

1.

TAFZ
Summer 2017

Born-Oppenheimer approximation

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role of the approximation is to separate electronic and nuclear motions. The idea is based on the electronic and nuclear mass difference.

- Non-relativistic molecular Hamiltonian (N electrons, M nuclei) (atomic units)

$$H = \underbrace{-\frac{1}{2} \sum_{i=1}^N \nabla_i^2 - \sum_{i=1}^N \sum_{A=1}^M \frac{Z_A}{|\vec{r}_i - \vec{R}_A|}}_{H_{el}(\vec{r}; \vec{R})} + \underbrace{\sum_{i>j} \frac{1}{|\vec{r}_i - \vec{r}_j|} + \sum_{A>B} \frac{Z_A Z_B}{|\vec{R}_A - \vec{R}_B|} - \sum_{A=1}^M \frac{1}{2M_A} \nabla_A^2}_{T_N(\vec{R})}$$

$$\vec{r} \equiv \{\vec{r}_1, \dots, \vec{r}_N\} ; \vec{R} \equiv \{\vec{R}_1, \dots, \vec{R}_M\}$$

$$T_N = \sum_{\alpha, \kappa} \frac{1}{2M_A} P_{\alpha\kappa} P_{\alpha\kappa} ; P_{\alpha\kappa} = -i \frac{\partial}{\partial R_{\alpha\kappa}}$$

$\alpha \in \{x, y, z\}$

- $H_{el}(\vec{r}; \vec{R})$ is fixed (clamped)-nuclei ~~Hamiltonian~~ electronic Hamiltonian

1.) $H_{el} \phi_k(\vec{r}; \vec{R}) = E_k(\vec{R}) \phi_k(\vec{r}; \vec{R})$

2.) Normalization $\langle \phi_k(\vec{r}; \vec{R}) | \phi_{k'}(\vec{r}; \vec{R}) \rangle_{\vec{r}} = \delta_{kk'}$... independent of \vec{R}

$$\Rightarrow P_{\alpha\kappa} \langle \phi_k(\vec{r}; \vec{R}) | \phi_k(\vec{r}; \vec{R}) \rangle_{\vec{r}} = 0 = \underbrace{\langle P_{\alpha\kappa}^+ \phi_k | \phi_k \rangle_{\vec{r}} + \langle \phi_k | P_{\alpha\kappa} \phi_k \rangle_{\vec{r}}}_{\langle \phi_k | P_{\alpha\kappa}^+ \phi_k \rangle_{\vec{r}} = \langle \phi_k | P_{\alpha\kappa} \phi_k \rangle_{\vec{r}} \text{ (real } \phi_k)} \left. \vphantom{\langle \phi_k | P_{\alpha\kappa}^+ \phi_k \rangle_{\vec{r}}} \right\} \langle \phi_k | P_{\alpha\kappa} \phi_k \rangle = 0$$

3.) $H \Psi = E \Psi$

$$\Psi(\vec{r}, \vec{R}) = \sum_k \phi_k(\vec{r}; \vec{R}) \chi_k(\vec{R})$$

Expansion is exact

$$\int d\vec{r} \phi_k^* (H_{el} + T_N) \sum_k \phi_k \chi_k = E \sum_k \phi_k \chi_k$$

$$(E_k(\vec{R}) - E) \chi_k + \sum_{k'} \langle \phi_k | T_N | \phi_{k'} \chi_{k'} \rangle_{\vec{r}} = 0$$

$$[E_k(\vec{R}) - E] \chi_k + \sum_{k'} \sum_{\alpha, \kappa} \frac{1}{2M_A} P_{\alpha\kappa} P_{\alpha\kappa} \chi_{k'} \delta_{kk'} + 2 \underbrace{\langle \phi_k | P_{\alpha\kappa} \phi_{k'} \rangle_{\vec{r}}}_{-U_{kk'}} P_{\alpha\kappa} \chi_{k'} + \langle \phi_k | P_{\alpha\kappa} P_{\alpha\kappa} \phi_{k'} \rangle_{\vec{r}} \chi_{k'} = 0$$

$$(1) \quad (T_N + E_k(\vec{R}) + U_{kk} - E) \chi_k(\vec{R}) = \sum_{k' \neq k} \left[\sum_{\alpha, \kappa} \frac{1}{M_A} \langle \phi_k | P_{\alpha\kappa} | \phi_{k'} \rangle_{\vec{r}} P_{\alpha\kappa} + \sum_{\alpha, \kappa} \frac{1}{2M_A} \langle \phi_k | P_{\alpha\kappa} P_{\alpha\kappa} \phi_{k'} \rangle_{\vec{r}} \right] \chi_{k'}(\vec{R})$$

BO diagonal correction

non-adiabatic coupling (1st order)

second order

(2.)

