS) and of the low energy section have been delivered to ELI-NP. The GBS is able to provide gamma-ray beams with continuously tunable energy in the range from 200 keV to 19.5 MeV with the number photons of  $10^9$ /sec. Production of quasi-monochromatic gamma beams is based on the Inverse Compton Scattering (ICS) process of laser light pulses off relativistic electron bunches. To generate high-brilliance beams, the solution adopted in the case of ELI-NP was to minimize the interaction volume and maximize the number of photons and electrons within the interaction volume by using low-emittance electron beams and high-brilliance laser pulses. The resulting GBS is based on a high-quality electron–photon collider with luminosity at the level of about  $2 \times 10^{35}$  cm<sup>-2</sup>.s<sup>-1</sup>, that is almost one order of magnitude better than what was achieved at LHC in CERN. The Gamma Beam System at ELI-NP will become the state-of-the-art system in the world delivering beams with photon spectral densities and bandwidth orders of magnitude better than what is today available. 8 independent experimental areas are constructed where users can prepare their experiments independent of the on-going



ones. The experimental chambers have typically the volume of 24 m<sup>3</sup> with 10<sup>-6</sup> mbar vacuum.

## Advances in Science and Technology to Meet Challenges

**High Power Laser System Experiment:** One typical challenge of the possible experiments is the strong QED in the laser electron collisions where 10PW laser pulse will be focused onto the relativistic electron flux created via. the wake field acceleration mechanism through the

$$\eta = \frac{E_{RF}}{E_S} \sim \frac{\gamma m_e c^2}{1 \text{ GeV}} \sqrt{\frac{I}{10^{22} \text{ W cm}^{-2}}} \quad (1)$$

gas jet target [3]. The effect is measured with the ratio of the laser electric field in the electron rest frame:  $E_{RF}$  and the Schwinger field:  $E_S 10^{16} \text{ Vcm}^{-1}$  at which the vacuum breaks into electron-positron pairs,

From this Eq. (1), it is advantageous to achieve high electron energy, namely  $\gamma$  for large  $\eta$ . We expect to have more than GeV electron energy through the wake field acceleration while the focused laser intensity is of the order of  $10^{22} \text{ W/cm}^2$ . A large flux of gamma rays will be created via. the inverse Compton scattering process, but the production mechanism should be affected with the non-linear quantum effect and will be rather different from the classical one.

**Gamma Beam System Experiment:** The NRF technique is a unique way for the investigating low-lying dipole excitations in atomic nuclei and provides a specific research niche for the ELI-NP facility. Gamma ray photons are scattered off bound nuclear states. In particular, the pencil-size bright gamma beams at ELI-NP will provide access to minimum weight targets and will open the actinide region for NRF studies. NRF experiments at present-day facilities require large-mass targets, of the order of grams, while at ELI-NP milligram targets will be used. Thus, the availability frontier will be widen further with the start of the facility. Other studies, which will benefit from the improvement of the availability frontier, are studies of p-process nuclei, which are available in small quantities. At a longer term, parity-doublet studies [4] would be of a high impact, addressing both the sensitivity and precision frontiers. An important part of the NRF program, which will be possible at the high energy beam line, are high-resolution studies of the E1, M1 or E2 strength in the region of the soft collective modes, below the neutron evaporation threshold. Day-one experiments related to photonuclear reactions will also be directed towards the studies of the photodisintegration of <sup>7</sup>Li with the ELISSA array, related to the Big-Bang nucleosynthesis, high-precision studies of photo-

neutron cross-sections (Camera et al 2016). The motivation for studying the <sup>7</sup>Li( $\gamma,t$ )<sup>4</sup>He reaction

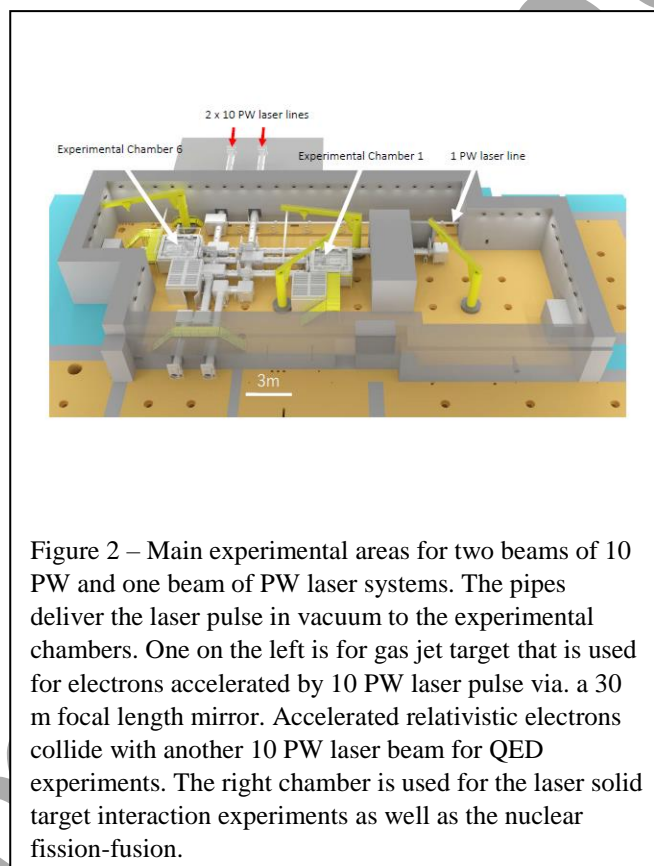


Figure 2 – Main experimental areas for two beams of 10 PW and one beam of PW laser systems. The pipes deliver the laser pulse in vacuum to the experimental chambers. One on the left is for gas jet target that is used for electrons accelerated by 10 PW laser pulse via. a 30 m focal length mirror. Accelerated relativistic electrons collide with another 10 PW laser beam for QED experiments. The right chamber is used for the laser solid target interaction experiments as well as the nuclear fission-fusion.

Serves for not only a final check on one of the reactions that could solve the so-called ‘cosmological Li problem’ but also to resolve a recent disagreement between theoretical models and the experimental data.

## Concluding Remarks

The current status of ELI-NP facility is reported here. The 2 beam HPLS of 10 PW outputs is under construction on time. Recent performance test demonstrated 3 PW output with an indication of very good focusability. The components of the GBS have been delivered to the ELI-NP building. One of the highlight of the GBS is to deliver the highest number of photons after the collimator. Various detectors for measuring gamma rays, neutrons etc. have been produced and tested. Early phase Day 1 experiments will start some time in 2020 including electron acceleration of up to GeV leading to the nonlinear quantum QED experiments with HPLS. NRF and photonuclear reaction experiments with <sup>7</sup>Li will be planned using GBS.

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