- Estimate for the second order:

$$\nabla_A \phi_k(\vec{r}; \vec{R}) \sim \nabla_i \phi_k(\vec{r}; \vec{R}) \dots \text{similar differences over the same distance scale}$$

$$\sim p_e \phi_k(\vec{r}; \vec{R}) \dots \text{electronic momentum } p_e$$

$$\Rightarrow \frac{1}{2M_A} \nabla_A^2 \phi_k(\vec{r}; \vec{R}) \sim \frac{p_e^2}{2M_A} \phi_k(\vec{r}; \vec{R}) \sim E_{el} \left(\frac{m_e}{M_A} \right) \phi_k(\vec{r}; \vec{R})$$

$\sim 10^4$

Born - Oppenheimer approximation (1927)

$$\left(T_N + E_L(\vec{R}) - E \right) \chi_k(\vec{R}) = 0$$

Nuclei move in the effective potential $E_L(\vec{R})$

For each $\phi_k(\vec{r}; \vec{R})$ we have $\chi_k(\vec{R})$ and $\Psi(\vec{r}; \vec{R}) = \phi_k(\vec{r}; \vec{R}) \chi_k(\vec{R})$

Is it always valid? When does it fail. Do we have some signature property of its failure?

$$\langle \phi_k | P_{A\alpha} \phi_k \rangle = E_k \langle \phi_k | P_{A\alpha} \phi_k \rangle$$

$$\begin{cases} \langle \phi_k | P_{A\alpha} \text{Hel} | \phi_{k'} \rangle = E_{k'} \langle \phi_k | P_{A\alpha} | \phi_{k'} \rangle \\ \langle \phi_k | \text{Hel} P_{A\alpha} | \phi_{k'} \rangle = E_k \langle \phi_k | P_{A\alpha} | \phi_{k'} \rangle \end{cases} \left\{ \langle \phi_k | P_{A\alpha} | \phi_{k'} \rangle = \frac{\langle \phi_k | [P_{A\alpha}, \text{Hel}] | \phi_{k'} \rangle}{E_{k'}(\vec{R}) - E_k(\vec{R})} \right.$$

Can we evaluate $\langle \phi_k | [P_{A\alpha}, \text{Hel}] | \phi_{k'} \rangle$?

$$1.) \langle \phi_k | P_{A\alpha} \frac{1}{R_{AB}} | \phi_{k'} \rangle - \langle \phi_k | \frac{1}{R_{AB}} P_{A\alpha} | \phi_{k'} \rangle = \langle \phi_k | \left(P_{A\alpha} \frac{1}{R_{AB}} - \frac{1}{R_{AB}} P_{A\alpha} \right) | \phi_{k'} \rangle = 0 \text{ for } k \neq k'$$

$$2.) \langle \phi_k | \left[P_{A\alpha}, - \sum_{i=1}^{N_e} \frac{z_A}{|\vec{r}_i - \vec{R}_A|} \right] | \phi_{k'} \rangle = i z_A \langle \phi_k | \sum_{i=1}^N \frac{(\vec{r}_i - \vec{R}_A)}{|\vec{r}_i - \vec{R}_A|^3} | \phi_{k'} \rangle$$

non-divergent, non-zero

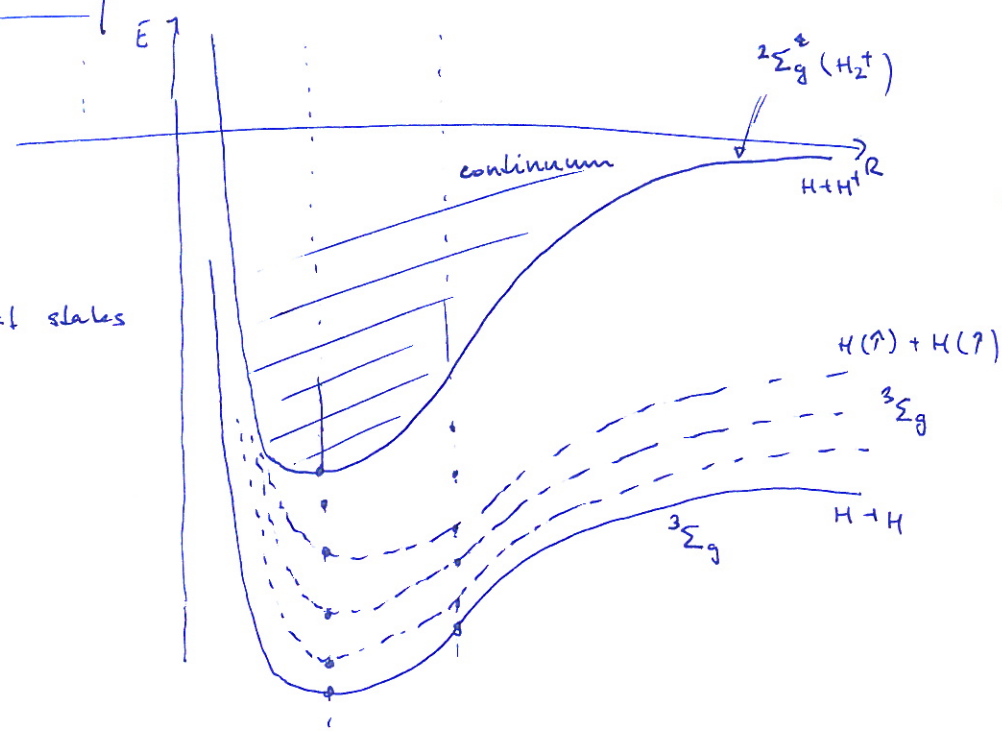
(3)

TAF2
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Born - Oppenheimer approximation

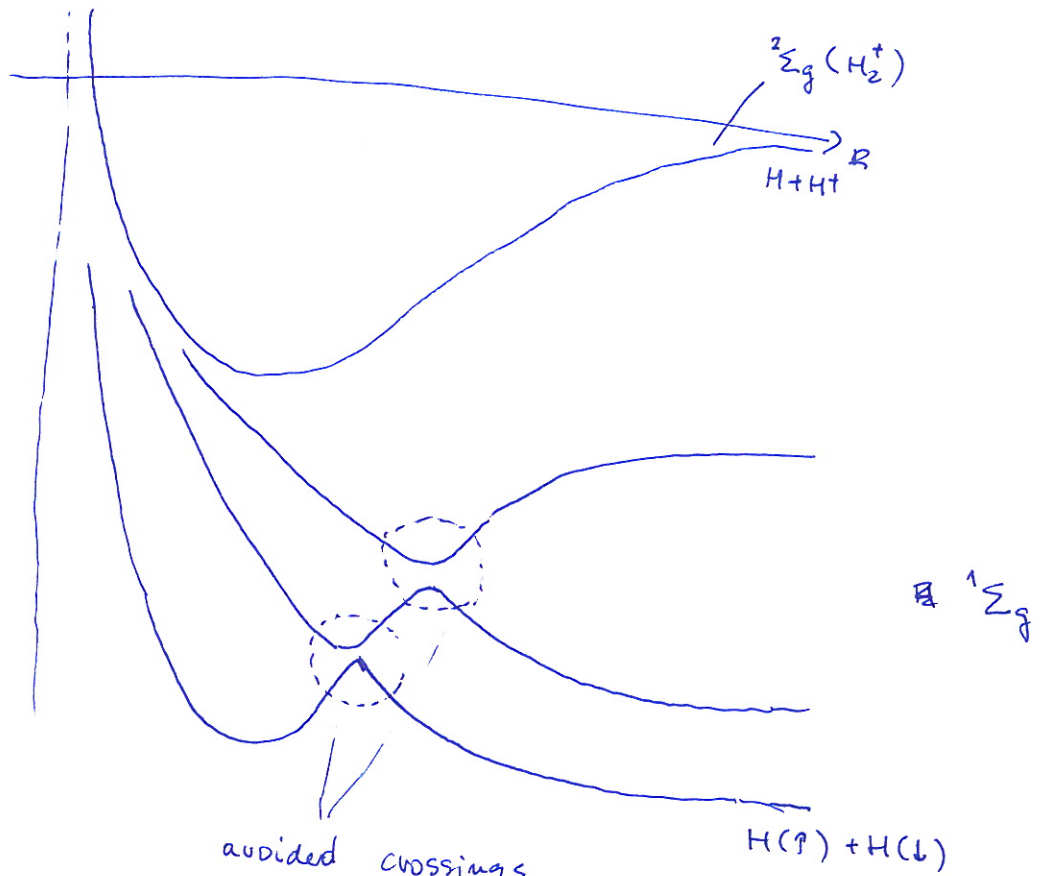
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H₂ triplet states



~~doublet~~
doublet
 $S = \frac{1}{2}$
② $\Sigma_g^- \dots$
m=0
gerade
(inversion from center)

H₂ singlet gerade states



avoided crossings
BOA breaks down and full vibronic coupled system (1) must be solved

